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Luigi Accardi

**Characterization of product states on
polynomial algebras in terms of scalar
products of n -particle vectors**

based on joint work with:

Abdon Ebang Ella and Yun Gang Lu

Email: accardi@volterra.uniroma2.it

WEB page: <http://volterra.uniroma2.it>

Introduction

Quadratic quantization

The RSWN n -particle scalar product is given by

$$\begin{aligned}
 \left\langle \phi^{\otimes n}, \psi^{\otimes n} \right\rangle_n &= \tag{1} \\
 \sum_{i_1+2i_2+\dots+ki_k=n} &\frac{n!\beta^n c_2^{i_1+\dots+i_k} \langle \phi, \psi \rangle^{i_1} \dots \langle \phi^k, \psi^k \rangle^{i_k}}{i_1! \dots i_k! \beta^{i_1} (2\beta)^{i_2} \dots (k\beta)^{i_k}} \\
 &= c_2^n \langle \phi, \psi \rangle^n + \\
 + \sum_{\substack{i_1 \neq n \\ i_1+2i_2+\dots+ki_k=n,}} &\frac{n!\beta^n c_2^{i_1+\dots+i_k} \langle \phi, \psi \rangle^{i_1} \dots \langle \phi^k, \psi^k \rangle^{i_k}}{i_1! \dots i_k! \beta^{i_1} (2\beta)^{i_2} \dots (k\beta)^{i_k}}
 \end{aligned}$$

The complex polynomial algebra in d Hermitean indeterminates

$$\mathbb{C}[X_1, \dots, X_d] \quad ; \quad (d \in \mathbb{N} \cup \{\infty\})$$

is a well studied object in mathematics since hundreds years.

For reasons that will be clear later, we want to give a **concrete representation** of this \ast -algebra **in terms of functions** on a real vector space.

The algebra \mathcal{P}_V of polynomial functions on a real vector space V

V a **real** vector space.

V^* algebraic **dual** of V

natural embedding $V \hookrightarrow (V^*)^*$

$$X : v \in V \rightarrow X_v \in (V^*)^*$$

defined by

$$X_v : u^* \in V^* \rightarrow X_v(u^*) := \langle u^*, v \rangle \in \mathbb{R} \quad (2)$$

$X_v(u^*) := \langle u^*, v \rangle :=$ **coordinate** of u^* along v

Thus we have the **coordinate maps**

$$X_v : V^* \rightarrow \mathbb{R}$$

The n -th powers of the X_v are the **monomial functions** (or simply **monomials**) on V of degree n

$$X_v^n : u^* \in V^* \rightarrow (X_v(u^*))^n \in \mathbb{R} \quad ; \quad v \in V, n \in \mathbb{N}$$

$$\mathcal{P}_V := \mathbb{C}\text{-linear span of monomials}$$

$$:= \text{polynomial algebra on } V$$

\mathcal{P}_V is a ***-algebra** for the pointwise operations.

Involution: complex conjugation.

X_v is real-valued as a function \Rightarrow is **Hermitean**
as element of \mathcal{P}_V :

$$X_v^* = X_v \quad ; \quad \forall v \in V$$

Classical random fields as vector valued random variables

Denote \mathcal{F}^* the σ -algebra on the real vector space V^* generated by the coordinate functions $\{X_v : V^* \rightarrow \mathbb{R} ; v \in V\}$.

If on \mathcal{F}^* it is given a **probability measure** φ , we obtain a **probability space**

$$(V^*, \mathcal{F}^*, \varphi)$$

and the coordinate functions

$$X_v : u^* \in V^* \rightarrow X_v(u^*) := \langle u^*, v \rangle \in \mathbb{R}$$

become real valued random variables on $(V^*, \mathcal{F}^*, \varphi)$. Therefore the map

$$X : v \in V \rightarrow X_v \tag{3}$$

can be interpreted as a

V -valued classical random variable.

or equivalently as a

real-valued classical random field on V .

If $V \equiv \mathbb{R}^d$ is **finite dimensional**, one usually identifies V with V^* fixing a linear basis on V . Through this identification, any probability measure φ **becomes a probability measure on V** .

Classical random fields with all moments

Definition 1 A classical random field X on V is said to have **finite moments of any order** if:

$$\varphi(|X_v|^n) < \infty \quad ; \quad \forall v \in V, \forall n \in \mathbb{N}$$

$$\varphi(|X_v|^n) := \int_{V^*} |X_v(\omega)|^n \varphi(d\omega)$$

Same symbol φ for measure and φ -integral.

From now on:

– **random field (or variable) on $V \equiv$**
 \equiv random field (or variable) on V with all
moments.

– We fix V (real vector space)

$$\mathcal{P} \equiv \mathcal{P}_V$$

States and pre-scalar products on \mathcal{P}_V

The restriction of the φ -integral on \mathcal{P}_V gives a **state** φ on \mathcal{P}_V

$$\varphi(Q) := \int Q d\varphi \quad ; \quad Q \in \mathcal{P}_V$$

\iff

- φ is positive on positive elements in \mathcal{P} .
- $\varphi(1) = 1$.
- φ defines a **pre-scalar product** on \mathcal{P}_V :

$$\langle P, Q \rangle := \varphi(P^*Q) \quad ; \quad Q \in \mathcal{P}_V$$

The pair

$$(\mathcal{P}_V, \varphi)$$

is a standard example of a **classical algebraic probability space**. **Classical** because the $*$ -algebra \mathcal{P}_V is **commutative**.

Theorem 1 For a semi-scalar product $\langle \cdot, \cdot \rangle$ on \mathcal{P} the following statements are **equivalent**:
 (i) There exists a state φ on \mathcal{P} such that:

$$\varphi(f^*g) = \langle f, g \rangle \quad ; \quad f, g \in \mathcal{P} \quad (4)$$

(ii) The semi-scalar product $\langle \cdot, \cdot \rangle$ satisfies

$$\langle 1_{\mathcal{P}}, 1_{\mathcal{P}} \rangle = 1 \quad (5)$$

and, for each $v \in V$, multiplication by the coordinate X_v is a **Hermitean** linear operator on \mathcal{P} with respect to $\langle \cdot, \cdot \rangle$, i.e.:

$$\langle X_v f, g \rangle = \langle f, X_v g \rangle \quad (6)$$

Taking the **cyclic representation** of the pair

$$(\mathcal{P}_V, \varphi) \mapsto (\mathcal{H}_\varphi, \eta_\varphi, \pi_\varphi, \Phi_0)$$

one obtains:

- a Hilbert space $\mathcal{H}_\varphi \equiv [\mathcal{P}_V \cdot \Phi_0]$,
- a quotient map $\eta_\varphi : \mathcal{P}_V \rightarrow \mathcal{H}$,
- a $*$ –representation π_φ of \mathcal{P}_V ,
- a $\pi_\varphi(\mathcal{P}_V)$ –cyclic vector (**vacuum vector**)

$$\Phi_0 := \eta_\varphi(1_{\mathcal{P}_V}) = \eta_\varphi(\text{constant function } = 1)$$

Simplified notations.

we **omit the representation symbol** $\pi_\varphi, \eta_\varphi$
and, for $Q \in \mathcal{P}_V$, we use the notations:

$$Q(X) = \pi_\varphi(Q) \text{ multiplication } \mathbf{operator}$$

$$Q \cdot \Phi_0 = \eta_\varphi(Q) = \pi_\varphi(P) \cdot \Phi_0 \text{ vector}$$

Note that

$$\mathcal{P}_V \cdot \Phi_0 := \{P \cdot \Phi_0 : P \in \mathcal{P}_V\} \subseteq \mathcal{H}$$

is a dense sub-space invariant under the action of $\pi_\varphi(\mathcal{P}_V)$.

Its closure

$$\mathcal{H}_\varphi = [\mathcal{P}_V \cdot \Phi_0] =: L^2_{pol}(V^*, \varphi) \subseteq L^2(V^*, \varphi)$$

The degree filtration in \mathcal{P}_V
 is the family of sub-spaces of \mathcal{P}_V : $n \in \mathbb{N}$

$$\mathcal{P}_{V,n]} := \quad (7)$$

$:= \{\text{linear span of **monomials of degree** } \leq n\}$

Closed if $\dim(V) = +\infty$.

It is increasing in the sense that

$$\mathcal{P}_{V,n]} \subset \mathcal{P}_{V,n+1]} \subset \mathcal{P}_V \quad ; \quad \forall n \in \mathbb{N} \quad (8)$$

moreover

$$\mathcal{P}_V = \bigcup_{n \in \mathbb{N}} \mathcal{P}_{V,n]} \quad (9)$$

Closed if $\dim(V) = +\infty$.

Since V is fixed, we write $\mathcal{P}_n], \mathcal{P}, \dots$

instead of $\mathcal{P}_{V,n]}, \mathcal{P}_V, \dots$

Abstract roots of generalized quantizations

Theorem 2 Let \mathcal{H} be an Hilbert space and let be given:

– an **increasing filtration** (sequence) $(\mathcal{H}_{n}]$ of closed sub-spaces of \mathcal{H} with corresponding orthogonal projectors

$$P_{n]} : \mathcal{H} \rightarrow \mathcal{H}_{n]} \quad ; \quad n \in \mathbb{N}$$

– a set D is and, for each $j \in D$, a linear Hermitean operator

$$Y_j = Y_j^* : \mathcal{H} \rightarrow \mathcal{H}$$

filtration increasing of degree $+1$

$$Y_j \left(\mathcal{H}_{n]} \right) \subseteq \mathcal{H}_{n+1]} \quad (10)$$

Then the sequence of mutually orthogonal projectors:

$$P_n := P_{n]} - P_{n-1]} \quad ; \quad \forall n \in \mathbb{N} \quad (11)$$

$$(P_{-1} := P_{-1]} := 0) \quad (12)$$

satisfies the 3-**diagonality relation**

$$P_m Y_j P_n = 0 \quad \text{if} \quad m \notin \{n-1, n, n+1\} \quad (13)$$

If, in addition

$$\lim_{n \rightarrow \infty} P_{n]} = 1_{\mathcal{H}} \quad (14)$$

in any operator topology, then (P_n) is a **partition of the identity** in the same topology and:

$$Y_j P_n = P_{n+1} Y_j P_n + P_n Y_j P_n + P_{n-1} Y_j P_n \quad (15)$$

or equivalently:

$$Y_j = \sum_{n \in \mathbb{N}} P_{n+1} Y_j P_n + \sum_{n \in \mathbb{N}} P_n Y_j P_n + \sum_{n \in \mathbb{N}} P_{n-1} Y_j P_n \quad (16)$$

In particular, defining the operators

$$\begin{cases} a_{Y_j}^+ := \sum_{n \in \mathbb{N}} P_{n+1} Y_j P_n \text{ (**creation**)} \\ a_{Y_j}^- := \sum_{n \in \mathbb{N}} P_{n-1} Y_j P_n \text{ (**annihilation**)} \\ a_{Y_j}^0 := \sum_{n \in \mathbb{N}} P_n Y_j P_n \text{ (**preservation**)} \end{cases} \quad (17)$$

(16) becomes:

$$Y_j = a_{Y_j}^+ + a_{Y_j}^0 + a_{Y_j}^- \quad ; \quad j \in D \quad (18)$$

Corollary 1 The following are equivalent:

(i) To give a triple

$$(\mathcal{H}, (P_n]_{n \in \mathbb{N}}, (Y_j)_{j \in D})$$

satisfying the conditions of Theorem (2).

(ii) To give:

(ii.1) an **orthogonal Hilbert space gradation**

$$(\mathcal{H}, \langle \cdot, \cdot \rangle) = \bigoplus_{n \in \mathbb{N}} (\mathcal{H}_n, \langle \cdot, \cdot \rangle_n) \quad (19)$$

(ii.2) for each $j \in D$ an **adjointable operator on \mathcal{H} gradation increasing** of degree $+1$, i.e. such that

$$a_j^+ : \mathcal{H}_n \rightarrow \mathcal{H}_{n+1} \quad (20)$$

(ii.3) for each $j \in D$ an **Hermitean gradation preserving** operator on \mathcal{H} (i.e. of degree 0):

$$(a_j^0)^* = a_j^0 : \mathcal{H}_n \rightarrow \mathcal{H}_n \quad (21)$$

Definition 2 Let D be a set. A **generalized Fock quantization with 1-particle space** $V \sim \mathbb{R}^{|D|}$ is given by;

$$(\mathcal{H}_n, \langle \cdot, \cdot \rangle_n)_{n \in \mathbb{N}} \quad ; \quad (a_j^+)_{j \in D} \quad ; \quad (a_j^0)_{j \in D}$$

Satisfying the conditions of Corollary 1.

- \mathcal{H}_n is called **generalized n -particle space**;
- The $(a_j^+)_{j \in D}$ are called **generalized creators**;
- The $(a_j^0)_{j \in D}$ are called **generalized preservation operators**;
- The adjoints of the creators

$$(a_j^+)^* := a_j^- : \mathcal{H}_n \rightarrow \mathcal{H}_{n-1} \quad (22)$$

are called **generalized annihilators**.

They are gradation decreasing of degree -1 .

- The **CAP operators** is the union of the 3 sets above.

– The identity

$$Y_j = a_j^+ + a_j^0 + a_j^- \quad ; \quad j \in D \quad (23)$$

is called **the quantum decomposition of** Y_j associated to the triple $(\mathcal{H}, (P_n]_{n \in \mathbb{N}}, (Y_j))_{j \in D}$. The equivalent formulation of the identity (23) given by

$$Y_j P_n = P_{n+1} Y_j P_n + P_n Y_j P_n + P_{n-1} Y_j P_n \quad ; \quad n \in \mathbb{N} \quad (24)$$

is called the **symmetric 3–diagonal (Jacobi) relation** associated to the triple $(\mathcal{H}, (P_n]_{n \in \mathbb{N}}, (Y_j))_{j \in D}$.

Theorem 3 The quantum decomposition (23) **is unique** up to unitary isomorphism.

Type I commutation relations

The products of CAP operators

$$a_{Y_j} a_{Y_k}^+ \quad , \quad a_{Y_k}^+ a_{Y_j} \quad , \quad j, k \in D$$

are well defined on \mathcal{D} , **gradation preserving**.
In particular satisfy the following
commutation relations:

$$[a_{Y_j}, a_{Y_k}^+] = \partial\Omega_{jk} \quad , \quad j, k \in D \quad (25)$$

where $(\partial\Omega_{jk})$ is a kernel with values into the gradation preserving adjointable linear operators $\mathcal{H} \rightarrow \mathcal{H}$.

Definition 3 The commutation relations (25) are called **type I commutation relations**.
(The reasons of this terminology will be clear in the following.)

Remark. Type I commutation relations only arise from the fact that both $a_{Y_j} a_{Y_k}^+$ and $a_{Y_k}^+ a_{Y_j}$ are **gradation preserving**:

Product states on polynomial algebras

Let V be a real topological vector space and let $e \equiv (e_n)_{n \in D}$ be a linear basis of V (D an at most countable set). Then

$$V = \dot{\sum}_{n \in D} \mathbb{C} \cdot e_n$$

where $\dot{\sum}_{n \in D}$ denotes algebraic direct sum of vector spaces.

Moreover for each $n \in D$,

$$\mathcal{P}_{\mathbb{C} \cdot e_n} \subseteq \mathcal{P} \equiv \mathcal{P}_V \quad ; \quad m \neq n \Rightarrow \mathcal{P}_{\mathbb{C} \cdot e_m} \cap \mathcal{P}_{\mathbb{C} \cdot e_n} = 1_{\mathcal{P}}$$

$$\mathcal{P} = \bigvee_{n \in D} \mathcal{P}_{\mathbb{C} \cdot e_n} = \text{lin. span} \left\{ \prod_{n \in D} q_{\mathbb{C} \cdot e_n} : \right.$$

$$\left. q_{\mathbb{C} \cdot e_n} \in \mathcal{P}_{\mathbb{C} \cdot e_n} , \quad q_{\mathbb{C} \cdot e_n} = 1_{\mathcal{P}_{\mathbb{C} \cdot e_n}} \quad \text{almost all } n \in D \right\}$$

Definition 4 A state φ on \mathcal{P}_V , is called a **product state** if there exists a linear basis $e \equiv (e_n)_{n \in D}$ of V such that, denoting

$$\varphi_n := \varphi \Big|_{\mathcal{P}_{\mathbb{C} \cdot e_n}} \quad (\text{restriction}) \quad , \quad n \in D$$

one has

$$\varphi \left(\prod_{n \in F} p_{\mathbb{C} \cdot e_n} \right) = \prod_{n \in F} \varphi_n(p_{\mathbb{C} \cdot e_n})$$

for any finite sub-set $F \subseteq I$ and any family of polynomials $p_{\mathbb{C} \cdot e_n} \in \mathcal{P}_{\mathbb{C} \cdot e_n}$ ($n \in F$).

The Boson canonical commutation relations

Recall the type I commutation relations

$$[a_v, a_u^+] = \partial\Omega(v, u) \quad , \quad j, k \in D$$

The simplest gradation preserving operators are the **multiples of identity**:

$$\partial\Omega(v, u) = \partial\omega(v, u) \cdot 1$$

with $\partial\omega(v, u) \in \mathbb{C}$.

In this case **the Type I CCR** become:

$$\partial\Omega(v, u) = [a_v^-, a_u^+] = \partial\omega(v, u) \cdot 1 \quad (26)$$

Hence their vacuum expectation value is

$$0 \leq \|a_v^+ \Phi_0\|^2 = \langle \Phi_0, \partial\omega(v, v) \Phi_0 \rangle = \partial\omega(v, v)$$

Therefore the sesqui-linear form

$$\langle v, u \rangle_{V_{\mathbb{C}}} := \partial\omega(v, u)$$

is positive definite and therefore it defines a **semi-scalar product** $\langle \cdot, \cdot \rangle_{V_{\mathbb{C}}}$ on $V_{\mathbb{C}}$.

Theorem 4 Let V be a real vector space and let $(v, u) \in V_{\mathbb{C}} \times V_{\mathbb{C}} \mapsto \partial\omega(v, u)$ be a non-identically zero semi-scalar product on $V_{\mathbb{C}} \equiv V + iV$ (the complexification of V) taking real values on $V \times V$.

For any classical random field X on V the following statements are equivalent.

(i) The CAP operators of X satisfy $a_v^0 = 0$, $\forall v \in V$ (i.e. X is symmetric) and

$$\partial\Omega(v, u) := [a_v^-, a_u^+] = \partial\omega(v, u) \cdot 1 \quad (27)$$

(ii) The quantum field canonically associated to X is the Fock Boson field over the Hilbert space $(V_{\mathbb{C}, \partial\omega}, \langle \cdot, \cdot \rangle_{\partial\omega})$, obtained by completing $V_{\mathbb{C}}$ with the semi-scalar product $\partial\omega$.

(iii) X is the standard classical Gaussian field on the real Hilbert space $(V_{\partial\omega}, \langle \cdot, \cdot \rangle_{\partial\omega_V})$, obtained by completing V with the semi-scalar product $\partial\omega_V$ where $\partial\omega_V$ is the restriction of $\partial\omega$ on V .

Remark. Note that the semi-scalar product $\langle \cdot, \cdot \rangle_{\partial\omega_V}$, on V , is **real valued** because φ is real on real polynomials.

The theorem above follows from the following more general result.

Product measures:

Statistics implies algebra

The canonical commutation relations of Boson quantum field theory (CCR)

$$[a_v^-, a_u^+] = \langle v, u \rangle_{V_{\mathbb{C}}} \cdot 1$$

imply in particular that

$$v \perp u \Rightarrow [a_v^-, a_u^+] = 0$$

Therefore, the canonical commutation relations of Boson quantum field theory enjoy the following property:

Property (C)

If $e \equiv (e_j)_{j \in D}$ is an orthogonal basis of $(V_{\mathbb{C}}, \langle \cdot, \cdot \rangle_{V_{\mathbb{C}}})$, one has:

$$[a_{e_j}^-, a_{e_k}^+] = 0 \quad ; \quad \forall j \neq k \quad (28)$$

The following question then arises naturally:
For which classical V -valued random variables X does there exist a linear basis $e \equiv (e_j)_{j \in D}$ of $V_{\mathbb{C}}$ for which condition (28) is satisfied?

Theorem 5 (Accardi, Kuo, Stan)

For a state φ on \mathcal{P} associated to a classical V -valued random variable X and a linear basis $e \equiv (e_j)_{j \in D}$ of $V_{\mathbb{C}}$, **the following statements are equivalent:**

(i) The classical real-valued random variables

$$X_{e_j} = a_{e_j}^+ + a_{e_j}^0 + a_{e_j}^- \quad ; \quad j \in D \quad (29)$$

are **mutually independent**.

(ii) φ is a **product state** with respect to the linear basis $e \equiv (e_j)_{j \in D}$ of $V_{\mathbb{C}}$.

(iii) The CAP operators in the quantum decomposition of the X_{e_j} **mutually commute**, i.e. for all $j \neq k \in D$, and all $\varepsilon, \eta \in \{+1, 0, -1\}$, one has

$$[a_{e_j}^{\varepsilon}, a_{e_k}^{\eta}] =: [a_j^{\varepsilon}, a_k^{\eta}] = 0 \quad (30)$$

Properties of product states

From Definition 4 it follows a state is a product state if and only if its restriction to an increasing sequence of finite dimensional sub-spaces is a product state.

This reduces the problem to the finite dimensional case.

Therefore from now on we suppose that for $d \in \mathbb{N}^*$,

$$D := \{1, \dots, d\}$$

Thus V is a real d -dimensional vector space, $e \equiv \{e_j\}_{j \in D}$ a linear basis of V

$$D_n := \left\{ \bar{m} := (m_1, \dots, m_d) \in \mathbb{N}^d : \sum_{j \in D} m_j = n \right\}$$

and φ a state on \mathcal{P}_V .

We denote

$$\langle a \cdot \Phi_0, b \cdot \Phi_0 \rangle := \varphi(a^*b)$$

the semi-scalar product defined by the state

$$\varphi := \langle \Phi_0, (\cdot) \Phi_0 \rangle$$

and $a_j^\varepsilon := a_{e_j}^\varepsilon$ ($\varepsilon \in \{+, 0, -\}$) the CAP operators associated with φ .

The n -particle vectors (in the e -basis) are then defined by

$$\Phi_{n, \bar{m}} := a_1^{+m_1} \cdots a_d^{+m_d} \Phi_0 ; \quad \bar{m} := (m_1, \cdots, m_d) \in D_n \quad (31)$$

where

$$\Phi_0 \equiv \Phi_{0, (0, \dots, 0)}$$

is the vacuum vector.

Theorem 6 Let $e \equiv \{e_j\}_{j \in D}$ be a linear basis of $V_{\mathbb{C}}$. Let φ be a product state on \mathcal{P} and let $a_{e_j}^+, a_{e_j}, a_{e_j}^0$ be its CAP operators in the e -basis. Then for any $d \in \mathbb{N}$ with $d \leq |D|$,

$$\langle a_{e_1}^{+m_1} \dots a_{e_d}^{+m_d} \Phi_0, a_{e_1}^{+n_1} \dots a_{e_d}^{+n_d} \Phi_0 \rangle = 0 \quad (32)$$

whenever $\bar{m} := (m_1, \dots, m_d) \neq \bar{n} := (n_1, \dots, n_d) \in \mathbb{N}^d$, while if $\bar{m} = \bar{n}$,

$$\|a_{e_1}^{+n_1} \dots a_{e_d}^{+n_d} \Phi_0\|^2 = \prod_{k \in D} \|a_{e_k}^{+n_k} \Phi_0\|^2 \quad (33)$$

It is natural to ask oneself if the converse of Theorem 6 is true, namely:

given a state φ on \mathcal{P}_V such that the associated semi-scalar product satisfies the orthogonality condition: $\bar{m} := (m_1, \dots, m_d) \neq \bar{n} := (n_1, \dots, n_d) \Rightarrow$

$$\langle a_{e_1}^{+m_1} \dots a_{e_d}^{+m_d} \Phi_0, a_{e_1}^{+n_1} \dots a_{e_d}^{+n_d} \Phi_0 \rangle = 0 \quad (34)$$

then is a product state.

Theorem 7 Let $(\mathcal{H}, V_{\mathbb{C}}, a^+, \Phi_0)$ be a **symmetric** interacting Fock space over $V_{\mathbb{C}}$. For any linear basis $e \equiv (e_j)_{j \in D}$ of V , $n \in \mathbb{N}$ and $\bar{m} = (m_1, \dots, m_d) \in D_n$, denote

$$\Phi_{n, \bar{m}} := a_{e_1}^{+m_1} \dots a_{e_d}^{+m_d} \Phi_0 := a_e^{\bar{m}} \Phi_0 \quad (35)$$

Then the following statements are equivalent.

(i) There exists a linear basis $e \equiv (e_j)_{j \in D}$ of V such that for all $n \in \mathbb{N}$ and $\bar{m} \neq \bar{n} \in D_n$,

$$\begin{aligned} & \langle \Phi_{n, \bar{m}}, \Phi_{n, \bar{n}} \rangle \\ &= \langle a_{e_1}^{+m_1} \dots a_{e_d}^{+m_d} \Phi_0, a_{e_1}^{+n_1} \dots a_{e_d}^{+n_d} \Phi_0 \rangle = 0 \end{aligned} \quad (36)$$

(ii) There exists a linear basis $e \equiv (e_j)_{j \in D}$ of V such that, with the convention that $0/0 := 0$, for all $j \in D$, $n \in \mathbb{N}$ and $\bar{m} \neq \bar{n} \in D_n$,

$$a_j \Phi_{n, \bar{m}} =$$

$$= \begin{cases} 0, & \text{if } m_j = 0 \\ \frac{\|\Phi_{n,\bar{m}}\|^2}{\|\Phi_{n,(m_1,\dots,m_j-1,\dots,m_d)}\|^2} \Phi_{n,(m_1,\dots,m_j-1,\dots,m_d)}, & \text{if } m_j > 0 \end{cases} \quad (37)$$

In particular, with the notation

$$\begin{aligned} (\bar{m})_{j,\pm} &= (m_1, \dots, m_j, \dots, m_d)_{j,\pm} \\ &:= (m_1, \dots, m_j \pm 1, \dots, m_d) \in \mathbb{N}^d \end{aligned} \quad (38)$$

(36) becomes

$$a_j a^{+\bar{m}} \Phi_0 = \begin{cases} 0, & \text{if } m_j = 0 \\ \frac{\|a^{+\bar{m}} \Phi_0\|^2}{\|a^{+(\bar{m})_{j,-}} \Phi_0\|^2} a^{+(\bar{m})_{j,-}} \Phi_0, & \text{if } m_j > 0 \end{cases} \quad (39)$$

Corollary 2 In the notations and assumptions of Theorem 7, the orthogonality condition (36) is equivalent to the following condition.

(iii) Introducing for $\bar{m} = (m_1, \dots, m_j, \dots, m_d) \in D_n$ and $j, k \in D$, the notation,

$$\pi^{j,k}(\bar{m}) := ((\bar{m})_{k,+})_{j,-} = (m_1, \dots, m_j - 1, \dots, m_k + 1, \dots) \quad (40)$$

where, if $j > k$, the order of $m_j - 1$ and $m_k + 1$ in (40) is inverted, one has

$$a_j a_k^+ \Phi_{n,\bar{m}} = \begin{cases} 0, & \text{if } j \neq k \text{ and } m_j = 0 \\ \frac{\|\Phi_{n+1,(\bar{m})_{k,+}}\|^2}{\|\Phi_{n,\pi^{j,k}(\bar{m})}\|^2} \Phi_{n,\pi^{j,k}(\bar{m})}, & \\ \text{if either } j = k \text{ or } m_j > 0 \end{cases} \quad (41)$$

A characterization of product states

Theorem 8 In the above notations, suppose that φ is a state on \mathcal{P}_V . Then φ is a product state on \mathcal{P} if and only if either condition (i) or (ii) of Theorem 7 or condition (iii) of Corollary 2 holds and for all $j \in D$, $n \in \mathbb{N}$ and $\bar{m} \neq \bar{n} \in D_n$, one has

$$\begin{aligned} & \frac{\|\Phi_{n+1,(m_1,\dots,m_j,\dots,m_k+1,\dots,m_d)}\|^2}{\|\Phi_{n,(m_1,\dots,m_j-1,\dots,m_k+1,\dots,m_d)}\|^2} \\ &= \frac{\|\Phi_{n,(m_1,\dots,m_j,\dots,m_k,\dots,m_d)}\|^2}{\|\Phi_{n-1,(m_1,\dots,m_j-1,\dots,m_k,\dots,m_d)}\|^2} \end{aligned} \quad (42)$$

Remark. Denoting, for each $j \in D$, $\{\omega_{j,k}\}_{k \in \mathbb{N}}$ the Jacobi sequence of the restriction of the state φ on the \ast -algebra $\mathcal{P}(a_{e_j}^\pm)$ generated by the CA operators $a_{e_j}^\pm$, **if φ is a product state** then one has the factorization condition:

$$\begin{aligned} \|\Phi_{n,\bar{m}}\|^2 &= \|\Phi_{(m_1,0,\dots,0)}\|^2 \cdots \|\Phi_{(0,\dots,0,m_d)}\|^2 \\ &= \prod_{j \in D} \omega_{j,m_j}! \end{aligned} \quad (43)$$

which implies that (42) is satisfied.

Proposition 1 Condition (42) is equivalent to the existence of a family $\{\omega_k^{(j)}\}_{k \in \mathbb{N}}$ ($j \in D$) of sequences of strictly positive numbers satisfying (43).