

Interlinks of dynamical thermalization

Kolmogorov turbulence, KAM and Anderson localization

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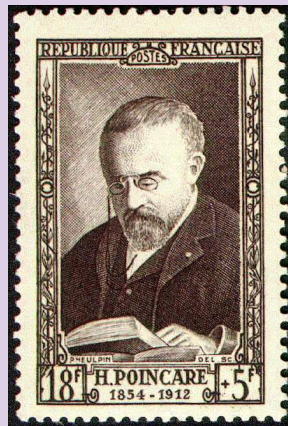
It is not so much important to be rigorous as to be right
A.N.Kolmogorov (from Chirikov Phys Rep (1979))



* Boltzmann - Loschmidt dispute:
irreversible statistical laws from reversible dynamical equations

Sitzungsberichte der Akademie der Wissenschaften, Wien,
II **73**, 128 (1876); **75**, 67 (1877)

Dynamical chaos (mathematics, physics)



- * Poincaré (1893);
- * Chirikov resonance overlap criterion for plasma experiments at Kurchatov Inst. (1959)

Chirikov talk (celebrating his 70th at Toulouse)

CHAOS in SIBERIA

A story for history

Short (40') account of a long (40 yrs) work
Boris Chirikov & many coworkers in
Budker Institute of Nuclear Physics
& elsewhere

My (accidental) start
($\approx$ 1958, Kurchatov Institute, Moscow)

a simple-looking but turned-out rich
Budker's problem: single particle confinement in
Budker's (adiabatic) magnetic trap
(for the great END – controlled nuclear fusion !)

Toulouse
16 July 1998

BACKGROUND (retrospectively)

Revival of intensive studies
into nonlinear dynamics,
surprising rediscovery of chaos (stochasticity) after
Boltzmann ... Poincare ...

- new applications:
strong focusing accelerators,
controlled nuclear fusion
- computers, NUMERICAL EXPERIMENTS !!

– Lehmer, 1951
pseudorandom number generators after
Galton Board (= Lorentz gas in external field !)
– Goward and Hine (CERN), 1953, accelerators

- ➔ Fermi, Pasta and Ulam, 1955
foundations of statistical mechanics
- Symon and Sessler, 1956, accelerators
- ➔ Kolmogorov, 1954, KAM-integrability
(in spite of Poincare theorem !)

Chirikov talk (celebrating his 70th at Toulouse)

instructive
examples

SIMPLE CHAOS:
most complex dynamics
BUT simple statistics

• MICROTRON, Veksler, 1944

$$E_{n+1} = E_n + V_0 \sin \phi_n$$

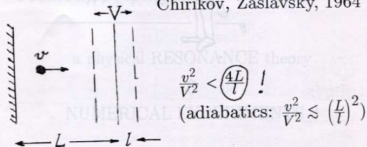
$$\phi_{n+1} = \phi_n + E_{n+1}/E_0, \quad F = 2 \text{ freedoms}$$

$$K = V_0/E_0 \gg 1$$

$$E_0 = \omega_0/2\pi\Omega, \quad e=m=c=1$$

• FERMI ACCELERATION, Ulam, 1961

Chirikov, Zaslavsky, 1964



• COMET HALLEY

diffusion backwards in time

$$t_L \sim 10 \text{ Myrs} \quad ?$$

Chirikov, Vecheslavov, 1989

Fermi's problem (1923)

ergodicity of nonlinear oscillation

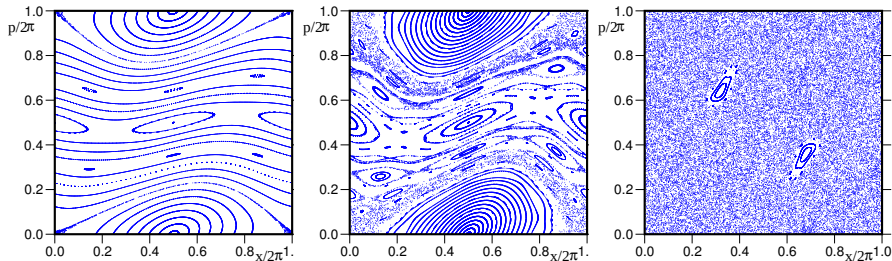
$$H = \sum_{i=1}^N \frac{\dot{x}_i^2}{2} + \frac{(x_i - x_{i-1})^2}{2} + (x_i - x_{i-1})^p, \quad p=3;4, \quad N \gg 1$$

Fermi, Pasta, Ulam, 1955

nonergodic (generally) $\rightarrow$ solitons Kruskal,
Miura,
Zabusky, 1966
 $\rightarrow$ dynamical chaos
Chirikov, Izrailev, 1966

De Luca, Lichtenberg, Lieberman, 1995
Shepelyansky, 1997

Chirikov standard map (1969, 1979)



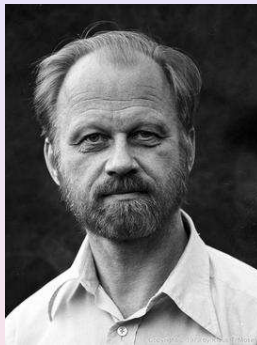
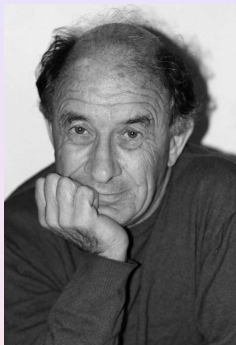
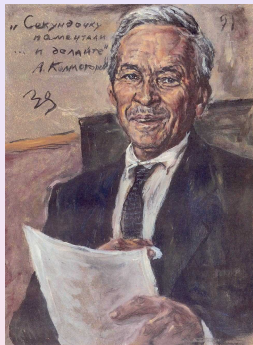
$$\bar{p} = p + K \sin x, \quad \bar{x} = x + K \sin x; \quad K = 0.5; 0.971635; 5$$

related systems: Frenkel-Kontorova model (1937); Veksler microtron (1944); Ulam map for Fermi acceleration (1961), Kepler map for microwave ionization of Rydberg atoms, Halley comet, dark matter dynamics in the Solar system ... (see Scholarpedia 3(3):3550 (2008))

Diffusion in momentum and energy $\langle p^2 \rangle = Dt \approx K^2 t/2$ for $K > 0.971635...$ ($E = p^2/2$)

Chirikov, DS Scholarpedia (2008-20017)

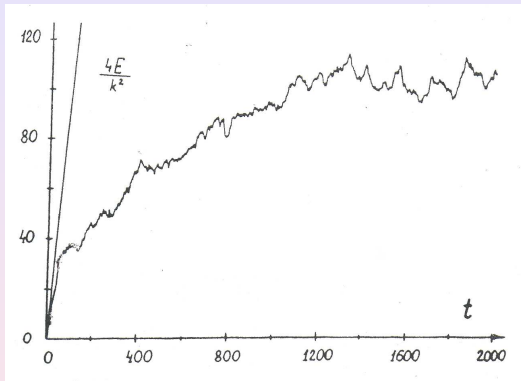
Kolmogorov-Arnold-Moser theory (1954-1963)



Chirikov criterion (1959) (e.g. for Chirikov standard map $K < K_c$)
term “KAM theory” coined by Chirikov (INP preprint 1969)

Quantum localization of chaos

Quantum Chirikov standard map (kicked rotator)



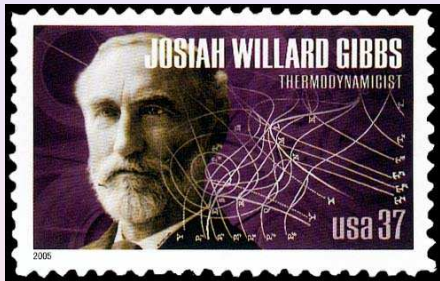
Chirikov, Izrailev, DS (1981) → Raizen cold atom experiment (1995):

$\bar{\psi} = \exp(-ik \cos x) \exp(-iTn^2/2)\psi$; $k = K/\hbar$, $T = \hbar$, $n = -i\partial/\partial x$;

localization length: $\ell \approx \Delta n \approx D \sim K^2/\hbar^2$

Analogy with Anderson localization (AL) in disordered solids (1958; Nobel prize 1977) shown by Fishman, Grempel, Prange (1982)

Thermalization (classical/quantum)

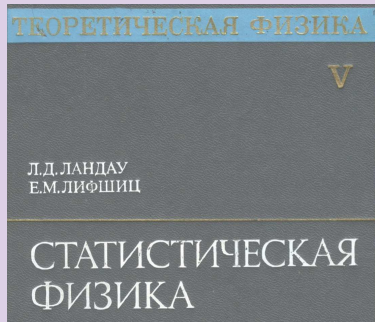


J.W.Gibbs (1874-1878) ;



M.Planck (1900) $\rightarrow \hbar$

Energy equipartition over degrees of freedom



§ 44]

ЗАКОН РАВНОРАСПРЕДЕЛЕНИЯ

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Соответственно теплоемкость $c_p = c_v + 1$ равна

$$c_p = \frac{l+2}{2}. \quad (44,2)$$

Таким образом, чисто классический идеальный газ должен обладать постоянной теплоемкостью. Формула (44,1) позволяет при этом высказать следующее правило: на каждую переменную в энергии $\varepsilon(p, q)$ молекулы приходится по равной доле $1/2$ в теплоемкости c_v газа ($k/2$ в обычных единицах), или, что то же, по равной доле $T/2$ в его энергии. Это правило называют *законом равнораспределения*.

Имея в виду, что от поступательных и вращательных степеней свободы в энергию $\varepsilon(p, q)$ входят только соответствующие им импульсы, мы можем сказать, что каждая из этих степеней свободы вносит в теплоемкость вклад, равный $1/2$. От каждой же колебательной степени свободы в энергию $\varepsilon(p, q)$ входит по две переменных (координата и импульс), и ее вклад в теплоемкость равен 1.

Nauka Moscow (1976)

Dynamical thermalization in finite isolated systems?

Fermi (1923); Fermi-Pasta-Ulam (FPU) problem (1955)

Chirikov criterion for onset of chaos (1959)

Novosibirsk => FPU: Izrailev, Chirikov Dokl. Akad. Nauk SSSR 166: 57 (1966)

Integrability of nonlinear Schrödinger equation

Zakharov, Shabat, Zh. Eksp. Teor. Fiz. 61: 118 (1971) + Toda lattice (1967)

Fermi-Pasta-Ulam problem (1955)

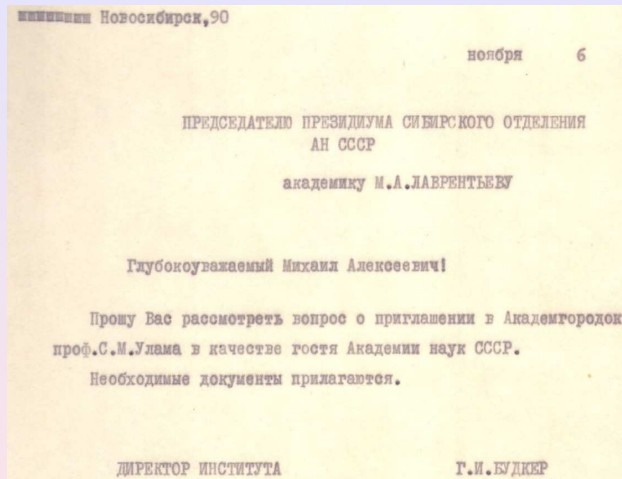


"In January 1951, Ulam and Teller came up with the Teller-Ulam design, which is the basis for all thermonuclear weapons. ...

"After the H-bomb was made," Bethe recalled, "reporters started to call Teller the father of the H-bomb. For the sake of history, I think it is more precise to say that Ulam is the father, because he provided the seed, and Teller is the mother, because he remained with the child. As for me, I guess I am the midwife." "

Wikipedia

Virtual visit of Ulam to Akademgorodok (1967)



invitation Ulam never got

Meeting of Chirikov and Ulam at Mathematical Congress, Moscow 1966

Invitation of Ulam to Novosibirsk for summer 1967 (letter of Budker to Lavrentiev Nov 1966 (initiated by Chirikov)) from Archive of Chirikov

<http://www.quantware.ups-tlse.fr/chirikov/archive/ulamprog.pdf>

Ulam method (1960) and Ulam networks (2010)

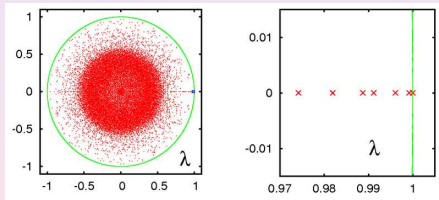
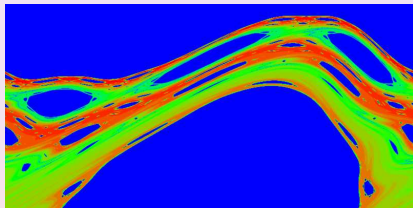
Outside note:

Ulam method for the Chirikov standard map $\rightarrow$ Markov chains, Google matrix

World Wide Web, PageRank algorithm $\rightarrow$ Brin, Page (1998);

Ulam networks $\rightarrow$ DS, Zhirov; Ermann, DS (2010);

Chirikov standard map $\rightarrow$ Frahm, DS (2010)



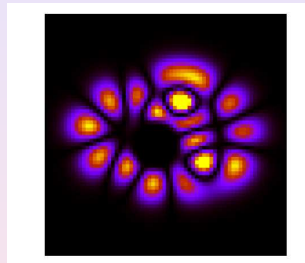
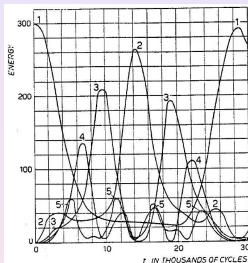
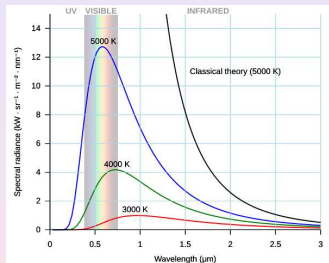
Google matrix analysis of directed networks: WWW, world trade, Wikipedia ...

PageRank, CheiRank

Chepelianski (2010);

Ermann, Frahm, DS Rev Mod Phys (2015)

Dynamical thermalization in finite systems



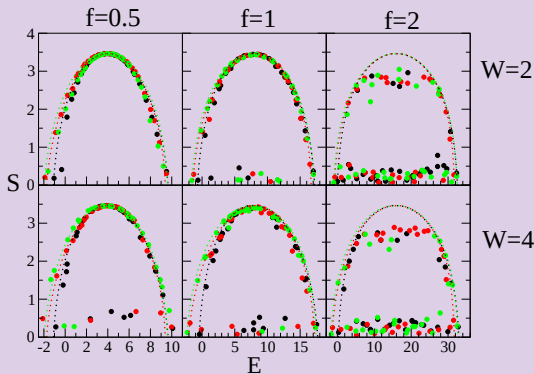
Planck's law (1900); Fermi-Pasta-Ulam problem (1955);
Bose-Einstein condensate (BEC) in Sinai-oscillator trap (2016) (left to right)
Ketterle group PRL (1995) $\Rightarrow$ BEC in Sinai oscillator (not understood; 3d)

Quantum Gibbs ansatz instead of equipartition

DANSE model (1D) $i\hbar\partial\psi_n/\partial t = E_n\psi_n + \beta|\psi_n|^2\psi_n + V(\psi_{n+1} + \psi_{n-1})$

$$H = \sum_{m=1}^N \epsilon_m |C_m|^2 + \beta \sum_{m_1 m_2 m_3 m} V_{m_1 m_2 m_3 m} C_{m_1} C_{m_2} C_{m_3}^* C_m^*$$

with $\sum_m |C_m|^2 = 1$, $V_{mm'm_1m'_1} \sim \ell^{-3/2}$, $E_n = f|n - n_0| + \delta E_n$, $|\delta E_n| < W/2$



Quantum Gibbs ansatz: $\rho_m = Z^{-1} \exp(-\epsilon_m/T)$, $Z = \sum_m \exp(-\epsilon_m/T)$
 $S(E) = -\sum_m \rho_m \ln \rho_m$; $\rho_m = |C_m|^2$ (Ermann, DS NJP (2013))

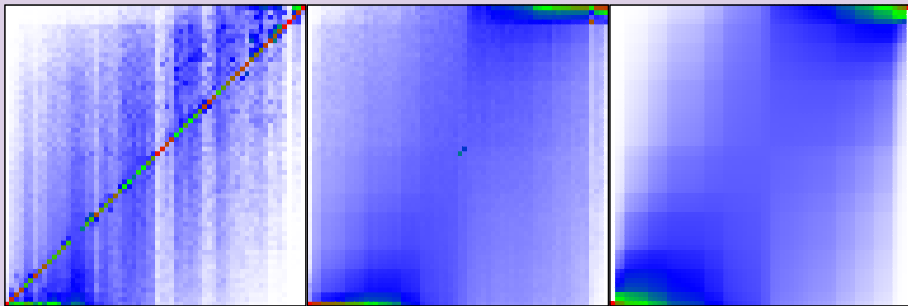
Dynamical thermalization in nonlinear chains

$$i \frac{\partial \psi_{n_x n_y}}{\partial t} = E_{n_x n_y} \psi_{n_x n_y} + \beta |\psi_{n_x n_y}|^2 \psi_{n_x n_y} + (\psi_{n_x+1 n_y} + \psi_{n_x-1 n_y} + \psi_{n_x n_y+1} + \psi_{n_x n_y-1}). \quad (6)$$

Periodic boundary conditions are used for the $N \times N$ square lattice with $-N/2 \leq n_x, n_y \leq N/2$. However, here we use the extended version of this model assuming that

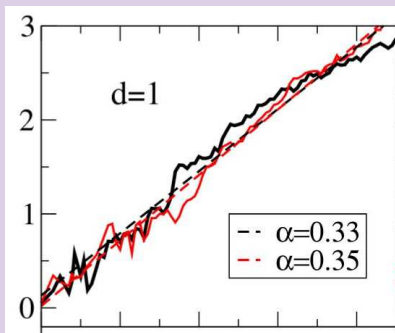
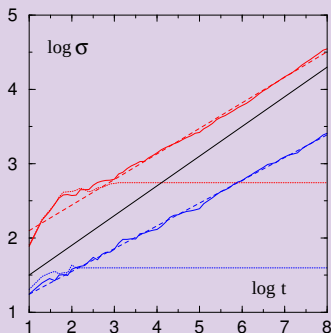
$$E_{n_x n_y} = \delta E_{n_x n_y} + f(n_x^2 + n_y^2), \quad -W/2 \leq \delta E_{n_x n_y} \leq W/2 \quad (M3). \quad (7)$$

This is the *M3* model with random values of energies $\delta E_{n_x n_y}$ in a given interval.



Nonlinear chains with disorder: 8×8 sites, $f = 1$, $W = 2$, $\beta = 1, 4$ (left, center); quantum Gibbs ansatz (right); $t = 2 \times 10^6$ (Ermann, DS NJP (2013))

KAM and spreading in infinite nonlinear chains



Left: DANSE at disorder $W = 2$ (red), 4 (blue); $\beta = 1, 0; f = 0, t \leq 10^8$;
 Right: kicked rotator with nonlinear phase shift $t \leq 10^9$; random (black) and $Tn^2/2$ (red) rotational phases, $\beta = 1$ (log-log scale)

DS (1993; 2012); Pikovsky, DS (2008); Ermann, DS JPA (2014)
 Flach *et al.* JPA (2014)

Spreading: $\sigma = \langle n^2 \rangle \propto t^\alpha$, $\alpha \approx 0.34 \approx 2/5 < 1/2$ (random phases)

KAM border for $\beta \ll \beta_c \sim 1$?

J.Bourgain, W.-M.Wang J. Eur. Math. Soc. (2008)

Bose-Einstein condensates (BEC) with cold atoms



BEC dynamics => Gross-Pitaevskii equation (GPE)

Ketterle BEC experiment (1995))

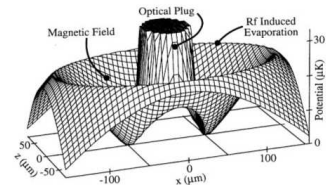
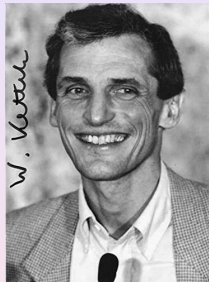


FIG. 1. Adiabatic potential due to the magnetic quadrupole field, the optical plug, and the rf. This cut of the three-dimensional potential is orthogonal to the propagation direction (y) of the blue-detuned laser. The symmetry axis of the quadrupole field is the z axis.

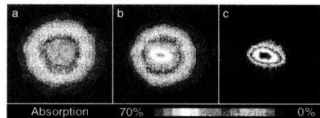


FIG. 2 (color). Two-dimensional probe absorption images, after 6 ms time of flight, showing evidence for BEC. (a) is the velocity distribution of a cloud cooled to just above the transition point, (b) just after the condensate appeared, and (c) after further evaporative cooling has left an almost pure condensate. (b) shows the difference between the isotropic thermal distribution and an elliptical core attributed to the expansion of a dense condensate. The width of the images is $870\ \mu\text{m}$. Gravitational acceleration during the probe delay displaces the cloud by only 0.2 mm along the z axis.

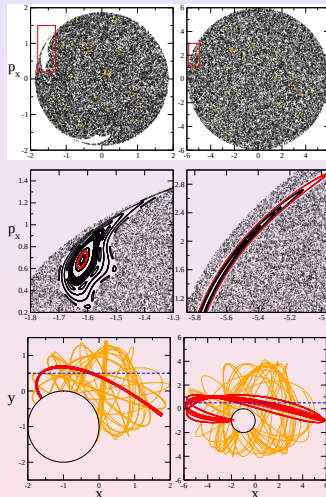
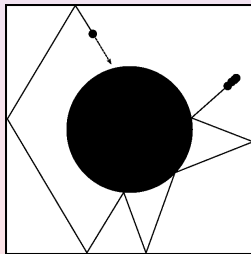
PRL (1995) (Nobel prize 2001)

Sinai billiard $\rightarrow$ Sinai oscillator for BEC

Sinai billiard (1963; 1970) (Abel prize 2014)

$$H = (p_x^2 + p_y^2)/2 + x^2/2 + y^2$$

$$r_d = 1, x_d = y_d = -1; E = 2; 18 \rightarrow$$



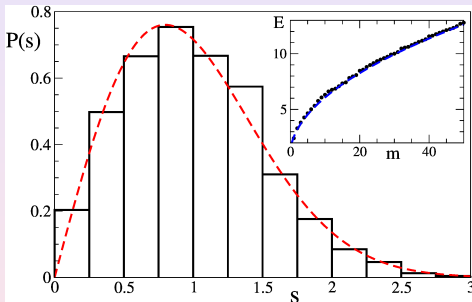
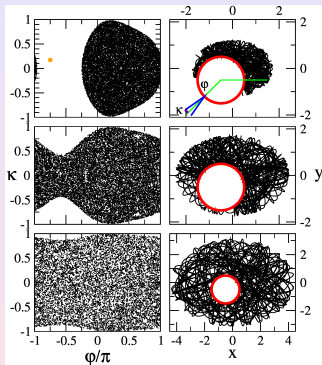
BEC dynamics in Sinai oscillator with GPE

Ermann, Vergini, DS (2015-2017)

Classical and quantum chaos in Sinai oscillator

Poincaré sections (left)

$H = (p_x^2 + p_y^2)/2 + x^2/2 + y^2$; $r_d = 1$; $\hbar = 1$; $x_d = y_d = -0.5$; $E = 1.5; 3; 10$;



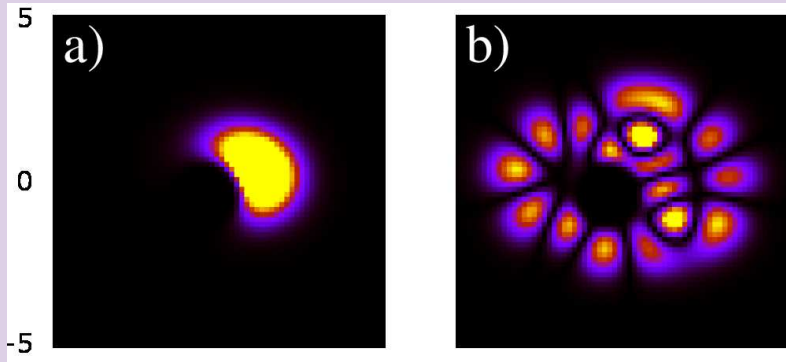
Bohigas-Giannoni-Schmit conjecture (1984); Ullmo Scholarpedia (2016)

Wigner-Dyson statistics of lowest 2500 energy levels unfolded (right)

Random matrix theory (Wigner (1967)); quantum chaos (e.g. Haake (2010))

Ermann, Vergini, DS (2015-2017)

Quantum chaos eigenstates in Sinai oscillator



Eigenstates at $\beta = 0$; ground state $m = 1$ and $m = 24$

Bose-Einstein THERMALIZATION anzats:

$$\rho_m = 1 / [\exp(E_m - E_g - \mu) / T] - 1];$$

$$\rho_m = \langle |\psi_m|^2 \rangle, \text{ energy } \sum_m E_m \rho_m = E,$$

$$\text{entropy } S = - \sum_m \rho_m \ln \rho_m \rightarrow S(E)$$

BEC in Sinai oscillator with GPE

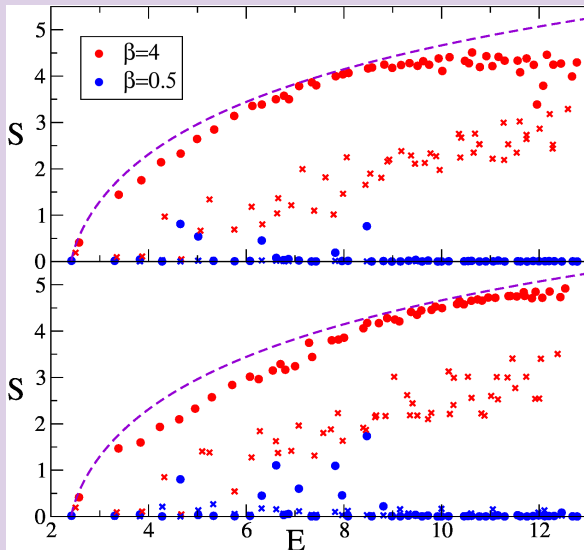
Gross-Pitaevskii equation (GPE or NSE) (Pitaevskii, Stringari (2003))

The BEC evolution in the Sinai oscillator trap is described by the GPE, which reads as

$$\begin{aligned} i\hbar \frac{\partial \psi(\vec{r}, t)}{\partial t} = & -\frac{\hbar^2}{2m} \vec{\nabla}^2 \psi(\vec{r}, t) \\ & + \left[\frac{m}{2} (\omega_x^2 x^2 + \omega_y^2 y^2) + V_d(x, y) \right] \psi(\vec{r}, t) \\ & + \beta |\psi(\vec{r}, t)|^2 \psi(\vec{r}, t). \end{aligned} \quad (2)$$

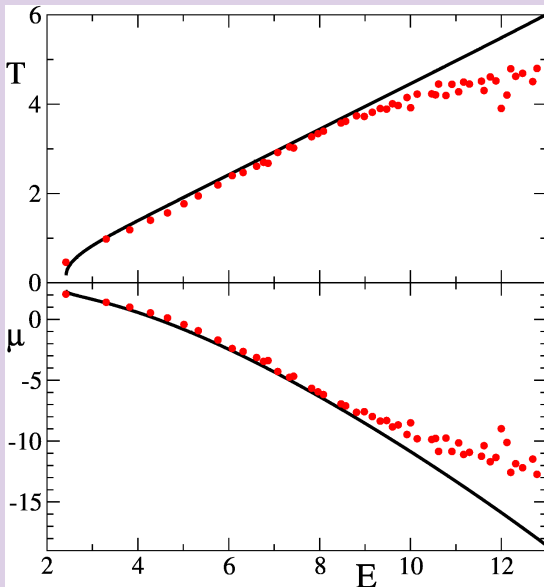
Here in (2), we use the same oscillator and disk parameters as in (1) and take $\hbar = 1$. The wave function is normalized to unity $W = \int |\psi(x, y)|^2 dx dy = 1$. Then, the parameter β describes the nonlinear interactions of atoms in BEC. All

Bose-Einstein anzats for dynamical thermalization



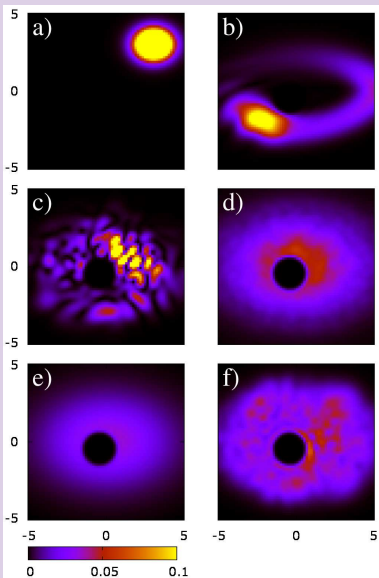
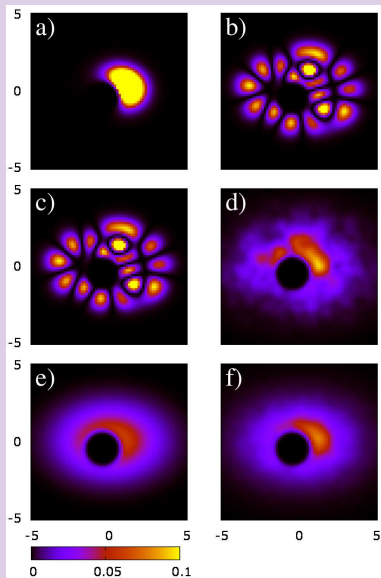
first 50 states; Sinai osc (dots), no disk (X); $500/1500 < t < 1500/2500$
(top/bottom); Bose-Einstein anzats (dashed) → no energy equipartition

Bose-Einstein anzats



temperature and chemical potential dependence on energy ($\beta \Rightarrow 4$)

BEC time evolution



$\beta = 4$ various initial states

Kolmogorov turbulence defeated by Anderson localization

532.507

ЛОКАЛЬНАЯ СТРУКТУРА ТУРБУЛЕНТНОСТИ В НЕСЖИМАЕМОЙ ВЯЗКОЙ ЖИДКОСТИ ПРИ ОЧЕНЬ БОЛЬШИХ ЧИСЛАХ РЕЙНОЛЬДСА *)

А. Н. Колмогоров

§ 1. Будем обозначать через

$$U_{\alpha}(P) = U_{\alpha}(x_1, x_2, x_3, t), \quad \alpha = 1, 2, 3,$$

компоненты скорости в момент времени t в точке с прямоугольными декартовыми координатами x_1, x_2, x_3 . При изучении турбулентности естественно считать компоненты скорости $U_{\alpha}(P)$ в каждой точке $P = (x_1, x_2, x_3, t)$

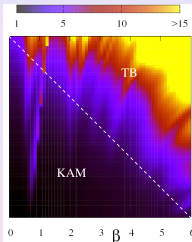
*) ДАН СССР 30 (4), 299 (1941).

V.E. Zakharov V.S. Lvov
G. Falkovich

Kolmogorov Spectra of Turbulence I Wave Turbulence

With 34 Figures

Springer-Verlag
Berlin Heidelberg New York
London Paris Tokyo
Hong Kong Barcelona
Budapest



Kolmogorov (1941) - energy flow over space scales $E_k \propto 1/k^{5/3}$;

NSU $\rightarrow$ Zakharov-Filonenko (1967) capillary waves $E_k \propto 1/k^{7/4} \rightarrow$:

“In the theory of weak turbulence nonlinearity of waves is assumed to be small; this enables us, using the hypothesis of the random nature of the phase of individual waves, to obtain the kinetic equation for the mean square of the wave amplitudes.”

extensions $\rightarrow$ Zakharov, L’vov, Falkovich(1992); Nazarenko(2011)

BEC in Sinai-oscillator trap with driving

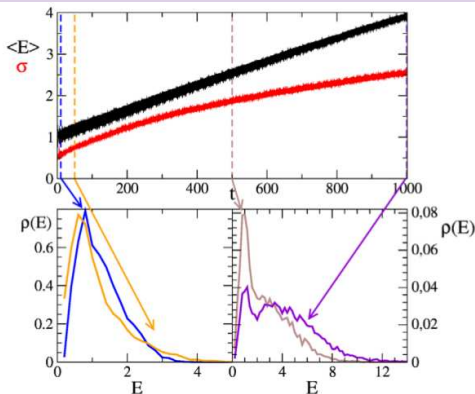


FIG. 1. Classical time evolution of average energy $\langle E \rangle$ and its standard deviation σ for $f = 0.4$. The data are obtained from 10^4 trajectories with random initial conditions at $\langle E \rangle = 1$ and $\sigma = 0.5$. Top panel: $\langle E(t) \rangle$ and $\sigma(t)$ are shown by black and red (gray) curves, respectively. Bottom panels show probability distribution of trajectories $\rho(E, t)$ for (a) $t = 10, 50$ [blue (black),

$$H = H_S(x, y, p_x, p_y) + fx \sin \omega t \text{ (classical dynamics)}$$

Energy flow to high modes

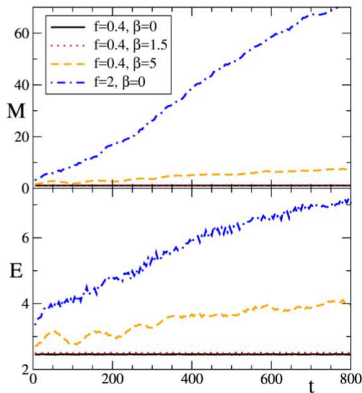


FIG. 2. Time evolution of M (top panel) and energy E (bottom panel) for GPE (2) averaged over time intervals $\Delta t = 1$. The initial state is the ground state of (2) at $\beta = 0, f = 0$ [see Fig. 5(a) in Ref. [28]]. Both panels show the cases of $f = 0.4, \beta = 0$ (black solid lines), $f = 0.4, \beta = 1.5$ [red (gray) dotted lines], $f = 0.4, \beta = 5$ [orange(gray) dashed lines], $f = 2, \beta = 0$ [blue (gray) dot-dashed lines].

$M = \sum_k k \rho_k; \ell_\phi \approx 2\pi \rho_c (D/\omega)^2 \approx 2f^2 \omega_x^2 E^{3/2} / \omega^4$ (ground state);

Anderson photonic localization $\rho_k \propto \exp(-2E/\omega \ell_\phi)$ for $\beta = 0, f < f_c \approx 1.5$

Energy flow to high modes

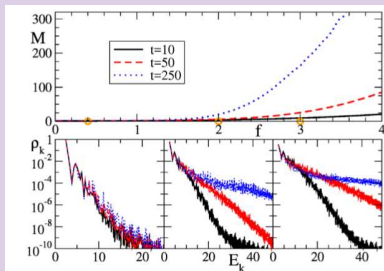


FIG. 3. Top panel shows M as a function of driven force f for linear case ($\beta = 0$). Bottom panels show probability distribution ρ_k , averaged over time interval $\Delta t = 5$, as a function of eigenenergies E_k with $t = 10$ in black solid lines, $t = 50$ in the red (gray) dashed lines, and $t = 250$ in blue (gray) dotted lines. Left, center, and right bottom panels show the cases of $f = 0.4, 2$, and 3 , respectively [highlighted with orange (gray) circles in top panel].

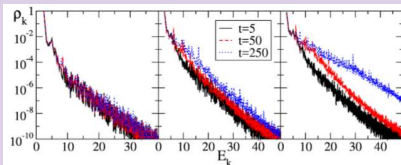


FIG. 4. Same as in bottom panels of Fig. 3 for $f = 0.3, \beta = 1.5$ (left panel); $f = 0.5, \beta = 5$ (center panel); $f = 1, \beta = 5$ (right panel).

M and probability distributions ρ_k ; left $\rightarrow \beta = 0$;
right $f = 0.3, \beta = 1.5, f = 0.5, \beta = 5, f = 1, \beta = 5$

Turbulence phase diagram

$$f_c r_d / \hbar \omega_x \approx 1.5 [1 - \beta_c / (6 \hbar \omega_x r_d^2)] \quad (3)$$

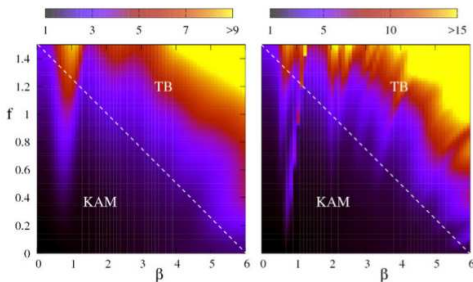
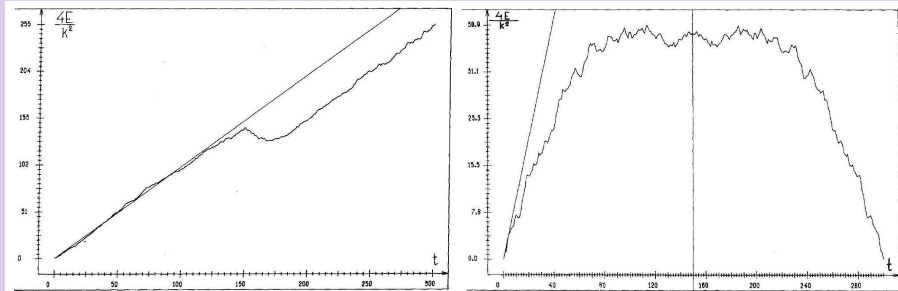


FIG. 5. Number of modes M shown by color (grayness) in the plane of parameters f and β (average is done in the time intervals $100 \leq t \leq 150$ and $250 \leq t \leq 300$ in left and right panels, respectively). The approximate separation of KAM or insulator phase (KAM) and delocalized turbulent or metallic phase (TB) is shown by the white line (3).

Thus there is a stability domain where the Kolmogorov flow from large to small scales is defeated by the Anderson localization and KAM-integrability

Quantum version of Boltzmann-Loschmidt dispute

* Time reversal for the Chirikov standard map



BESM-6 computation, rescaled energy or squared momentum vs. time t :

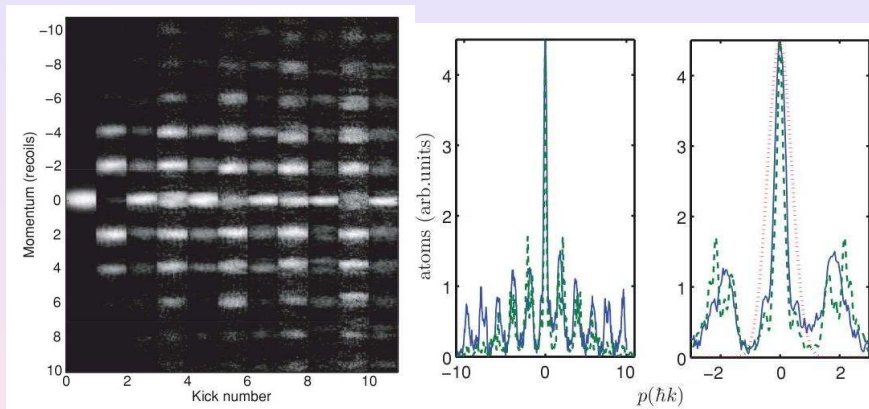
$K = 5$, $\hbar = 0$ (left), $\hbar = 1/4$ (right), DS (1983)

proposal for BEC in kicked optical lattice (Martin, Georgeot, DS (2008))

$k = K/\hbar$, $\hbar = T = 4\pi + \epsilon$ (forward), $4\pi - \epsilon$ (back);

$\cos x \rightarrow -\cos x$ by a π shift in x

Time reversal of atomic matter waves by Hoogerland group (2011)

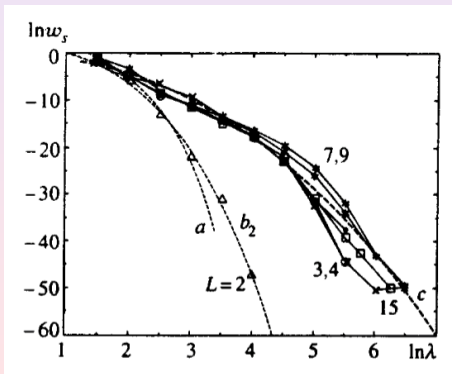


- $k = 2 - 3$, $T = 4\pi \pm \epsilon$ (Ullah, Hoogerland (2011))
- Ultracold ^{67}Rb atoms BEC:** 10^4 atoms, $T_{\text{cool}} = 50\text{nK}$,
 $\lambda = 2\pi/k_L = 760\text{nm}$, $\epsilon = 1$; 5 + 5 kicks; right panel shows zoom near
 initial distribution shown by red dotted curve (initial/final width is
 0.43/0.21 recoils; full/dashed curve for experiment/numerics).

Discussion (open problems)

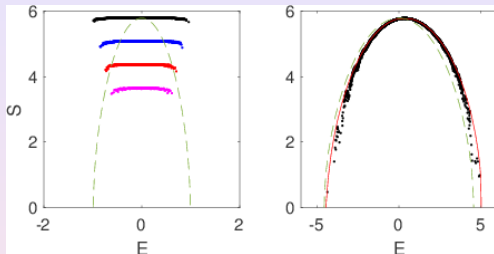
Classical systems with many freedoms:

- Quantum Gibbs instead of equipartition
- KAM border for pure-point spectrum of Anderson localization
- fast Arnold diffusion with many freedoms (fig, $\lambda = 1/\sqrt{K}$)
 $D \sim w_s^2 \propto K^{6.6}$ coupled standard maps (Chirikov, Vecheslavov (1997))
chaos measure of triplet resonances $\mu \sim K$ (Mulansky et al (2011))

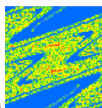


Discussion (open problems)

Quantum systems with many freedoms:



- Dynamical thermalization in quantum many-body systems:
Åberg criterion PRL (1990); Jacquod, DS (1997)
Two-body random interaction model (TBRIM) for
quantum dots and SYK black holes ..., (Kolovsky, DS (2017))
quantum dot $\rightarrow \delta E \approx g^{2/3} \Delta$; SYK black hole $\rightarrow ?$
- Time reversal of atoms, waves, sound ...

[illegible]

I
 $K \sim S^2 = \left(\frac{\Delta \omega r}{\Omega d} \right)^2 > 1$
 (1959)

II
 $\beta E_0 > N/k$ /FPU/
 (1966/73)

III
 $\bar{I} = I + K \sin \theta$ $\bar{\theta} = \theta + \bar{I}$
 $> K_c = \frac{\pi}{4} S^2 = 1$; $h = \ln(K/2) > 0$
 /KAM/ (1969/79)

IV
 $= 4\pi e \lambda^3 \exp(-\pi \lambda/2)$, $\lambda = \Omega/\omega_0$
 $D_A = D_0 \exp(-\pi \lambda) \sim \omega_s^2$
 (1969/79)

V
 $\bar{\psi} = \exp(-i \frac{Tn^2}{2}) \exp(-ik \cos \hat{\theta}) \psi$
 $(\Delta n)^2 \sim k^2 t^4 < \infty$ /KR/
 (1979)

$$K - S^2 = \left(\frac{\Delta \omega_r}{\Omega_d} \right)^2 > 1 \quad (1959)$$

II	(1966/73)
$\beta E_0 > N/k$	/FPU/
III	(1966/73)

III (1966/73)
 $\bar{I} = I + K \sin \theta, \quad \bar{\theta} = \theta + \bar{I}$
 $K > K_c = \frac{\pi^2}{4} S^2 = 1; \quad h = \ln(K/2) > 0$
 /KAM/ /KS/ (1969/79)

$$\begin{aligned} \text{IV} \\ \omega_s = 4\pi\epsilon\lambda^3 \exp(-\pi\lambda/2), \quad \lambda = \Omega/\omega_0 \\ D_A = D_0 \exp(-\pi\lambda) \sim \omega_s^2 \\ (196g/7g) \end{aligned}$$

$$\bar{\psi} = \exp(-i\frac{\pi n^2}{2}) \exp(-ik \cos \hat{\theta}) \psi$$

$$(\Delta n)^2 \sim k^2 t^* < \infty / KR /$$

(1979)

VI
 $h > \Lambda_m = 0.38H^{1/4} > 0$
 /YM/ (1981)

VII
 $P(\tau) \sim 1/\tau^p$, $p \approx 1.5$ (1981)
 $p = 3$ ($\tau \rightarrow \infty$) (1998)

$$t^* \sim \Delta n \sim D \sim k^2 \gg t_E \sim \ln k / h \quad (1981)$$

IX
 $\varepsilon_q = 0.5 \omega_0^{7/6} / \sqrt{n_0} \cdot H /$
 (1984)

$\bar{\omega} = \omega + F(x)$, $\bar{x} = x + \bar{\omega}^{-5/2}$
 $h > 0$, $t_D \approx 4 \times 10^6$ yrs
 /CH/ (1989)

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