

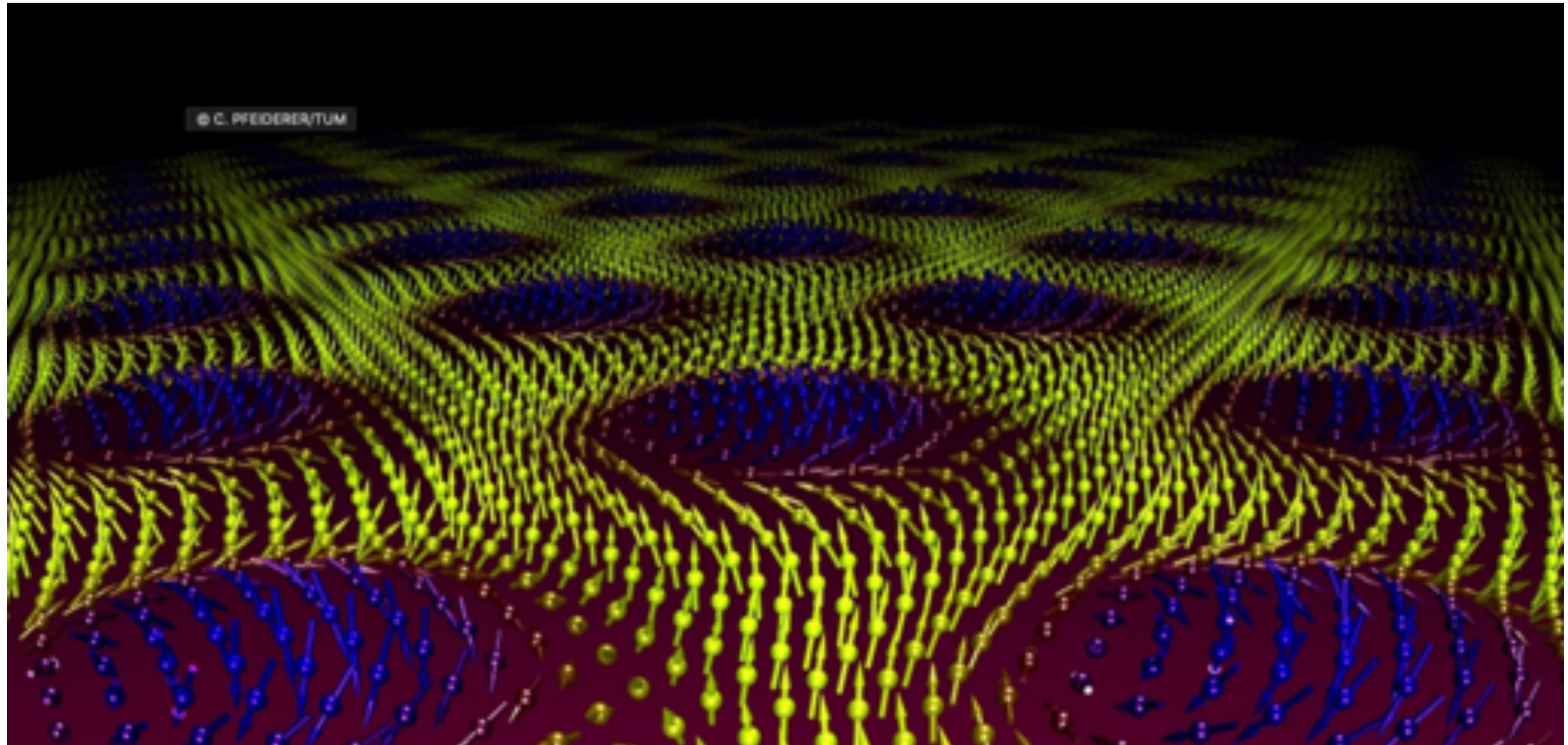
Hydrodynamics and holography of the spin current

Frontiers-2, 21.10.2020

Umut Gürsoy
ITP, Utrecht University

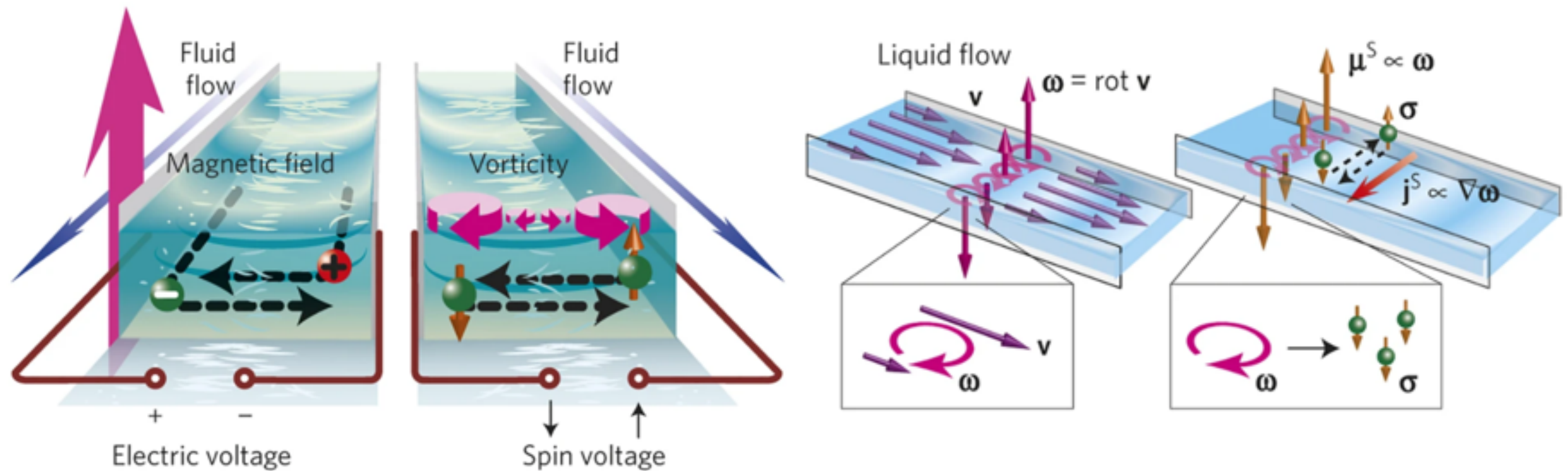
Based on: with Domingo Gallegos, arXiv:2004.05148
ongoing with Domingo Gallegos, Amos Yarom

Spin transport



- Spintronics, liquid spintronics
- Graphene, Dirac/Weyl semimetals
- Astrophysics
- Quark gluon plasma

Liquid spintronics

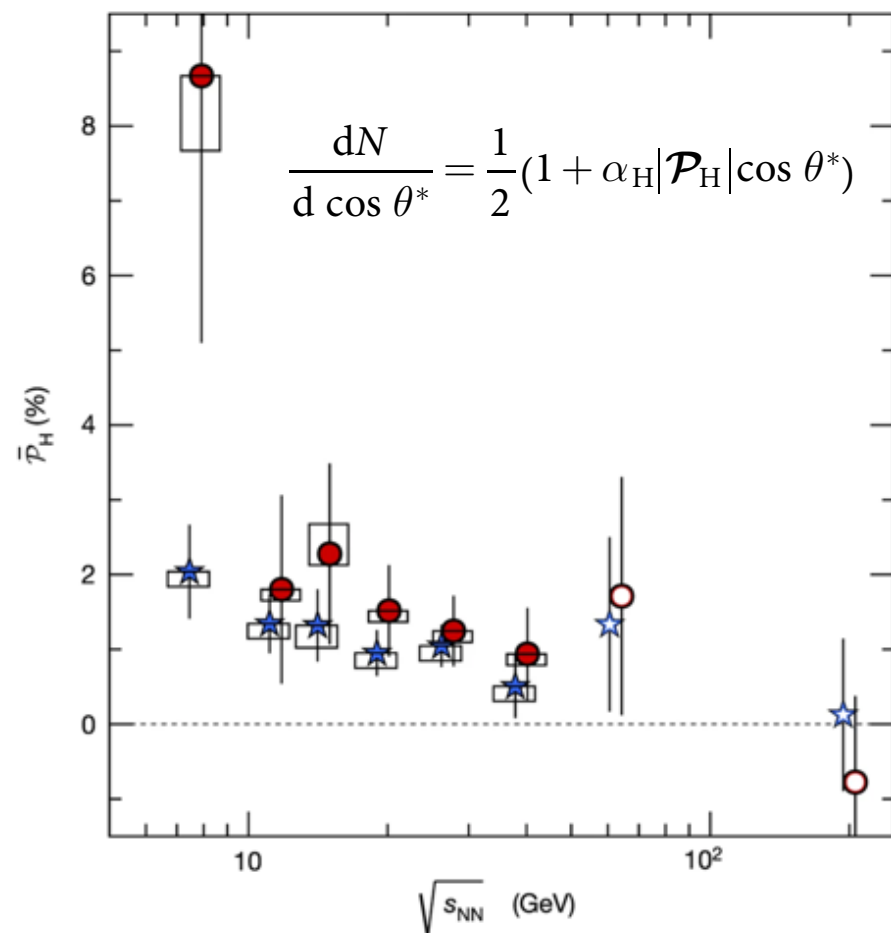
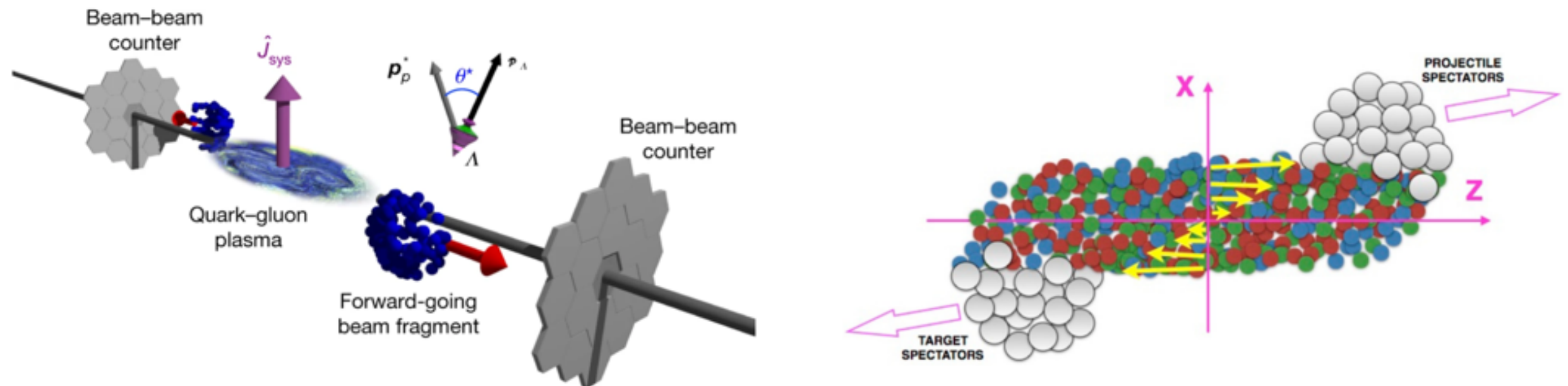


R. Takahashi et al. '16

Generation of spin density by vorticity in liquid metals (Hg, GaInSn)

$$\mu_s \propto \omega, J_s \propto \nabla \omega$$

Quark gluon plasma



Global hyperon polarization at RHIC
by spin-orbit coupling $\vec{S} \cdot \vec{J}$

QGP: most vortical fluid: $\omega \sim 10^{22} \text{ s}^{-1}$

F. Becattini et al. '17
STAR collaboration '19

Hydrodynamics with spin current

- All possible sources of spin density/current
- Relation of spin flow to energy/charge/chirality flow
- Classification of spin transport coefficients

Hydrodynamics with spin current

- All possible sources of spin density/current
- Relation of spin flow to energy/charge/chirality flow
- Classification of spin transport coefficients

$\Rightarrow$ construction of hydrodynamics with spin

$$T_{\mu\nu} \quad S_{\mu\nu}^{\lambda}$$

Outline

- Definition of spin current
- Hydrodynamics with spin current and torsion
- Holographic description: first order gravity

Lovelock-Chern-Simons theory

- Analytic blackhole solutions
- Spin transport coefficients
- Discussion

Ambiguity in the spin current

Total angular momentum

$$J^{\lambda\mu\nu} = x^\mu T^{\lambda\nu} - x^\nu T^{\lambda\mu} + S^{\lambda\mu\nu}$$

Conservation laws

$$\partial_\mu T^{\mu\nu} = 0, \quad \partial_\lambda J^{\lambda\mu\nu} = 0$$

Preserved by

$$T'^{\mu\nu} = T^{\mu\nu} + \frac{1}{2} \partial_\lambda (\Phi^{\lambda\mu\nu} - \Phi^{\mu\lambda\nu} - \Phi^{\nu\lambda\mu}),$$
$$S'^{\lambda\mu\nu} = S^{\lambda\mu\nu} - \Phi^{\lambda\mu\nu}$$

Belinfante-Rosenfeld

$$\Phi^{\lambda\mu\nu} = S^{\lambda\mu\nu}$$

Removes spin current completely, renders energy-momentum tensor symmetric, compatible with GR

Spin in effective action

Consider quantum field in a nontrivial Lorentz representation

$$e^{iW[e,\omega]} = \int D\Psi e^{iI[e,\omega,\Psi]}$$

Variations define the energy-momentum and spin current

$$T^{\mu\nu} = \frac{\delta W}{\delta e_\mu^a} e_a^\nu, \quad S_{ab}^\lambda = \frac{\delta W}{\delta \omega_\lambda^{ab}}$$

Spin in effective action

Consider quantum field in a nontrivial Lorentz representation

$$e^{iW[e,\omega]} = \int D\Psi e^{iI[e,\omega,\Psi]}$$

Variations define the energy-momentum and spin current

$$T^{\mu\nu} = \frac{\delta W}{\delta e_\mu^a} e_a^\nu, \quad S_{ab}^\lambda = \frac{\delta W}{\delta \omega_\lambda^{ab}}$$

Dependent in the absence of **torsion**:

$$T^a = de^a + \omega_b^a \wedge e^b$$

$\Rightarrow$ spin current unambiguous when sourced by torsion

Spin in effective action

Consider quantum field in a nontrivial Lorentz representation

$$e^{iW[e,\omega]} = \int D\Psi e^{iI[e,\omega,\Psi]}$$

Variations define the energy-momentum and spin current

$$T^{\mu\nu} = \frac{\delta W}{\delta e^a_\mu} e^\nu_a, \quad S^\lambda_{ab} = \frac{\delta W}{\delta \omega_\lambda^{ab}}$$

Dependent in the absence of **torsion**:

$$T^a = de^a + \omega_b^a \wedge e^b$$

$\Rightarrow$ **spin current unambiguous** when sourced by torsion

Easier to formulate in terms of **contorsion**:

$$\omega_\mu^{ab} = \mathring{\omega}_\mu^{ab} + K_\mu^{ab}, \quad \mathring{\omega} \sim \partial e$$

Lowest order source of spin flow in flat space $e = 1$

Hydrodynamic equations

Require invariance of $W[e,\omega]$ under

1. Diffeomorphisms

$$\delta_\xi e^a = \mathcal{L}_\xi e^a, \quad \delta_\xi \omega^{ab} = \mathcal{L}_\xi \omega^{ab}$$

2. Local Lorentz transformations

$$\delta_\lambda e^a = -\lambda^a_b e^b, \quad \delta_\lambda \omega^{ab} = D\lambda^{ab}$$

Hydrodynamic equations

Require invariance of $W[e,\omega]$ under

1. Diffeomorphisms

$$\delta_\xi e^a = \mathcal{L}_\xi e^a, \quad \delta_\xi \omega^{ab} = \mathcal{L}_\xi \omega^{ab}$$

2. Local Lorentz transformations

$$\delta_\lambda e^a = -\lambda^a_b e^b, \quad \delta_\lambda \omega^{ab} = D\lambda^{ab}$$

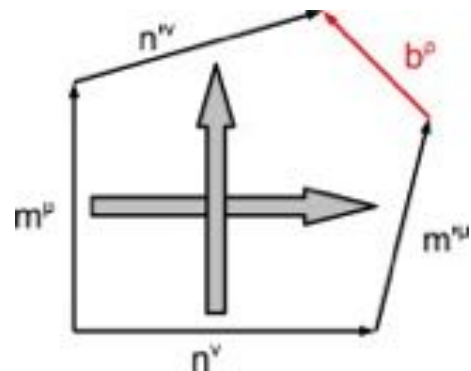
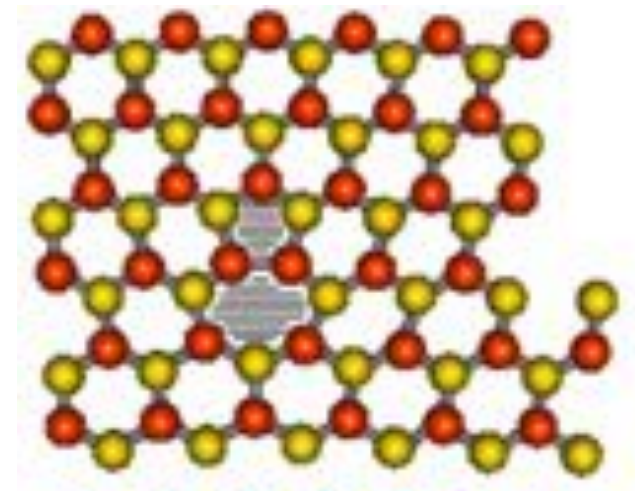
Relativistic hydrodynamics with spin current

$$\begin{aligned} \overset{\circ}{\nabla}_\mu T^{\mu\nu} &= R^{\rho\sigma\nu\lambda} S_{\lambda\rho\sigma} - T_{[\rho\sigma]} K^{\rho\sigma\nu}, & \text{analogous to EM work } \mathbf{j} \cdot \mathbf{E} \\ \overset{\circ}{\nabla}_\lambda S^\lambda_{\mu\nu} &= T_{[\mu\nu]} - 2S^\lambda_{\rho[\mu} K^\rho_{\nu]\lambda} & \text{antisymm. EMT generates S} \end{aligned}$$

Torsion in condensed matter

Kondo '52; Katanaev, Volovich '92; Furtado, Moraes '99; Mesaros, Sadri, Zaanen '09;
de Juan, Cortijo, Vozmediano '10

e.g. dislocations in graphene

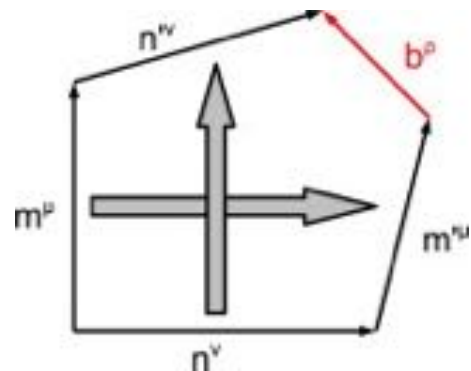
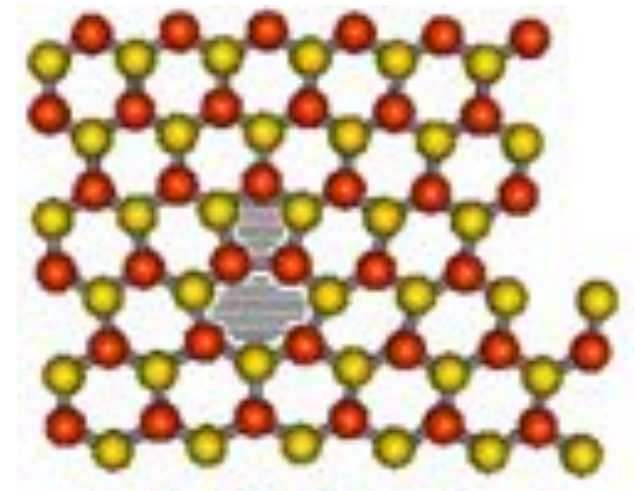


$$m'^{\lambda} - n'^{\lambda} = [\Gamma_{\mu\nu}^{\lambda} - \Gamma_{\nu\mu}^{\lambda}] m^{\mu} n^{\nu} = T_{\nu\mu}^{\lambda} m^{\mu} n^{\nu}$$

Torsion in condensed matter

Kondo '52; Katanaev, Volovich '92; Furtado, Moraes '99; Mesaros, Sadri, Zaanen '09;
de Juan, Cortijo, Vozmediano '10

e.g. dislocations in graphene



$$m'^{\lambda} - n'^{\lambda} = [\Gamma_{\mu\nu}^{\lambda} - \Gamma_{\nu\mu}^{\lambda}] m^{\mu} n^{\nu} = T_{\nu\mu}^{\lambda} m^{\mu} n^{\nu}$$

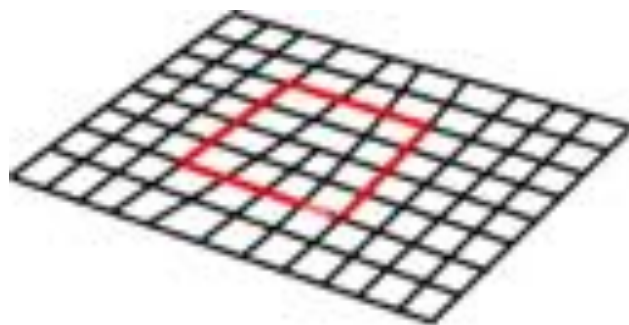
Torsion = continuum description of dislocations in the background

Classification of torsion sources

Assume flat 3+1D metric with nontrivial torsion $T_{\mu\nu}^a = (\omega_b^a)_\mu \wedge \delta_\nu^b$

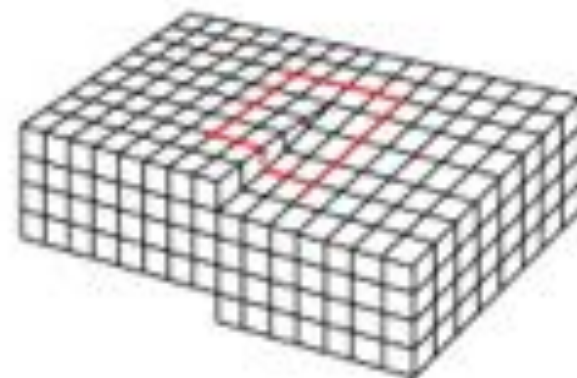
In 3D: ω_i^{ab} $\begin{matrix} \rightarrow \text{loop plane} \\ \rightarrow \text{Burger's vector} \end{matrix}$ $\Rightarrow$ 9 degrees of freedom

$i \parallel (ab)$



edge dislocation (6 d.o.f)

$i \perp (ab)$



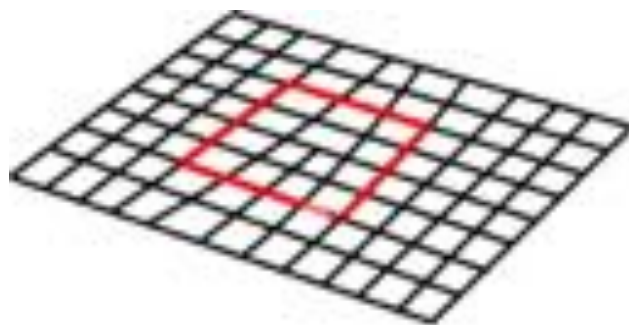
screw dislocation (3 d.o.f)

Classification of torsion sources

Assume flat 3+1D metric with nontrivial torsion $T_{\mu\nu}^a = (\omega_b^a)_\mu \wedge \delta_\nu^b$

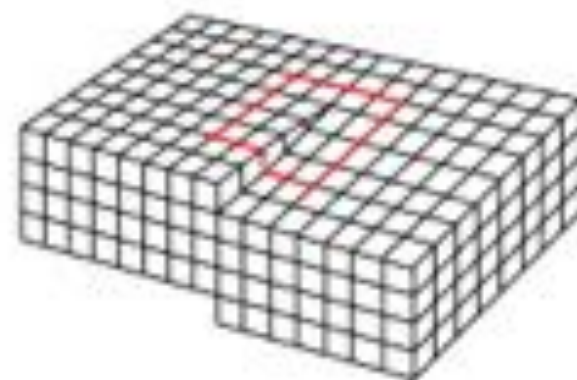
In 3D: ω_i^{ab} $\begin{matrix} \rightarrow \text{loop plane} \\ \rightarrow \text{Burger's vector} \end{matrix} \implies 9 \text{ degrees of freedom}$

$i \parallel (ab)$



edge dislocation (6 d.o.f)

$i \perp (ab)$



screw dislocation (3 d.o.f)

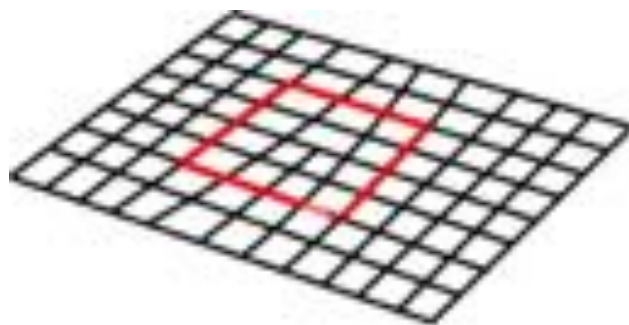
In 4D: $\omega_\mu^{ab} \implies \text{Wilson-edge (6 d.o.f.)} + \text{Wilson-screw (6 d.o.f.)} \\ + \text{Polyakov-edge (6 d.o.f.)} + \text{Polyakov-screw (6 d.o.f.)}$

Classification of torsion sources

Assume flat 3+1D metric with nontrivial torsion $T_{\mu\nu}^a = (\omega_b^a)_\mu \wedge \delta_\nu^b$

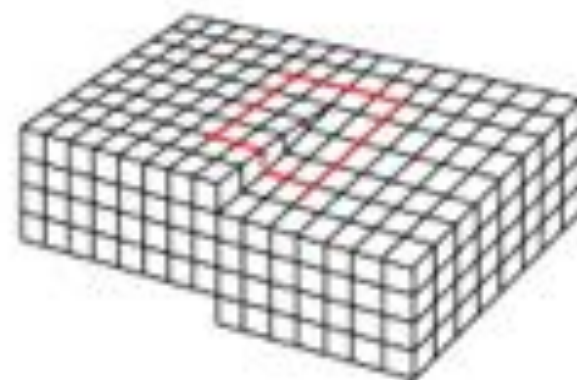
In 3D: ω_i^{ab} $\begin{matrix} \longrightarrow & \text{loop plane} \\ \longrightarrow & \text{Burger's vector} \end{matrix} \implies 9 \text{ degrees of freedom}$

$i \parallel (ab)$



edge dislocation (6 d.o.f)

$i \perp (ab)$



screw dislocation (3 d.o.f)

In 4D: $\omega_\mu^{ab} \implies \text{Wilson-edge (6 d.o.f.)} + \text{Wilson-screw (6 d.o.f.)} \\ + \text{Polyakov-edge (6 d.o.f.)} + \text{Polyakov-screw (6 d.o.f.)}$

Not a Lorentz invariant decomposition

Classification of torsion sources

Lorentz invariant decomposition: $\omega^{ab} = \omega_{\text{vector}}^{ab} + \omega_{\text{axial}}^{ab} + \omega_{\text{tensor}}^{ab}$

$$(\omega_{\text{vector}}^{ab})_{\mu} = \delta_{\mu}^a \tilde{V}^b - \delta_{\mu}^b \tilde{V}^a \quad (\omega_{\text{axial}}^{ab})_{\mu} = \epsilon^{abc}{}_d \tilde{A}_c \delta_{\mu}^d$$

In hydro with flow velocity u^{μ} further decompose: Gallegos, UG , '20

$$\omega_{\beta}^{ab} = e_{\mu}^a e_{\nu}^b \left[-\epsilon^{\mu\nu}{}_{\rho\sigma} u^{\rho} (\mu_A \delta_{\beta}^{\sigma} + \mathcal{C}_{\beta}^{\sigma} - \mathcal{A}_1^{\sigma} u_{\beta}) + 2 (\mu_V u^{[\mu} - \mathcal{V}_2^{[\mu} \Delta_{\beta}^{\nu]}) \right. \\ \left. + 2 u^{[\mu} \epsilon^{\nu]}{}_{\rho\sigma\beta} u^{\rho} \mathcal{A}_2^{\sigma} - 2 u^{[\mu} (\mathcal{V}_1^{\nu]} u_{\beta} + \mathcal{H}_{\beta}^{\nu]} \right]$$

1 scalar, 1 pseudo-scalar, 2 transverse vectors, 2 axial transverse vectors,
2 traceless-transverse symmetric tensors

Constitutive relations

Energy-momentum tensor:

$$T^{\mu\nu} = \varepsilon u^\mu u^\nu - p \Delta^{\mu\nu} + \underbrace{\bar{q}^\mu u^\nu + u^\mu q^\nu}_{\text{heat currents}} + \underbrace{\pi^{\mu\nu}}_{\text{symmetric} \ni \text{shear}} + \underbrace{\tau^{\mu\nu}}_{\text{intrinsic torque}}$$

Spin current:

$$\begin{aligned} S^{\lambda\mu\nu} = & u^\lambda \left(u^\nu n_V^\mu - u^\mu n_V^\nu + \epsilon^{\mu\nu}{}_{\rho\sigma} u^\rho n_A^\sigma \right) + \rho_A \epsilon^{\lambda\mu\nu\rho} u_\rho + \rho_V \left(\Delta^{\lambda\nu} u^\mu - \Delta^{\lambda\mu} u^\nu \right) \\ & + N^{\lambda\mu} u^\nu - N^{\lambda\nu} u^\mu - \epsilon^{\mu\nu\rho\sigma} u_\rho \bar{N}_\sigma^\lambda + \Delta^{\lambda\nu} \bar{n}_V^\mu - \Delta^{\lambda\mu} \bar{n}_V^\nu \\ & + \epsilon^{\lambda\mu}{}_{\rho\sigma} u^\nu u^\rho \bar{n}_A^\sigma - \epsilon^{\lambda\nu}{}_{\rho\sigma} u^\mu u^\rho \bar{n}_A^\sigma \end{aligned}$$

Constitutive relations

Energy-momentum tensor:

$$T^{\mu\nu} = \varepsilon u^\mu u^\nu - p \Delta^{\mu\nu} + \underbrace{\bar{q}^\mu u^\nu + u^\mu q^\nu}_{\text{heat currents}} + \underbrace{\pi^{\mu\nu}}_{\text{symmetric} \ni \text{shear}} + \underbrace{\tau^{\mu\nu}}_{\text{intrinsic torque}}$$

Spin current:

$$\begin{aligned} S^{\lambda\mu\nu} = & u^\lambda (u^\nu n_V^\mu - u^\mu n_V^\nu + \epsilon^{\mu\nu}{}_{\rho\sigma} u^\rho n_A^\sigma) + \rho_A \epsilon^{\lambda\mu\nu\rho} u_\rho + \rho_V (\Delta^{\lambda\nu} u^\mu - \Delta^{\lambda\mu} u^\nu) \\ & + N^{\lambda\mu} u^\nu - N^{\lambda\nu} u^\mu - \epsilon^{\mu\nu\rho\sigma} u_\rho \bar{N}_\sigma^\lambda + \Delta^{\lambda\nu} \bar{n}_V^\mu - \Delta^{\lambda\mu} \bar{n}_V^\nu \\ & + \epsilon^{\lambda\mu}{}_{\rho\sigma} u^\nu u^\rho \bar{n}_A^\sigma - \epsilon^{\lambda\nu}{}_{\rho\sigma} u^\mu u^\rho \bar{n}_A^\sigma = u^\lambda \rho^{\mu\nu} + S_\perp^{\lambda\mu\nu} \end{aligned}$$

Constitutive relations

Energy-momentum tensor:

$$T^{\mu\nu} = \varepsilon u^\mu u^\nu - p \Delta^{\mu\nu} + \underbrace{\bar{q}^\mu u^\nu + u^\mu q^\nu}_{\text{heat currents}} + \underbrace{\pi^{\mu\nu}}_{\text{symmetric} \ni \text{shear}} + \underbrace{\tau^{\mu\nu}}_{\text{intrinsic torque}}$$

Spin current:

$$\begin{aligned} S^{\lambda\mu\nu} = & u^\lambda (u^\nu n_V^\mu - u^\mu n_V^\nu + \epsilon^{\mu\nu}{}_{\rho\sigma} u^\rho n_A^\sigma) + \rho_A \epsilon^{\lambda\mu\nu\rho} u_\rho + \rho_V (\Delta^{\lambda\nu} u^\mu - \Delta^{\lambda\mu} u^\nu) \\ & + N^{\lambda\mu} u^\nu - N^{\lambda\nu} u^\mu - \epsilon^{\mu\nu\rho\sigma} u_\rho \bar{N}_\sigma^\lambda + \Delta^{\lambda\nu} \bar{n}_V^\mu - \Delta^{\lambda\mu} \bar{n}_V^\nu \\ & + \epsilon^{\lambda\mu}{}_{\rho\sigma} u^\nu u^\rho \bar{n}_A^\sigma - \epsilon^{\lambda\nu}{}_{\rho\sigma} u^\mu u^\rho \bar{n}_A^\sigma = u^\lambda \rho^{\mu\nu} + S_\perp^{\lambda\mu\nu} \end{aligned}$$

Spin Source	μ_V	μ_A	$\mathcal{V}_1^\mu$	$\mathcal{A}_1^\mu$	$\mathcal{V}_2^\mu$	$\mathcal{A}_2^\mu$	$\mathcal{H}^{\mu\nu}$	$\mathcal{C}^{\mu\nu}$
Spin current	ρ_V	ρ_A	n_V^μ	n_A^μ	$\bar{n}_V^\mu$	$\bar{n}_A^\mu$	$N^{\mu\nu}$	$\bar{N}^{\mu\nu}$
Degrees of Freedom	1	1	3	3	3	3	5	5

Spin transport from effective action

Define **spin chemical potential**:

$$\mu^{ab} = u^\mu \omega_\mu^{ab} = u^a \mathcal{V}_1^b - u^b \mathcal{V}_1^a - \epsilon^{abcd} u_c \mathcal{A}_{1d}$$

Spin transport from effective action

Define **spin chemical potential**:

$$\mu^{ab} = u^\mu \omega_\mu^{ab} = u^a \mathcal{V}_1^b - u^b \mathcal{V}_1^a - \epsilon^{abcd} u_c \mathcal{A}_{1d}$$

Hydrostatics from effective action:

$$W = \int d^4x \sqrt{-g} P[T, \mu^i] + \sum_i \sigma_i(T) \mathcal{S}^i + \dots$$

$$\mu^i = \{\mathcal{V}_1 \cdot \mathcal{V}_1, \mathcal{A}_1 \cdot \mathcal{A}_1, \mathcal{V}_1 \cdot \mathcal{A}_1\} \quad S^i = \{\mathcal{V}_1 \cdot \omega, \mathcal{V}_1 \cdot a, \mathcal{A}_1 \cdot \omega, \mathcal{A}_1 \cdot a\}$$

Spin transport from effective action

Define **spin chemical potential**:

$$\mu^{ab} = u^\mu \omega_\mu^{ab} = u^a \mathcal{V}_1^b - u^b \mathcal{V}_1^a - \epsilon^{abcd} u_c \mathcal{A}_{1d}$$

Hydrostatics from effective action:

$$W = \int d^4x \sqrt{-g} P[T, \mu^i] + \sum_i \sigma_i(T) \mathcal{S}^i + \dots$$

$$\mu^i = \{\mathcal{V}_1 \cdot \mathcal{V}_1, \mathcal{A}_1 \cdot \mathcal{A}_1, \mathcal{V}_1 \cdot \mathcal{A}_1\} \quad \mathcal{S}^i = \{\mathcal{V}_1 \cdot \omega, \mathcal{V}_1 \cdot a, \mathcal{A}_1 \cdot \omega, \mathcal{A}_1 \cdot a\}$$

Polarization by **spin chemical potential, vorticity and acceleration**

$$\begin{aligned} S_{ab}^\lambda = & u^\lambda \left\{ q_1 \mathcal{V}_{1[a} u_{b]} + q_2 \epsilon_{abcd} u^c \mathcal{A}_1^d + q^3 \tilde{\mu}_{ab} \right\} \\ & + u^\lambda \left\{ \sigma_1 \omega_{[a} u_{b]} + \sigma_2 a_{[a} u_{b]} + \sigma_3 \Omega_{ab} + \sigma_4 \epsilon_{abcd} a^c u^d \right\} \end{aligned}$$

Bjorken flow

Boost invariant expansion in z-direction: $u^\tau = 1$

No spin chemical potential and vorticity

Bjorken flow

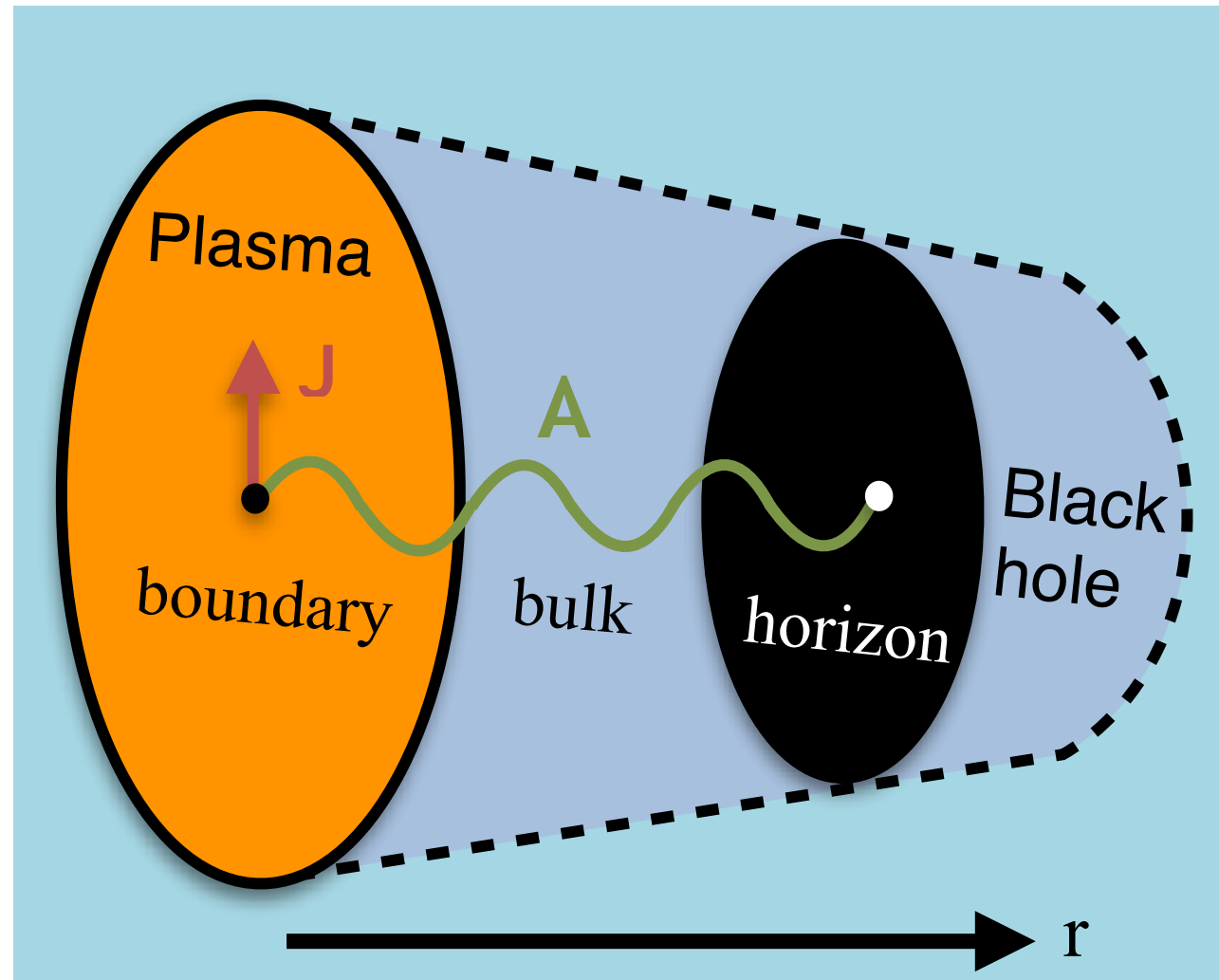
Boost invariant expansion in z-direction: $u^\tau = 1$

No spin chemical potential and vorticity

Acceleration polarizes the fluid in the z-direction (in the lab frame):

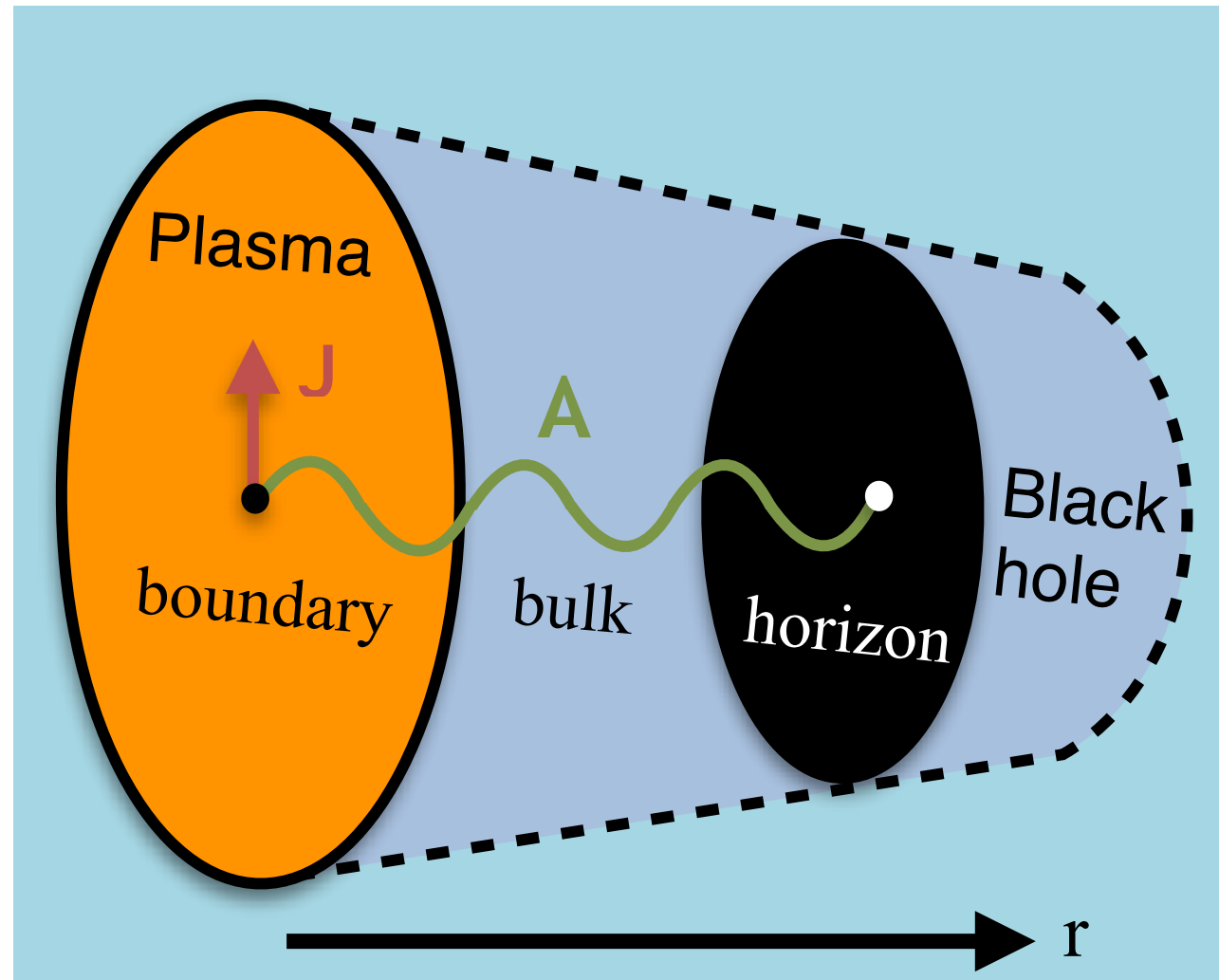
$$S_{xy}^\tau = \sigma_4(T) \frac{\sinh 2y}{2\tau}$$

Holography



- Current J is mapped onto fluctuations on the horizon
- Generic properties of near horizon geometry
 $\Rightarrow$ universal transport e.g. $\eta/s = 1/4\pi$ Policastro, Son, Starinets '01
Kovtun, Son, Starinets '04

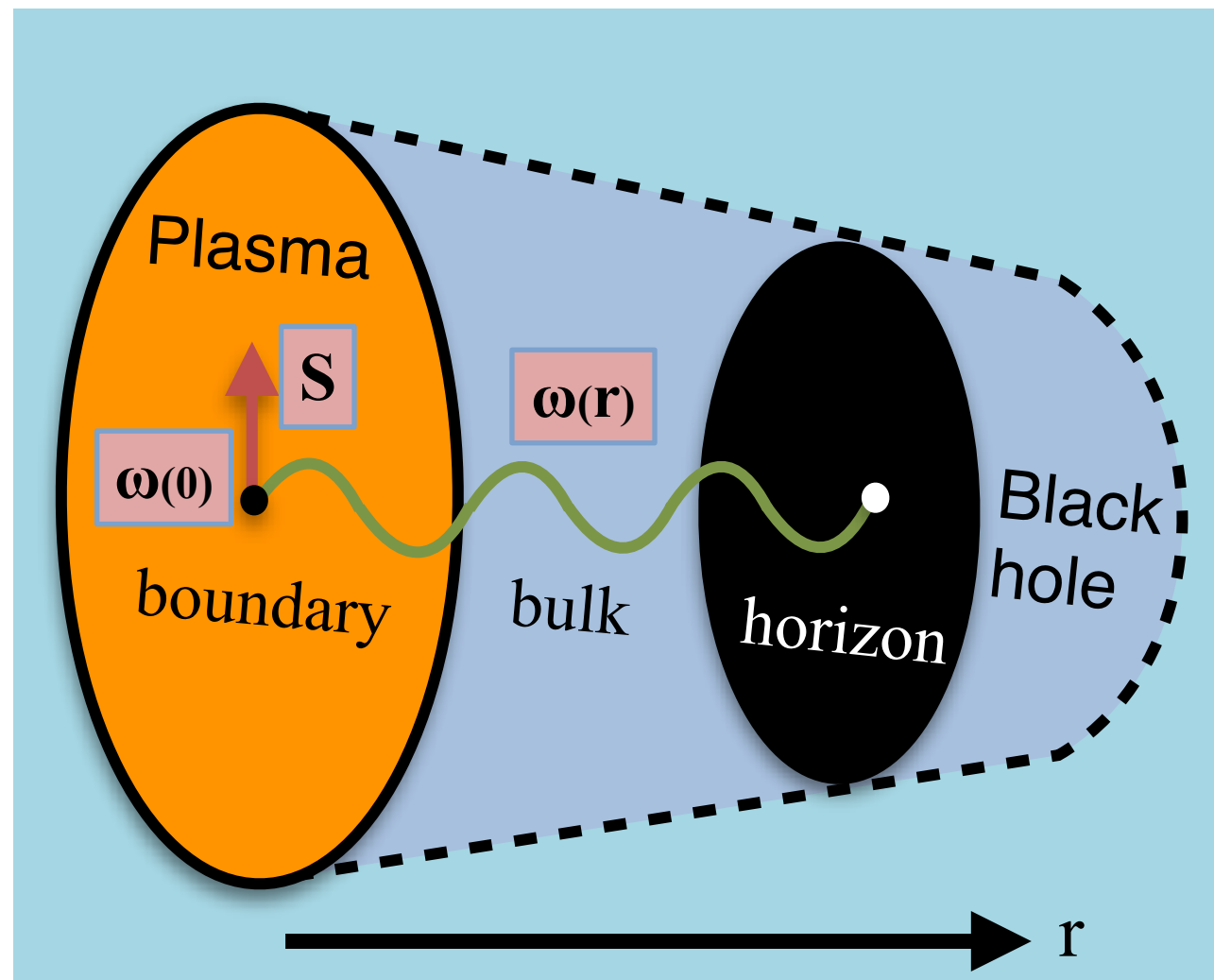
Holography



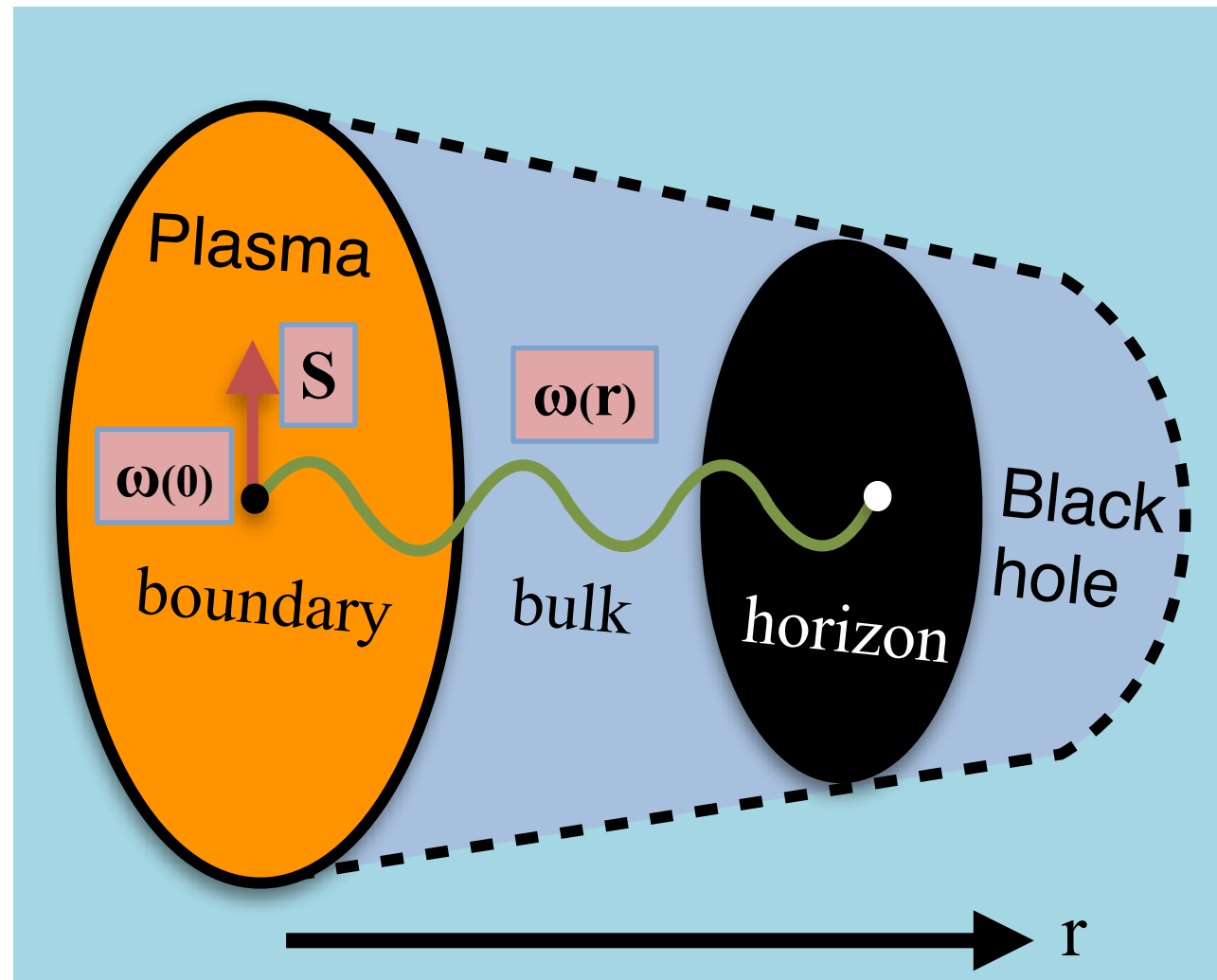
- Current J is mapped onto fluctuations on the horizon
- Generic properties of near horizon geometry
 $\Rightarrow$ universal transport e.g. $\eta/s = 1/4\pi$

Similar universality in spin transport?

Holography of spin flow



Holography of spin flow



- Need a first order formulation with independent e^a and ω^{ab}
- Practical to have analytic black-hole solutions

Holographic description

Toy model: Lovelock-Chern-Simons AdS-gravity in 5D

Chamseddine '89; Banados, Garay, Henneaux '96; Zanelli 05;
Banados, Miskovic, Theisen '06; Cvetkovic, Miskovic, Simic '17

$$S = \int \text{Tr} \left\{ \mathcal{F} \wedge \mathcal{F} \wedge \mathcal{A} - \frac{1}{2} \mathcal{F} \wedge \mathcal{A} \wedge \mathcal{A} \wedge \mathcal{A} + \frac{1}{10} \mathcal{A} \wedge \mathcal{A} \wedge \mathcal{A} \wedge \mathcal{A} \wedge \mathcal{A} \right\}$$

based on the $G = \text{SO}(4,2)$ group algebra

- CS gravity for non-Abelian G , $D \geq 5$ has local degrees of freedom unlike 3D
- Boundary symmetry $\simeq$ WZW₄ conformal theory with G

Nair, Schiff '90; Losev, Nekrasov, Moore, Shatashvili '95

Lovelock-Chern-Simons

Relation to Einstein-Cartan in 5D:

$$\mathcal{J}_{A6} = P_A, \mathcal{J}_{AB} = J_{AB}$$

$$\mathcal{A} \equiv \hat{e}^A P_A + \frac{1}{2} \hat{\omega}^{AB} J_{AB},$$

$$\mathcal{F} = \hat{T}^A P_A + \frac{1}{2} (R^{AB} + \hat{e}^A \hat{e}^B) J_{AB}$$

First order gravity in 5D:

$$S_{LCS} = \kappa \int_{M_5} \epsilon_{ABCDE} \left[\underbrace{\hat{R}^{AB} \hat{R}^{CD} \hat{e}^E}_{\text{Gauss-Bonnet with } \lambda_{\text{GB}}=1/4} + \frac{2}{3} \underbrace{\hat{R}^{AB} \hat{e}^C \hat{e}^D \hat{e}^E}_{\text{Ricci}} + \frac{1}{5} \underbrace{\hat{e}^A \hat{e}^B \hat{e}^C \hat{e}^D \hat{e}^E}_{\text{Cosmological Constant}} \right]$$

Lovelock-Chern-Simons

Relation to Einstein-Cartan in 5D: $\mathcal{J}_{A6} = P_A, \mathcal{J}_{AB} = J_{AB}$

$$\mathcal{A} \equiv \hat{e}^A P_A + \frac{1}{2} \hat{\omega}^{AB} J_{AB},$$

$$\mathcal{F} = \hat{T}^A P_A + \frac{1}{2} (R^{AB} + \hat{e}^A \hat{e}^B) J_{AB}$$

First order gravity in 5D:

$$S_{LCS} = \kappa \int_{M_5} \epsilon_{ABCDE} \left[\underbrace{\hat{R}^{AB} \hat{R}^{CD} \hat{e}^E}_{\text{Gauss-Bonnet with } \lambda_{GB}=1/4} + \frac{2}{3} \underbrace{\hat{R}^{AB} \hat{e}^C \hat{e}^D \hat{e}^E}_{\text{Ricci}} + \frac{1}{5} \underbrace{\hat{e}^A \hat{e}^B \hat{e}^C \hat{e}^D \hat{e}^E}_{\text{Cosmological Constant}} \right]$$

$\Rightarrow$ Standard AdS solution for vanishing torsion

New solutions with propagating torsion Gallegos, UG , '20

Charges and thermodynamics

Holographic renormalization in first order gravity: Banados, Miskovic, Teisen '06

$$\beta\mathcal{F} = \lim_{\epsilon \rightarrow 0} [S_{\text{on-shell}}(\epsilon) - V(\epsilon)] \quad \text{counterterm action}$$


Entropy as a charge from the Lee-Wald formula at horizon:

$$\mathcal{S} = 4\pi \int_{M_3^h} \epsilon_{ABCD F} \left(\hat{R}^{CD} \hat{e}^F + \frac{1}{3} \hat{e}^C \hat{e}^D \hat{e}^F \right) n^{AB}, \quad n^{AB} \equiv D^{[A} \xi^{B]}$$

Total mass of the black hole: Gallegos, UG, '20

$$M = M \equiv M_0 - \mu_I Q^I = \int_{M_3} d^3x \left[n_\mu \xi_\nu T^{\mu\nu} + \frac{1}{2} n_\lambda S^\lambda_{\mu\nu} \omega^{\mu\nu}{}_\alpha \xi^\alpha \right]$$

Satisfy the Smarr relation and the first-law:

$$\mathcal{F} = M_0 - T\mathcal{S} - \mu_I Q^I, \quad d\mathcal{F} = -\mathcal{S}dT - Q^I d\mu_I,$$

Solution I: single axial

Thermodynamics:

$$\varepsilon = 96\pi^4 \kappa T^4 \left(1 + \frac{\mathcal{A}_2^2}{6\pi^2 T^2} \right),$$

$$p = -32\pi^4 \kappa T^4 \left(1 + \frac{\mathcal{A}_2^2}{6\pi^2 T^2} \right),$$

$$\pi^{\mu\nu} = 16\pi^2 \kappa T^2 \left(\Delta_{(\rho}^{\mu} \Delta_{\sigma)}^{\nu} - \frac{1}{3} \Delta^{\mu\nu} \Delta_{\rho\sigma} \right) \mathcal{A}_2^{\rho} \mathcal{A}_2^{\sigma}$$

Spin current:

$$S^{\lambda\mu\nu} = -32\pi^2 \kappa T^2 \left(u^{\lambda} \epsilon^{\mu\nu\alpha\beta} + \epsilon^{\lambda\alpha\beta[\mu} u^{\nu]} \right) u_{\alpha} \mathcal{A}_{2\beta}$$

Solution I: single axial

Thermodynamics:

$$\begin{aligned}\varepsilon &= 96\pi^4 \kappa T^4 \left(1 + \frac{\mathcal{A}_2^2}{6\pi^2 T^2} \right), \\ p &= -32\pi^4 \kappa T^4 \left(1 + \frac{\mathcal{A}_2^2}{6\pi^2 T^2} \right), \\ \pi^{\mu\nu} &= 16\pi^2 \kappa T^2 \left(\Delta_{(\rho}^{\mu} \Delta_{\sigma)}^{\nu} - \frac{1}{3} \Delta^{\mu\nu} \Delta_{\rho\sigma} \right) \mathcal{A}_2^{\rho} \mathcal{A}_2^{\sigma}\end{aligned}$$

Spin current:

$$S^{\lambda\mu\nu} = -32\pi^2 \kappa T^2 \left(u^{\lambda} \epsilon^{\mu\nu\alpha\beta} + \epsilon^{\lambda\alpha\beta[\mu} u^{\nu]} \right) u_{\alpha} \mathcal{A}_{2\beta}$$

- Conformal fluid with axial vector spin source
- Non-trivial spin current *even for symmetric stress tensor!*

Solution II: Vector-axial-tensor

Thermodynamics:

$$p = -\epsilon/3 = -32\pi^4 \kappa T^4 \left(1 + \frac{\mathcal{C}^{\alpha\beta} \mathcal{C}_{\alpha\beta} - 2\mathcal{V}_2^2 - 6\mu_A^2}{12\pi^2 T^2} \right)$$

$$q^\mu = -16\pi^2 \kappa T^2 (\mathcal{C}_\alpha^\mu \mathcal{A}_1^\alpha - 2\mu_A \mathcal{A}_1^\mu + \epsilon^{\mu\alpha\rho\sigma} u_\alpha \mathcal{A}_{1\rho} \mathcal{V}_{2\sigma}) , \bar{q}^\mu = 0$$

$$\pi^{\mu\nu} = -16\pi^2 \kappa T^2 \left(C^{\mu\alpha} C_\alpha^\nu + \mathcal{V}_2^\mu \mathcal{V}_2^\nu - \frac{1}{3} (\mathcal{C}_\beta^\alpha \mathcal{C}_\alpha^\beta + \mathcal{V}_2^2) \Delta^{\mu\nu} - \mu_A C^{\mu\nu} \right)$$

$$\tau^{\mu\nu} = 16\pi^2 \kappa T^2 u_\rho \mathcal{V}_2^\beta \epsilon^{\mu\nu\rho\sigma} (\mathcal{C}_\beta^\sigma + \mu_A \delta_\beta^\sigma)$$

Spin current:

$$S^{\lambda\mu\nu} = 32\pi^2 \kappa T^2 \left[\mu_A \epsilon^{\lambda\mu\nu\alpha} u_\alpha - u_\alpha \epsilon^{\lambda\alpha\beta[\mu} (\mathcal{C}_{\beta}^{\nu]} + u^{\nu]} \mathcal{A}_{1\beta}) + \Delta^{\lambda[\mu} \mathcal{V}_2^{\nu]} + 2u^\lambda u^{[\mu} \mathcal{V}_2^{\nu]} \right]$$

Solution II: Vector-axial-tensor

Thermodynamics:

$$\begin{aligned} p &= -\epsilon/3 = -32\pi^4 \kappa T^4 \left(1 + \frac{\mathcal{C}^{\alpha\beta} \mathcal{C}_{\alpha\beta} - 2\mathcal{V}_2^2 - 6\mu_A^2}{12\pi^2 T^2} \right) \\ q^\mu &= -16\pi^2 \kappa T^2 (\mathcal{C}_\alpha^\mu \mathcal{A}_1^\alpha - 2\mu_A \mathcal{A}_1^\mu + \epsilon^{\mu\alpha\rho\sigma} u_\alpha \mathcal{A}_{1\rho} \mathcal{V}_{2\sigma}) , \bar{q}^\mu = 0 \\ \pi^{\mu\nu} &= -16\pi^2 \kappa T^2 \left(C^{\mu\alpha} C_\alpha^\nu + \mathcal{V}_2^\mu \mathcal{V}_2^\nu - \frac{1}{3} (\mathcal{C}_\beta^\alpha \mathcal{C}_\alpha^\beta + \mathcal{V}_2^2) \Delta^{\mu\nu} - \mu_A \mathcal{C}^{\mu\nu} \right) \\ \tau^{\mu\nu} &= 16\pi^2 \kappa T^2 u_\rho \mathcal{V}_2^\beta \epsilon^{\mu\nu\rho\sigma} (\mathcal{C}_\beta^\sigma + \mu_A \delta_\beta^\sigma) \end{aligned}$$

Spin current:

$$S^{\lambda\mu\nu} = 32\pi^2 \kappa T^2 \left[\mu_A \epsilon^{\lambda\mu\nu\alpha} u_\alpha - u_\alpha \epsilon^{\lambda\alpha\beta[\mu} (\mathcal{C}_{\beta}^{\nu]} + u^{\nu]} \mathcal{A}_{1\beta}) + \Delta^{\lambda[\mu} \mathcal{V}_2^{\nu]} + 2u^\lambda u^{[\mu} \mathcal{V}_2^{\nu]} \right]$$

- Conformal fluid with spin sources
- Non-trivial spin current *even for symmetric stress tensor*
- Non-vanishing intrinsic torque and energy currents

Solutions at first order

Fluid-gravity correspondence: $O(\partial)$ transport to corrections to background

Bhattacharyya, Hubeny, Minwalla, Rangamani '08

Boosted black brane with corrections:

$$ds^2 = \frac{dr^2}{r^2 g^2} + r^2 \left[-f^2 u_\mu u_\nu + h^2 \Delta_{\mu\nu} + j \left(2h c_I u_\mu v_\nu^I - 2f b_I u_\mu v_\nu^I + h d_I t_{\mu\nu}^I \right) \right] dx^\mu dx^\nu$$

$\{c_I v_\mu^I, b_I v_\mu^I, d_I t_{\mu\nu}^I\}$ are linearly independent combinations of $u^\mu, T, \omega_\mu^{ab}$

For simplicity we only consider $\mu_A \neq 0$ and $\mathcal{V}_2^\mu \neq 0$

Transport for $\mu_A \neq 0$

Energy transport:

$$\bar{q}^\mu = -\frac{64\pi^4\kappa T^4}{\mu_A}\omega^\mu,$$
$$q^\mu = -\frac{64\pi^4\kappa T^4}{\mu_A}\omega^\mu \left(1 - \frac{\mu_A^2}{2\pi^2 T^2}\right),$$
$$\pi^{\mu\nu} = 0,$$
$$\tau^{\mu\nu} = -16\pi^2\kappa T^2\epsilon^{\mu\nu\alpha\beta}u_\alpha\partial_\beta\mu_A$$

Spin transport:

$$J_A^\mu = 32\pi^2\kappa T^2 (\mu_A u^\mu + \omega^\mu) .$$

Transport for $\mu_A \neq 0$

Energy transport:

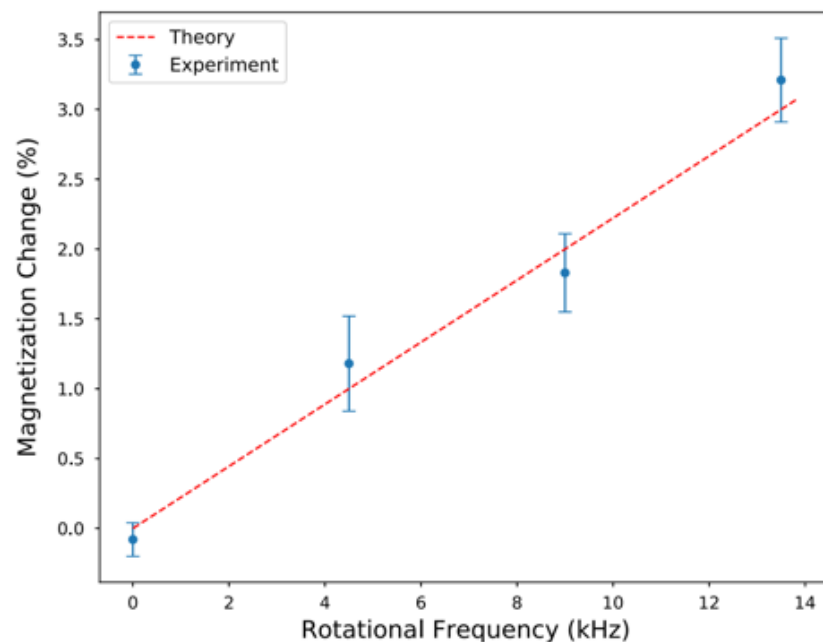
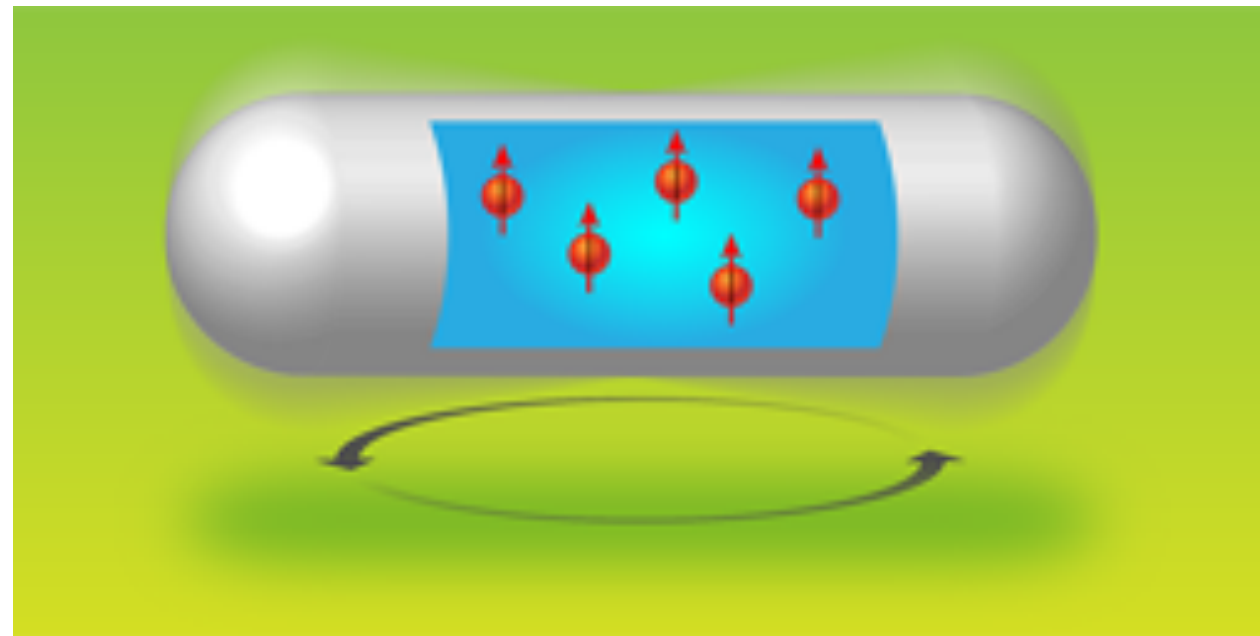
$$\bar{q}^\mu = -\frac{64\pi^4\kappa T^4}{\mu_A}\omega^\mu,$$
$$q^\mu = -\frac{64\pi^4\kappa T^4}{\mu_A}\omega^\mu \left(1 - \frac{\mu_A^2}{2\pi^2 T^2}\right),$$
$$\pi^{\mu\nu} = 0,$$
$$\tau^{\mu\nu} = -16\pi^2\kappa T^2\epsilon^{\mu\nu\alpha\beta}u_\alpha\partial_\beta\mu_A$$

Spin transport: $J_A^\mu = 32\pi^2\kappa T^2 (\mu_A u^\mu + \omega^\mu) .$

- Transport similar to chiral vortical separation effect: *generation of chiral imbalance by vorticity*.
- Here also an intrinsic torque
- More akin to “*chiral torsional effect*” Imaki and Yamamoto, ‘19

The Barnett effect

Magnetization of an uncharged body by rotation S. Barnett, 1915



$$J_A^\mu = 32\pi^2 \kappa T^2 (\mu_A u^\mu + \omega^\mu) .$$

Nuclear Barnett effect

Arabgol and Sleator '19

Transport for $\mathcal{V}_2^\mu \neq 0$

Novel possibilities for spin transport:

$$J_V^\mu = \frac{32}{3} \pi^2 T^2 (\nabla \cdot u) u^\mu ,$$

$$\bar{J}_V^\mu = 32 \pi^2 \kappa T^2 \mathcal{V}_2^\mu ,$$

$$\bar{J}_A^\mu = 16 \pi^2 \kappa T^2 \left(\eta^{\mu\nu} - \frac{\mathcal{V}_2^\mu \mathcal{V}_2^\nu}{\mathcal{V}_2^2} \right) \omega_\nu$$

- Spin current from divergence
- “Vortical-Hall” spin current

Summary

- Strong demand to develop relativistic hydrodynamics with spin
- Torsion acts as a source for the spin current
- Derived the hydro equations and constructed constitutive relations
- Novel effects from hydrodynamic effective action
- Lovelock-Chern-Simons provides a fruitful toy theory
- Characterization of spin transport coefficients from holography

Outlook

- A systematic study of relativistic spin transport:
Entropy generation, Onsager relations, etc.
- Analysis at weak coupling $\implies$ universal transport coefficients?
- Derivation in more general holographic theories: e.g. supergravity
- Observable effects of the novel spin transport?