

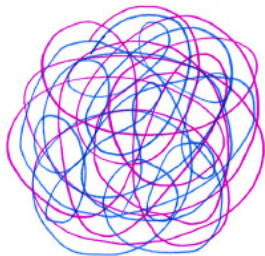
On integrability of geodesic flows on 3-dimensional manifolds

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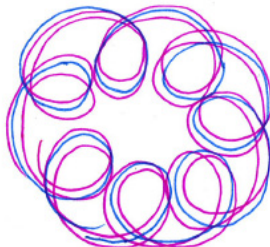
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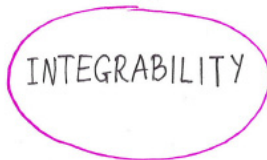
Chaos and Integrability

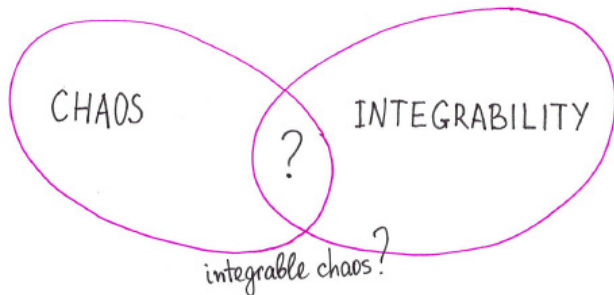


chaotic



regular





Let M be a manifold with a Riemannian metric $g = \sum g_{ij}(x) dx^i dx^j$.

Definition

Geodesic line on (M, g) is a trajectory of a point moving on M freely, i.e., by inertia:

$$\nabla_{\frac{d\gamma}{dt}} \frac{d\gamma}{dt} = 0. \quad (1)$$

Hamiltonian form of the equation of geodesics.

In canonical coordinates (x, p) on T^*M , the equations of geodesics (1) can be rewritten in Hamiltonian form:

$$\frac{dp_i}{dt} = \frac{\partial H}{\partial x^i}, \quad \frac{dx^i}{dt} = -\frac{\partial H}{\partial p_i}, \quad (2)$$

where $H = \frac{1}{2} \sum g^{ij}(x) p_i p_j$.

Properties of geodesics:

- ▶ if the metric is Euclidean, then the geodesics are straight lines.
- ▶ existence and uniqueness theorem;
- ▶ geodesic completeness and Hopf-Rinow theorem;
- ▶ existence of closed geodesics in each homotopy class;
- ▶ kinetic energy $H = \sum g_{ij}(x) \dot{x}^i \dot{x}^j$ as a first integral.

Integrability and Liouville-Arnold Theorem

Consider the geodesic flow on a Riemannian manifold (M, g) .

Complete integrability: There exist $F_1 = H, F_2, \dots, F_n$ which

- ▶ are first **integrals** (i.e., preserved by the flow),
- ▶ pairwise **commute** $\{F_i, F_j\} = 0$,
- ▶ are functionally **independent almost everywhere**.

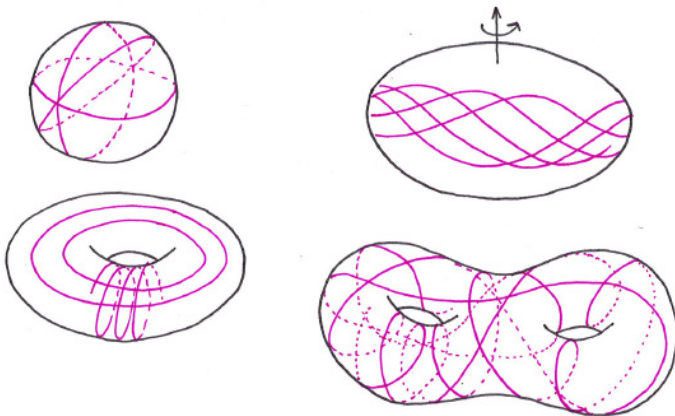
Theorem (Liouville-Arnold)

Let $X = \{F_1 = c_1, \dots, F_n = c_n\}$ be a regular, compact and connected integral surface. Then X is an n -dimensional torus and the dynamics on this torus is quasi-periodic.

Superintegrability: in addition to n commuting independent integrals, there are some more first integrals so that invariant tori have dimension **less than n** .

General problem: existence, construction and topological obstructions for integrable geodesic flows.

Two-dimensional case



Theorem (Kozlov, 1979, real analytic case)

On a 2-dimensional surface of genus $g \geq 2$, there exist no integrable geodesic flows.

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$$S^3, \quad T^3, \quad \mathbb{H}^3/\Gamma, \quad S^2 \times S^1, \quad M_g^2 \times \mathbb{R}, \quad Nil/\Gamma, \quad Sol/\Gamma, \quad \widetilde{SL(2, \mathbb{R})}/\Gamma$$

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$$5 \qquad 4 \qquad 3 \qquad ? \qquad ? \qquad ? \qquad 1 \qquad none$$

Observation. Let G be a three-dimensional Lie group, then any left-invariant Hamiltonian system on T^*G is Liouville integrable (with real analytic first integrals polynomials in momenta). Moreover, this system is superintegrable (i.e., invariant tori are two-dimensional).

The integrals are H, f_1, f_2, f_3 where f_i are linear integrals that correspond to basis right-invariant vector fields.

More generally, similar property holds true for every Lie group with two-dimensional coadjoint orbits.

Integrability of left-invariant geodesic flows on 3-dimensional Lie groups

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Question. What happens to the integrals after taking quotient $G \mapsto G/\Gamma$?

Answer. Instead of f_1, f_2, f_3 we should consider $h(f_1, f_2, f_3)$, where h is an invariant of the (coadjoint) action of Γ on \mathfrak{g}^* .

Theorem (Dinaburg, 1974)

If $\pi_1(M^n)$ is of exponential growth then $h_{\text{top}} > 0$ for any geodesic flow.

Theorem (Taimanov, 1987, geometrically simple case)

If M^n admits an integrable geodesic flow, then $\pi_1(M^n)$ is almost abelian.

Moreover, $\dim H_1(M^n, \mathbb{R}) < n$.

In particular, this holds true in the real analytic case.

Nil case

$$M_{Nil}^3 = G/\Gamma, \text{ where } G = \left\{ \begin{pmatrix} 1 & z & x \\ 0 & 1 & y \\ 0 & 0 & 1 \end{pmatrix} \right\}, \Gamma = \left\{ \begin{pmatrix} 1 & k & m \\ 0 & 1 & n \\ 0 & 0 & 1 \end{pmatrix}, k, n, m \in \mathbb{Z} \right\}$$

- ▶ integrable: $F_1 = p_x, F_2 = \exp\left(-\frac{1}{p_x^2}\right) \sin \frac{p_y}{p_x},$
- ▶ not real analytic,
- ▶ $\pi_1(M_{Nil}^3)$ is not almost commutative,
- ▶ $h_{top} = 0.$

Sol case

$$M_{Sol}^3 = G/\Gamma, \text{ where } G = \left\{ \begin{pmatrix} A^z & x \\ 0 & 0 & 1 \end{pmatrix} \right\}, \Gamma = \left\{ \begin{pmatrix} A^k & n \\ 0 & 0 & 1 \end{pmatrix}, k, n, m \in \mathbb{Z} \right\},$$

where $A^z = \exp(zB), \exp B = A = \begin{pmatrix} 2 & 1 \\ 1 & 1 \end{pmatrix}.$

- ▶ integrable: $F_1 = p_x^2 - p_x p_y - p_y^2, F_2 = \exp\left(-\frac{1}{F_1^2}\right) \cdot \sin\left(2\pi \frac{\log|p_x - \frac{1+\sqrt{5}}{2} p_y|}{\log \lambda}\right).$
- ▶ not real analytic,
- ▶ $\pi_1(M_{Sol}^3)$ is not almost commutative and has exponential growth,
- ▶ $h_{top} > 0.$

Integrability and chaos in $SL(2, \mathbb{R})$ -geometry

Consider the group $SL(2, \mathbb{R}) = \left\{ X = \begin{pmatrix} a & b \\ c & d \end{pmatrix}, \det X = 1 \right\}$ and its Lie algebra $sl(2, \mathbb{R}) = \left\{ X = \begin{pmatrix} \alpha & \beta \\ \gamma & -\alpha \end{pmatrix} = \alpha E_1 + \beta E_2 + \gamma E_3 \right\}$.

Denote by ξ_1, ξ_2, ξ_3 the **left-invariant vector fields** corresponding to the basis E_1, E_2, E_3 .

Similarly, let η_1, η_2, η_3 be the corresponding **right-invariant vector fields** on $SL(2, \mathbb{R})$.

The left-invariant **Hamiltonian**: $H = (\xi_1^2 + \xi_2\xi_3) + 2(\xi_2 - \xi_3)^2$.

Before taking quotient w.r.t. Γ , we have 3 first integrals η_1, η_2, η_3 .

Natural identification $PSL(2, \mathbb{R})$ with the unit tangent bundle $S\mathbb{H}^2$ gives the following formula in local coordinates x, y, ϕ :

$$ds^2 = \frac{dx^2 + dy^2}{y^2} + \left(d\phi + \frac{dx}{y} \right)^2$$

Let $\Gamma \subset PSL(2, \mathbb{R})$ is a Fuchsian group that acts on \mathbb{H}^2 freely and with compact quotient M_g^2 and consider M_{sl}^3 to be $PSL(2, \mathbb{R})/\Gamma$. Topologically, $M_{sl}^3 = SM_g^2$.

Apart from the Hamiltonian H the geodesic flow on M_{sl}^3 possesses one more quadratic integral (Casimir function):

$$\Delta = \eta_1^2 + \eta_2\eta_3 = \xi^2 + \xi_2\xi_3.$$

But we need one more!

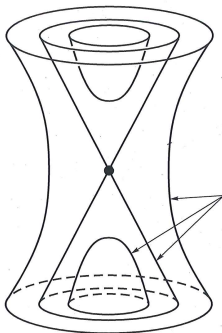
Theorem

Geodesic flow on M_{sl}^3 is *Liouville integrable* in analytic sense in the open (invariant) region of the phase space $T^*M_{sl}^3$ defined by $\Delta < 0$.

In the region with $\Delta > 0$ there are *no smooth integrals* independent from H and Δ . At the fixed integral level of H and $\Delta > 0$, the system has positive topological entropy $h_{top} \geq \sqrt{1 - C}$, where $C = \frac{H-2\Delta}{H+2\Delta}$.

Explanation of this phenomenon

In $\mathbb{R}^3(\eta_1, \eta_2, \eta_3) = T^*SL(2, \mathbb{R}) \simeq \mathfrak{sl}(2, \mathbb{R})^*$, consider the level surfaces of the integral $\Delta = \eta_1^2 + \eta_2\eta_3$.



The Fuchsian group Γ naturally acts on this space (by means of the (co)adjoint representation).

What about the orbit space and invariant functions of this action?

Let us relax the geometric simplicity condition (real analyticity) condition, i.e., assume that the first integrals of a geodesic flow on M^3 are such that the complement to the regular set filled by 3D Liouville tori is a timely-embedded polyhedron in T^*M^3 .

Theorem (L. Butler, 2005)

If an integrable geodesic flow on a closed manifold M^3 satisfies the above condition, then $\pi_1(M^3)$ is almost polycyclic.

“Conversely”, if $\pi_1(M^3)$ is polycyclic, then M^3 admits a real-analytic metric with integrable geodesic flow.

Thank you for your attention