

Weyl Cosmology

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Selected Topics in Mathematical Physics — 2021

A little bit of mathematics

Differential geometry

Metric tensor $g_{\mu\nu} \Rightarrow$ interval

$$ds^2 = g_{\mu\nu}(x) dx^\mu dx^\nu \quad (g_{\mu\nu} g^{\nu\lambda} = \delta_\mu^\lambda)$$

Connections $\Gamma_{\mu\nu}^\lambda \Rightarrow$ covariant derivative

$$\nabla_\mu l^\nu = l^\nu_{;\mu} + \Gamma_{\sigma\mu}^\nu l^\sigma, \dots$$

Curvature tensor $R_{\nu\lambda\sigma}^\mu$

$$R_{\nu\lambda\sigma}^\mu = \frac{\partial \Gamma_{\nu\sigma}^\mu}{\partial x^\lambda} - \frac{\partial \Gamma_{\nu\lambda}^\mu}{\partial x^\sigma} + \Gamma_{\kappa\lambda}^\mu \Gamma_{\nu\sigma}^\kappa - \Gamma_{\kappa\sigma}^\mu \Gamma_{\nu\lambda}^\kappa$$

Ricci tensor $R_{\mu\nu} = R_{\mu\lambda\nu}^\lambda$

Curvature scalar $R = g^{\mu\lambda} R_{\mu\lambda}$

$$g_{\mu\nu}(x), \quad \Gamma_{\mu\nu}^\lambda(x)$$

A little bit of mathematics

Differential geometry II

Three tensors $g_{\mu\nu}(x)$, $S_{\mu\nu}^{\lambda}$, $Q_{\lambda\mu\nu}$

Torsion $S_{\mu\nu}^{\lambda} = \Gamma_{\mu\nu}^{\lambda} - \Gamma_{\nu\mu}^{\lambda}$

Nonmetricity $Q_{\lambda\mu\nu} = \nabla_{\lambda} g_{\mu\nu}$

Connections $\Gamma_{\mu\nu}^{\lambda} = C_{\mu\nu}^{\lambda} + K_{\mu\nu}^{\lambda} + L_{\mu\nu}^{\lambda}$

$$C_{\mu\nu}^{\lambda} = \frac{1}{2} g^{\lambda\kappa} (g_{\kappa\mu,\nu} + g_{\kappa\nu,\mu} - g_{\mu\nu,\kappa})$$

$$Q_{\lambda\mu\nu} = \nabla_{\lambda} g_{\mu\nu} \quad \Rightarrow \quad Q_{\lambda\mu\nu} = Q_{\lambda\nu\mu}$$

$$K_{\mu\nu}^{\lambda} = \frac{1}{2} (S_{\mu\nu}^{\lambda} - S_{\mu}^{\lambda}{}_{\nu} - S_{\nu}^{\lambda}{}_{\mu})$$

$$L_{\mu\nu}^{\lambda} = \frac{1}{2} (Q_{\mu\nu}^{\lambda} - Q_{\mu}^{\lambda}{}_{\nu} - Q_{\nu}^{\lambda}{}_{\mu})$$

A little bit of mathematics

Differential geometry III

Riemann geometry

$$S_{\mu\nu}^{\lambda} = 0, \quad Q_{\lambda\mu\nu} = 0$$

$$S_{\mu\nu}^{\lambda} = 0, \quad Q_{\lambda\mu\nu} = 0, \quad \Rightarrow \quad \Gamma_{\mu\nu}^{\lambda} = C_{\mu\nu}^{\lambda}$$

Weyl geometry

$$S_{\mu\nu}^{\lambda} = 0 \quad \Rightarrow \quad \Gamma_{\mu\nu}^{\lambda} = \Gamma_{\nu\mu}^{\lambda}$$

$$Q_{\lambda\mu\nu} = A_{\lambda}(x)g_{\mu\nu}(x)$$

$$\Gamma_{\mu\nu}^{\lambda} = C_{\mu\nu}^{\lambda} + W_{\mu\nu}^{\lambda}$$

$$W_{\mu\nu}^{\lambda} = -\frac{1}{2}(A_{\mu}\delta_{\nu}^{\lambda} + A_{\nu}\delta_{\mu}^{\lambda} - A^{\lambda}g_{\mu\nu})$$

$A_{\lambda}(x)$ — “Weyl vector”

Local conformal transformation

$$ds^2 = \Omega^2(x) d\hat{s}^2 = \Omega^2 \hat{g}_{\mu\nu}(x) dx^\mu dx^\nu$$

$$g_{\mu\nu} = \Omega^2 \hat{g}_{\mu\nu}, \quad g^{\mu\nu} = \frac{1}{\Omega^2} \hat{g}^{\mu\nu}$$

$$\sqrt{-g} = \Omega^4(x) \sqrt{-\hat{g}}$$

$$C_{\mu\nu}^\lambda = \hat{C}_{\mu\nu}^\lambda + \left(\frac{\Omega_{,\nu}}{\Omega} \delta_\mu^\lambda + \frac{\Omega_{,\mu}}{\Omega} \delta_\nu^\lambda - g^{\lambda\kappa} \frac{\Omega_{,\kappa}}{\Omega} \hat{g}_{\mu\nu} \right)$$

If A_μ = gauge field

$$A_\mu = \hat{A}_\mu + 2 \frac{\Omega_{,\mu}}{\Omega} \Rightarrow$$

$$\Gamma_{\mu\nu}^\lambda = \hat{\Gamma}_{\mu\nu}^\lambda \Rightarrow$$

$$R^\mu{}_{\nu\lambda\sigma} = \hat{R}^\mu{}_{\nu\lambda\sigma}, \quad R_{\mu\nu} = \hat{R}_{\mu\nu}$$

Strength tensor

$$F_{\mu\nu} = \nabla_\mu A_\nu - \nabla_\nu A_\mu = A_{\nu,\mu} - A_{\mu,\nu}$$

Weyl gravity

$$\mathcal{L}_W = \alpha_1 R_{\mu\nu\lambda\sigma} R^{\mu\nu\lambda\sigma} + \alpha_2 R_{\mu\nu} R^{\mu\nu} + \alpha_3 R^2 + \alpha_4 F_{\mu\nu} F^{\mu\nu}$$

$$S_W = \int \mathcal{L}_W \sqrt{-g} d^4x, \quad \frac{\delta S_W}{\delta \Omega} = 0$$

Dynamical variables: $g_{\mu\nu}(x)$, $A_\mu x$

Matter fields $S_{\text{tot}} = S_W + S_m$, $S_m = \int \mathcal{L}_m \sqrt{-g} d^4x$

S_m is not necessary conformal invariant

δS_m does!

$$\begin{aligned} \delta S_m \stackrel{\text{def}}{=} & -\frac{1}{2} \int T^{\mu\nu} (\delta g_{\mu\nu}) \sqrt{-g} d^4x - \int G^\mu (\delta A_\mu) \sqrt{-g} d^4x \\ & + \int \frac{\delta \mathcal{L}_m}{\delta \Psi} (\delta \Psi) \sqrt{-g} d^4x = 0 \end{aligned}$$

Ψ — collective dynamical variable for matter fields, $\delta \mathcal{L}_m / \delta \Psi = 0$

$T^{\mu\nu}$ — energy-momentum tensor, G^μ — “Weyl current”

Weyl gravity II

$$\delta g_{\mu\nu} = \frac{2}{\Omega} g_{\mu\nu} (\delta\Omega) + \Omega^2 (\delta \hat{g}_{\mu\nu})$$

$$\delta A_\mu = (\delta A_\mu) + 2(\delta(\log \Omega))_{,\mu}$$

$$\delta \hat{g}_{\mu\nu} : \frac{\delta S_m}{\delta \hat{g}_{\mu\nu}} \stackrel{\text{def}}{=} -\frac{1}{2} \hat{T}_{\mu\nu} \Rightarrow \hat{T}_{\mu\nu} = \Omega^2 T_{\mu\nu}$$

$$(\hat{T}_\mu^\nu = \Omega^4 T_\mu^\nu, \hat{T}^{\mu\nu} = \Omega^6 T^{\mu\nu})$$

$$\delta \hat{A}_\mu : \frac{\delta S_m}{\delta \hat{A}_\mu} \stackrel{\text{def}}{=} -\hat{G}^\mu \Rightarrow \hat{G}^\mu = \Omega^4 G^\mu$$

Self-consistency condition

$$\delta\Omega : 2G_{;\mu}^\mu = \text{Trace}[T^{\mu\nu}]$$

“ $;$ ” — covariant derivative with $C_{\mu\nu}^\lambda$

Cosmology

Cosmology = homogeneity and isotropy
Robertson-Walker metric

$$ds^2 = dt^2 - a^2(t) \left(\frac{dr^2}{1 - kr^2} + r^2(d\theta^2 + \sin^2 \theta d\varphi^2) \right), \quad k = 0, \pm 1$$

$$A_\mu = (A_0(t), 0, 0, 0) \Rightarrow$$

$$F_{\mu\nu} \equiv 0$$

$$T^\mu_\nu = (T^0_0(t), T^1_1(t) = T^2_2 = T^3_3)$$

Special gauge

$$A_0(t) = 0$$

Field equations

$$\delta A_\mu : \quad -6\gamma\dot{R} = G^0, \quad \gamma = \frac{1}{3}(\alpha_1 + \alpha_2 + 3\alpha_3)$$

$$R = -6 \left(\frac{\ddot{a}}{a} + \frac{\dot{a}^2 + k}{a^2} \right)$$

$$\delta g_{\mu\nu} : \quad -12\gamma \left\{ \frac{\dot{a}}{a} \dot{R} + R \left(\frac{R}{12} + \frac{\dot{a}^2 + k}{a^2} \right) \right\} = T_0^0$$

$$-4\gamma \left\{ \ddot{R} + 2\frac{\dot{a}}{a} \dot{R} - R \left(\frac{R}{12} + \frac{\dot{a}^2 + k}{a^2} \right) \right\} = T_1^1$$

Self-consistency condition:

$$2 \frac{(G^0 a^3)'}{a^3} = T_0^0 + 3T_1^1$$

Vacuum solutions

$$G^0 = 0 \Rightarrow \dot{R} = 0 \Rightarrow R = R_0 \Rightarrow$$

$$a^2 + k = -\frac{R_0}{12}a^2 + \frac{Q_0}{a^2} \Rightarrow$$

$$R_0 Q_0 = 0$$

1 $Q_0 = 0$

$R_0 < 0 \Rightarrow$ de Sitter

$R_0 > 0 \Rightarrow$ AdS

$R_0 = 0 \Rightarrow$ Milne universe

2 $Q_0 \neq 0 \quad R_0 = 0$

$4Q_0(t - t_0)^2 = a^2, \quad k = 0 \quad (Q_0 > 0)$

$a^2 + (t - t_0)^2 = Q_0, \quad k = \pm 1 \quad (Q_0 > 0)$

$(t - t_0)^2 - a^2, \quad k = -1$

New possibilities

Single particle u^μ

Riemann geometry

$$S_{\text{part}} = -m \int ds$$

Weyl geometry — new invariant $B = A_\mu u^\mu$

$$S_{\text{part}} = \int f_1(B) ds + \int f_2(B) d\tau = \int \{ f_1(B) \sqrt{g_{\mu\nu} u^\mu u^\nu} + f_2(B) \} d\tau$$

Equations of motion

$$\begin{aligned} f_1(B) u_{\lambda;\mu} u^\mu &= \left((f_1'(B) + f_2''(B)) A_\lambda - f_1'(B) u_\lambda \right) B_{,\mu} u^\mu \\ &+ (f_1'(B) + f_2'(B)) F_{\lambda\mu} u^\mu \end{aligned}$$

Since $F_{\lambda\mu} u^\lambda u^\mu \equiv 0$ and $u_{\lambda;\sigma} u^\lambda \equiv 0$

Either $(f_1' + f_2'')B = f_1'$ or $B_{,\mu} u^\mu = 0$

Independent observer

Conformal invariance

Main problem — gauge fixing

Independent observer

Independent = not a dynamical variable

Simplest choice $\Rightarrow$

$$S_{\text{obs}} = \int A_{\mu} \xi^{\mu} \sqrt{-g} d^4x,$$

ξ^{μ} — proportional to the four-velocity

$$T_{\text{obs}}^{\mu\nu} = -(A_{\lambda} \xi^{\lambda}) g^{\mu\nu}$$

Cosmology: it singles out the gauge $A_{\mu} = 0$

Vacuum cosmological solutions

$$A_\mu = 0, \quad G^\mu = 0, \quad T^{\mu\nu} = 0$$

Field equations

$$-6\gamma\dot{R} = \xi^0$$

$$-12\gamma\left\{\frac{\dot{a}}{a}\dot{R} + R\left(\frac{R}{12} + \frac{\dot{a}^2 + k}{a^2}\right)\right\} = 0$$

$$-4\gamma\left\{\ddot{R} + 2\frac{\dot{a}}{a}\dot{R} - R\left(\frac{R}{12} + \frac{\dot{a}^2 + k}{a^2}\right)\right\} = 0$$

$$R = -6\left(\frac{\ddot{a}}{a} + \frac{\dot{a}^2 + k}{a^2}\right)$$

Self-consistency condition:

$$2\frac{(G^0 a^3)'}{a^3} = 0 \quad \Rightarrow \quad \dot{R} = \frac{b}{a^3}, \quad b = \text{const}$$

Vacuum cosmological solutions II

$$A_\mu = 0, G^\mu = 0, T^{\mu\nu} = 0$$

$$\ddot{a}^2 a^2 = (\dot{a}^2 + k)^2 - \frac{1}{3} \dot{a} b$$

$$\dot{a} > 0, a \gg a_0 \Rightarrow \dot{a} \propto a \Rightarrow \text{exponential growth!}$$

The End