

# Holographic Anisotropic Model for Heavy Quarks in Anisotropic Hot Dense QGP: Magnetic Catalysis

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# My collaborators:

Holographic Anisotropic Model for Heavy Quarks  
in Anisotropic Hot Dense QGP: Magnetic Catalysis

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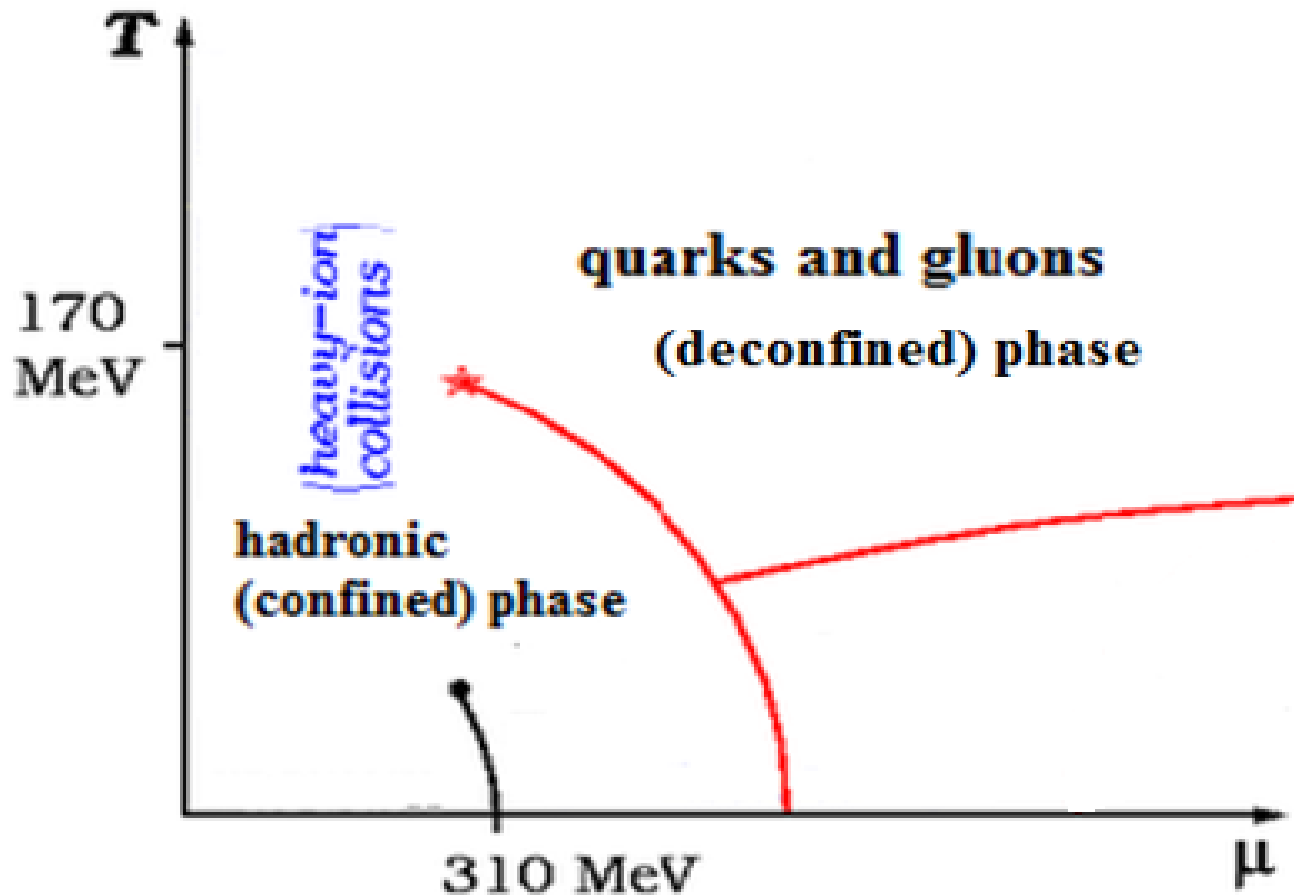
**Pavel Slepov** (MIAN)

Work is in progress.

# Outline:

- Introduction
- Set up a Question?
- Approach: AdS/CFT or Gauge/Gravity Duality
- Results
- Summary

# Introduction: (QCD phase diagram)



One of the most important open problems in the study of QCD at finite temperature and density is the determination of the **phase diagram** of the theory.

# Introduction: (Heavy Ion collision)

**RHIC**

Au-Au collision

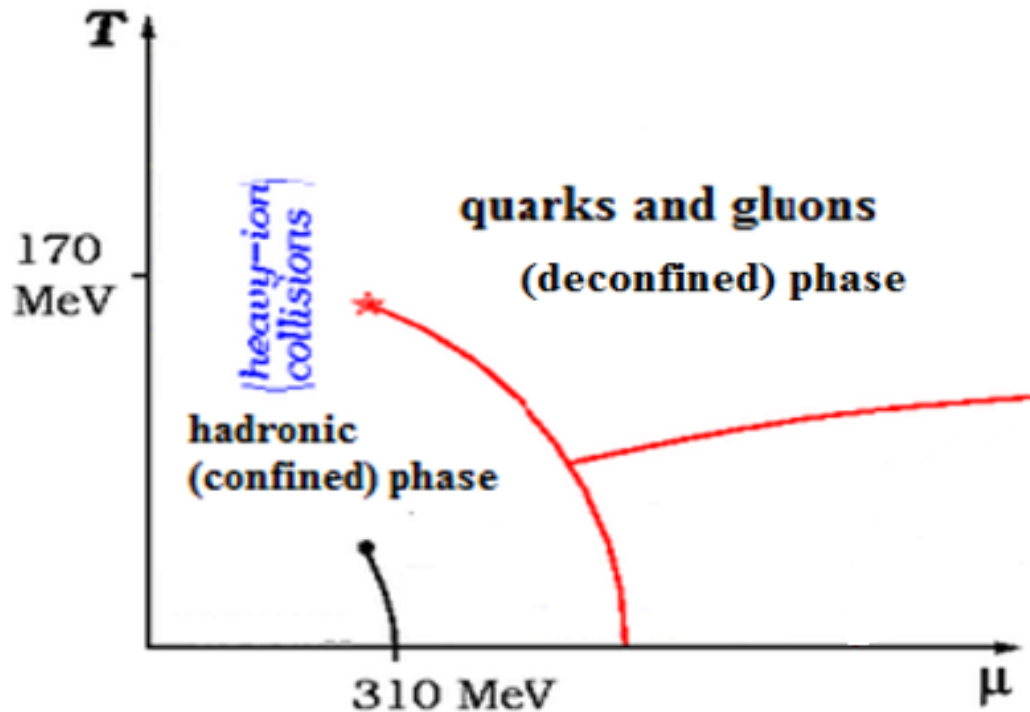
2Tc

**LHC**

Pb-Pb collision

5Tc

$T_c \sim 170 \text{ MeV}$



Experimental results show that:  
**QGP is strongly coupled**

$$\frac{\eta}{s} \sim 1/4\pi$$

# Introduction:

Heavy ion collisions (HIC):

Can teach us about properties of the high temperature **phase of QCD**.

Noncentral relativistic HIC  **Anisotropic** Plasma

There is a very strong **magnetic field**, in the early stages of relativistic heavy ion collision.

# Set up a Question:


What is the effect of the **magnetic field** on the phase transition temperature?  
(confinement/deconfinement)

Two phenomenon:

- 1- Inverse Magnetic Catalysis (IMC) (Mamo. '15, Aref'eva et al. '20 ,  
Dudal et al. '19, ....)
- 2- **Magnetic Catalysis (MC)** .....??? ( Gusynin et al. '94, He et al. '20, ...)

## 2nd Question:

What is the effect of the **anisotropy** on the phase (confinement/deconfinement) transition temperature?

Noncentral relativistic HIC  **Anisotropic** quark gluon plasma



## Approach:

QGP as a phase of QCD is **strongly coupled**.

So, we use **Non-Perturbative** approach:

**AdS/CFT duality** (That is an example of holographic principle)

# AdS/CFT correspondence: (Maldacena)

IIB String Theory on  $\text{AdS}_5 \times S^5 \Leftrightarrow \mathcal{N} = 4, d = 4$  SYM theory

Interesting limit:

Classical gravity  $\longleftrightarrow$  Strongly coupled QFT

*Gravity*  $\Leftrightarrow$  Gauge

# Methods:

Top-down models: Directly constructed from string theory:

D3-D7 model

J. Babington, J. Erdmenger, N. J. Evans, Z. Guralnik , I. Kirsch,  
M. Kruczenski, D. Mateos, R. C. Myers and D. J. Winters,...

D4-D8 model

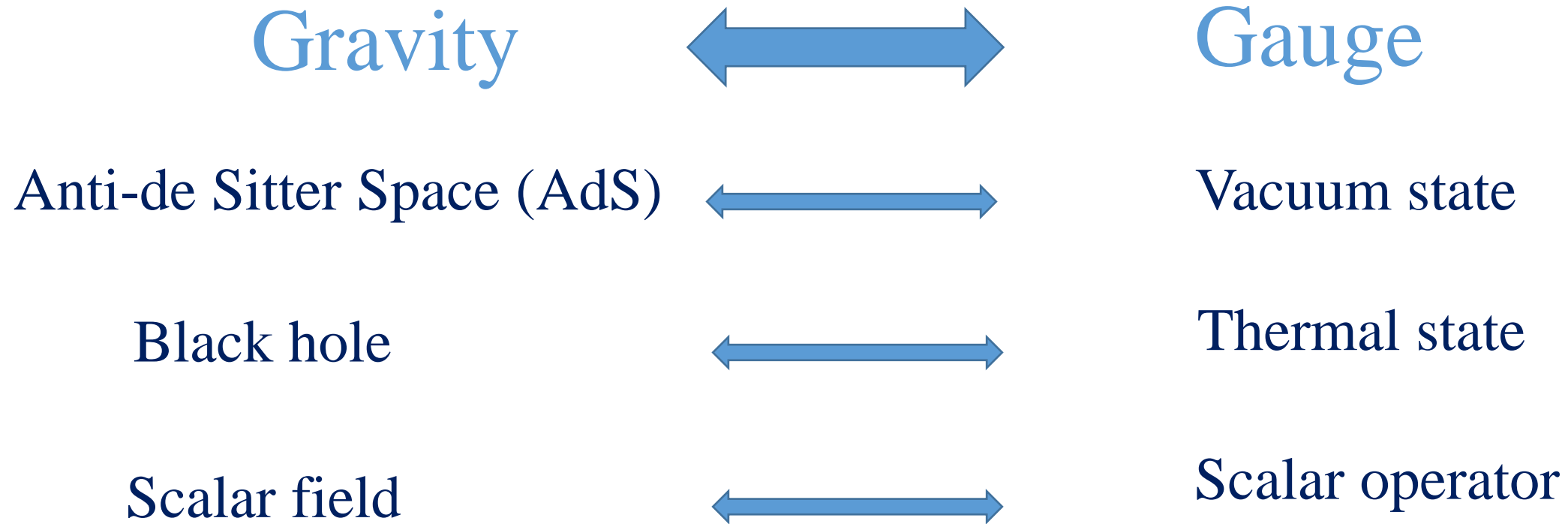
T. Sakai and S. Sugimoto

Bottom-up models: (phenomenological)

Introduce a dilaton field

J. Erlich, E. Katz, D. T. Son and M. A. Stephanov,  
A. Karch, B. Batell and T. Gherghetta, U. Gursoy, E. Kiritsis,...

# Duality:



# Our 1st Question:

What is the effect of the **magnetic field** on the phase transition temperature?  
(confinement/deconfinement)

**How can we set up a holographic approach?**

# Our Model:

$$\mathcal{L} = \sqrt{-g} \left[ R - \frac{f_0(\phi)}{4} F_0^2 - \frac{f_1(\phi)}{4} F_1^2 - \frac{f_3(\phi)}{4} F_3^2 - \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - V(\phi) \right]$$

$$\phi = \phi(z),$$

Electric ansatz  $F_0$ :  $A_0 = A_t(z)$ ,  $A_{i=1,2,3,4} = 0$ ,

Magnetic ansatz  $F_k$ :  $F_1 = q_1 dx^2 \wedge dx^3$ ,  $F_3 = q_3 dx^1 \wedge dx^2$ .

$F_0$   $\longleftrightarrow$  Chemical potential

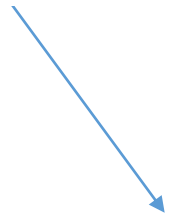
$F_1$   $\longleftrightarrow$  Primary anisotropy

$F_3$   $\longleftrightarrow$  Magnetic field

Our ansatz for the metric:

$$ds^2 = \frac{L^2}{z^2} \mathfrak{b}(z) \left[ -g(z) dt^2 + dx_1^2 + \left( \frac{z}{L} \right)^{2-\frac{2}{\nu}} dx_2^2 + e^{c_B z^2} \left( \frac{z}{L} \right)^{2-\frac{2}{\nu}} dx_3^2 + \frac{dz^2}{g(z)} \right]$$

$$\mathfrak{b}(z) = e^{2\mathcal{A}(z)}$$



Warp factor

isotropic

$$\nu = 1$$

Anisotropic

$$\nu = 4.5$$

Varying the Lagrangian with the ansatz of metric:

Equations of motion:

$$A_t'' + A_t' \left( \frac{b'}{2b} + \frac{f_0'}{f_0} + \frac{\nu - 2}{\nu z} + c_B z \right) = 0$$

$$g'' + g' \left( \frac{3b'}{2b} - \frac{\nu + 2}{\nu z} + c_B z \right) - \left( \frac{z}{L} \right)^2 \frac{f_0 (A_t')^2}{b} - \left( \frac{z}{L} \right)^{\frac{2}{\nu}} \frac{q_3^2 f_3}{b} = 0.$$

$$b'' - \frac{3(b')^2}{2b} + \frac{2b'}{z} - \frac{4b}{3\nu z^2} \left( 1 - \frac{1}{\nu} + \left( 1 - \frac{3\nu}{2} \right) c_B z^2 - \frac{\nu c_B^2 z^4}{2} \right) + \frac{b (\phi')^2}{3} = 0$$

$$2g' \left( 1 - \frac{1}{\nu} \right) + 3g \left( 1 - \frac{1}{\nu} \right) \left( \frac{b'}{b} - \frac{4(\nu + 1)}{3\nu z} + \frac{2c_B z}{3} \right) + \left( \frac{L}{z} \right)^{1 - \frac{4}{\nu}} \frac{L e^{-c_B z^2} q_1^2 f_1}{b} = 0$$



# Solving EOMs:

1<sup>st</sup> gauge field:

$$A_t'' + A_t' \left( \frac{b'}{2b} + \frac{f_0'}{f_0} + \frac{\nu - 2}{\nu z} + c_B z \right) = 0$$

Blackening function:

$$g'' + g' \left( \frac{3b'}{2b} - \frac{\nu + 2}{\nu z} + c_B z \right) - \left( \frac{z}{L} \right)^2 \frac{f_0 (A_t')^2}{b} - \left( \frac{z}{L} \right)^{\frac{2}{\nu}} \frac{q_3^2 f_3}{b} = 0$$

$$f_3 = 2 \left( \frac{L}{z} \right)^{\frac{2}{\nu}} b g \frac{c_B z}{q_3^2} \left( \frac{g'}{g} + \frac{3b'}{2b} - \frac{2}{\nu z} + c_B z \right)$$

$$g'' + g' \left( \frac{3b'}{2b} - \frac{\nu + 2}{\nu z} - c_B z \right) - 2g \left( \frac{3b'}{2b} - \frac{2}{\nu z} + c_B z \right) c_B z - \left( \frac{z}{L} \right)^2 \frac{f_0 (A_t')^2}{b} = 0.$$

# Magnetic catalysis for heavy quarks:

Warp factor:

$$b(z) = e^{2A(z)} = e^{-cz^2/2 - 2pz^4}$$

Gauge coupling  
function:

$$f_0 = e^{-(R_{gg} + \frac{cBq_3}{2})z^2} \frac{z^{-2+\frac{2}{\nu}}}{\sqrt{b}}$$

# Boundary conditions:

1<sup>st</sup> gauge field:

$$A_t(0) = \mu, \quad A_t(z_h) = 0$$

Blackening function:

$$g(0) = 1, \quad g(z_h) = 0$$

Dilaton field:

$$\phi(z_0) = 0$$

# In search of Magnetic Catalysis: (MC)

We need to find:  $g(z)$



Temperature and Entropy:

$$T = \frac{\sqrt{g_{tt}' g^{zz'}}}{4\pi} \Big|_{z=z_h} = \frac{\sqrt{g_{00}' g^{55'}}}{4\pi} \Big|_{z=z_h} = \frac{|g'|}{4\pi} \Big|_{z=z_h}$$

$$s = \frac{\sqrt{g_{xx} g_{y_1 y_1} g_{y_2 y_2}}}{4} \Big|_{z=z_h} = \frac{\sqrt{g_{11} g_{22} g_{33}}}{4} \Big|_{z=z_h}$$

$$s = \frac{1}{4} \left( \frac{L}{z_h} \right)^{1+\frac{2}{\nu}} e^{-(2R_{gg}-c_B)\frac{z_h^2}{2}-3pz_h^4}$$



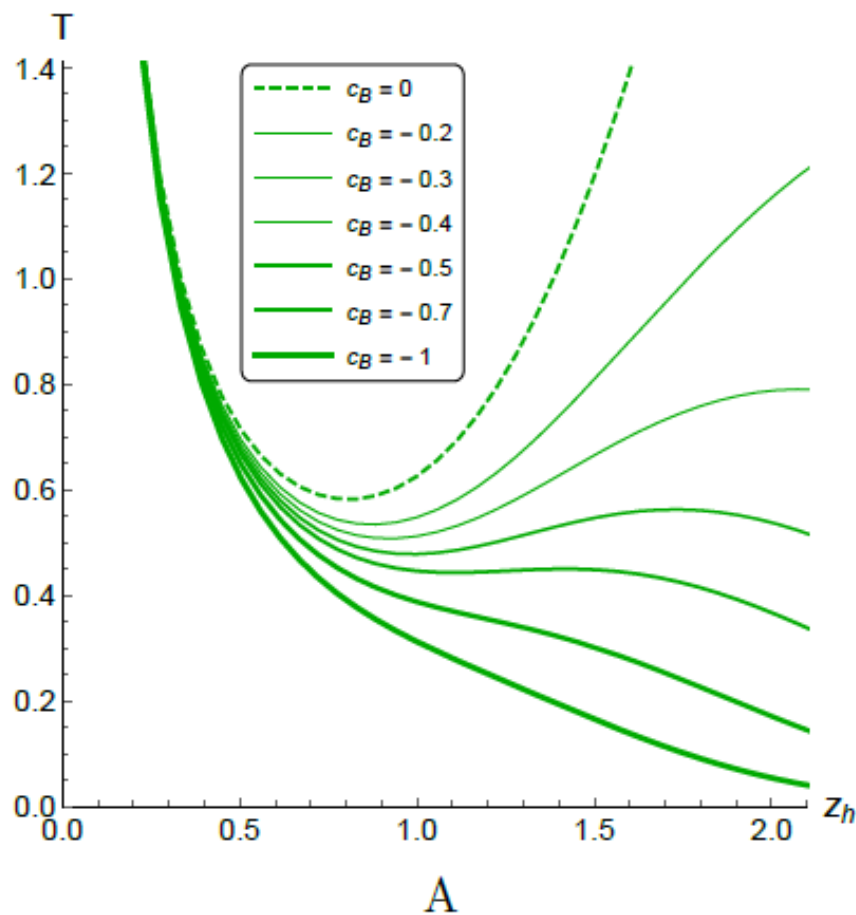
Free energy:

$$F = - \int s dT = \int_{z_h}^{\infty} s T' dz.$$

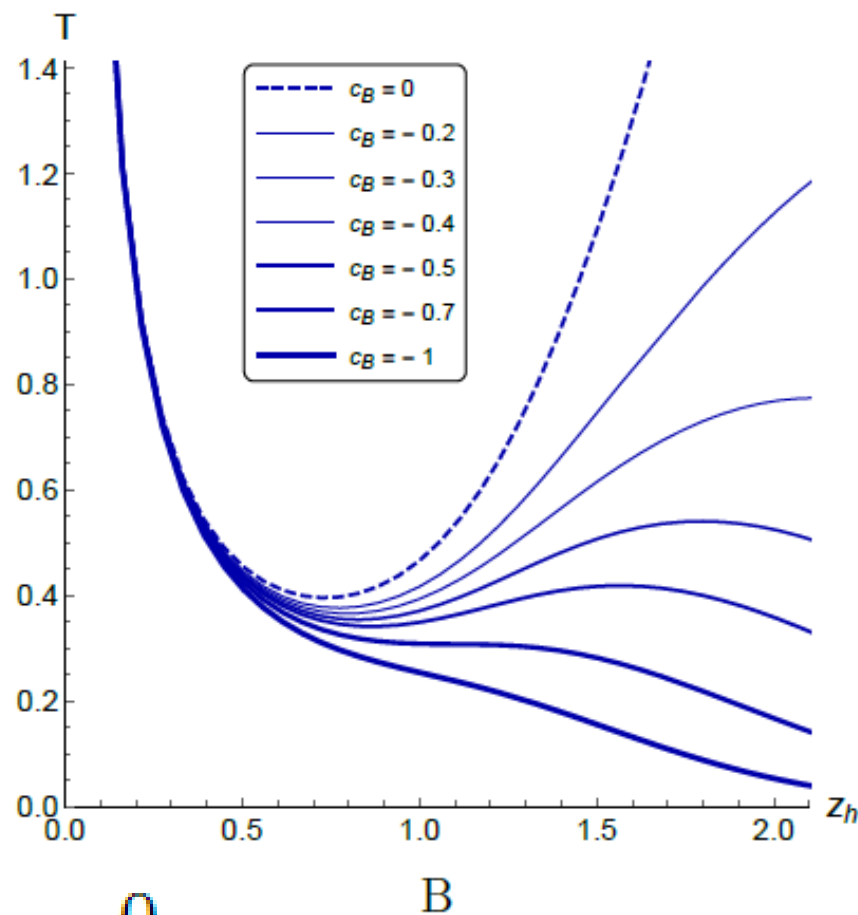
# Temperature:

Warp factor:

$$b(z) = e^{2\mathcal{A}(z)} = e^{-cz^2/2 - 2pz^4}$$



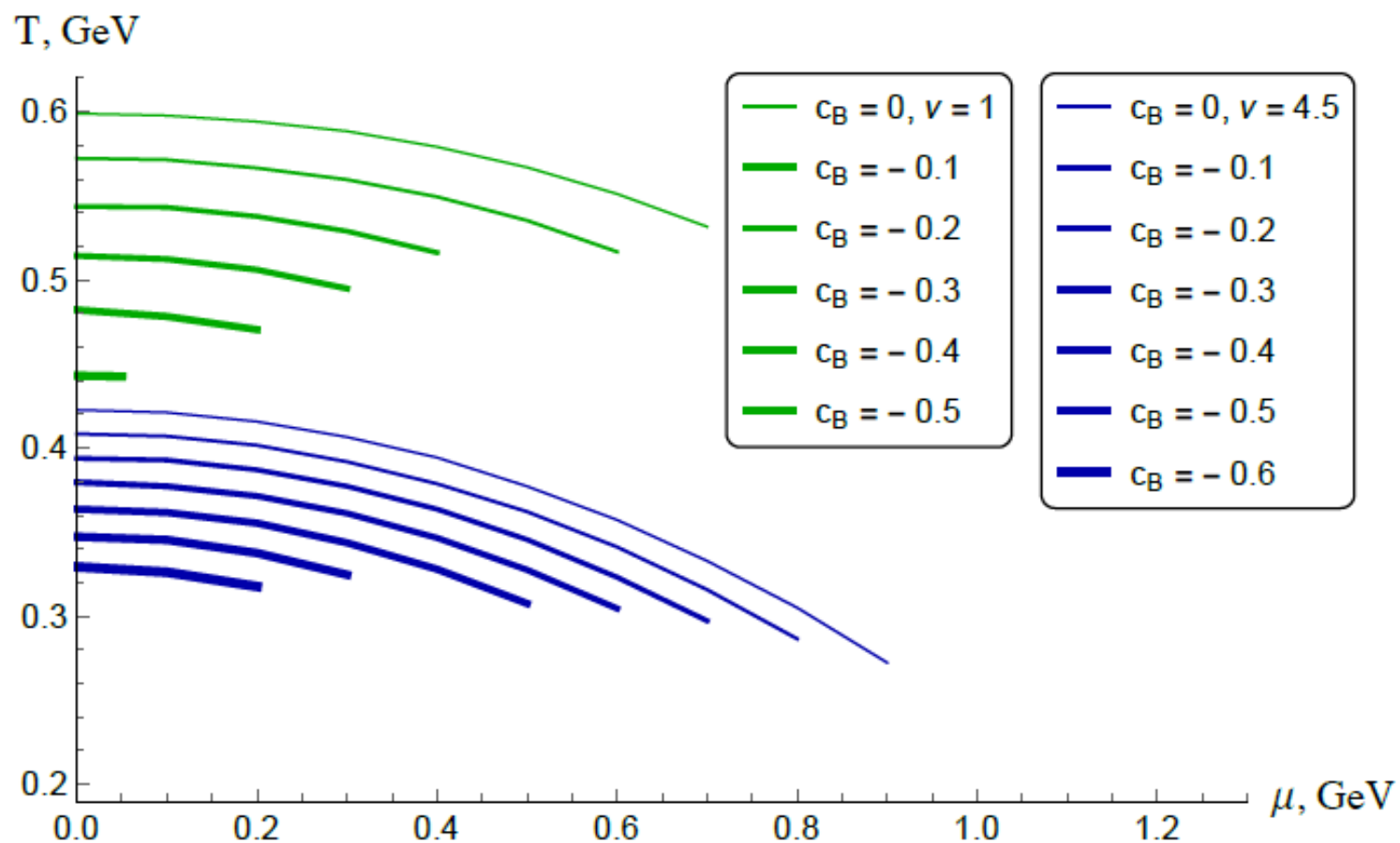
$$\mu = 0$$



# IMC:

Warp factor:

$$b(z) = e^{2\mathcal{A}(z)} = e^{-cz^2/2 - 2pz^4}$$



Warp factor:  $\mathfrak{b}(z) = e^{2\mathcal{A}(z)} = e^{-cz^2/2 - 2pz^4}$

NO MC phenomenon was observed!!!!

New Warp factor:  $\mathfrak{b}(z) = e^{2\mathcal{A}(z)} = e^{-cz^2/2 - 2(p - c_B q_3)z^4}$

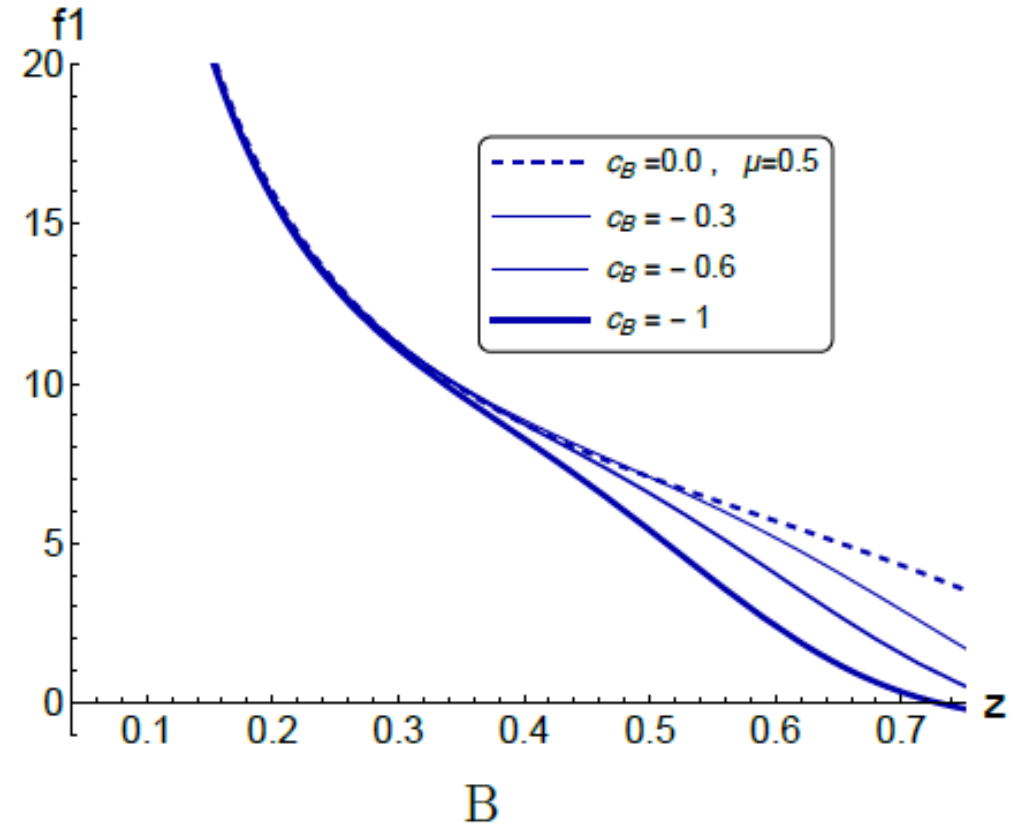
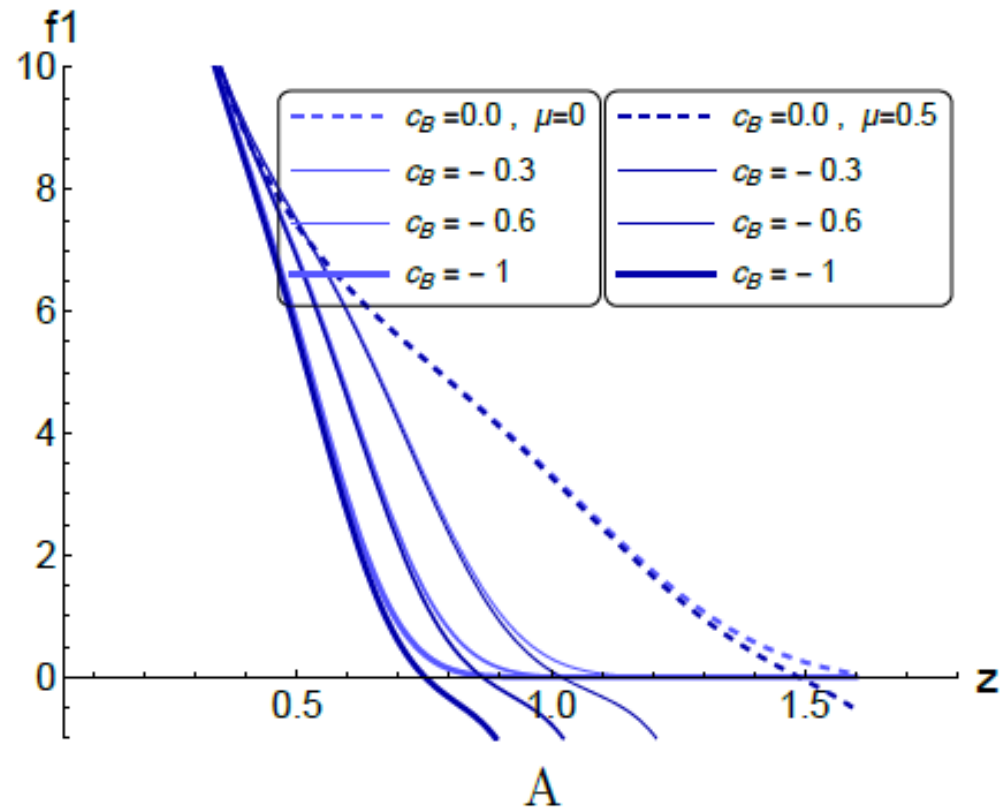
$$ds^2 = \frac{L^2}{z^2} \mathfrak{b}(z) \left[ -g(z) dt^2 + dx_1^2 + \left(\frac{z}{L}\right)^{2-\frac{2}{\nu}} dx_2^2 + e^{c_B z^2} \left(\frac{z}{L}\right)^{2-\frac{2}{\nu}} dx_3^2 + \frac{dz^2}{g(z)} \right]$$

Gauge coupling function:  $f_0 = e^{-(R_{gg} + \frac{c_B q_3}{2})z^2} \frac{z^{-2+\frac{2}{\nu}}}{\sqrt{\mathfrak{b}}}$

Our model can possess the Linear Regge trajectory for meson mass spectra.

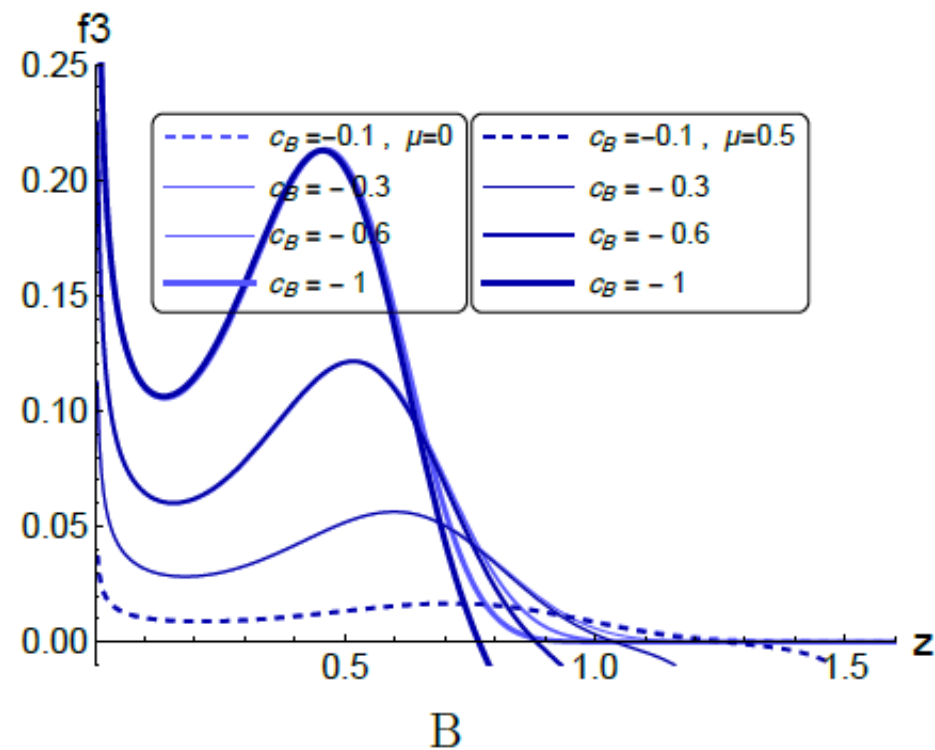
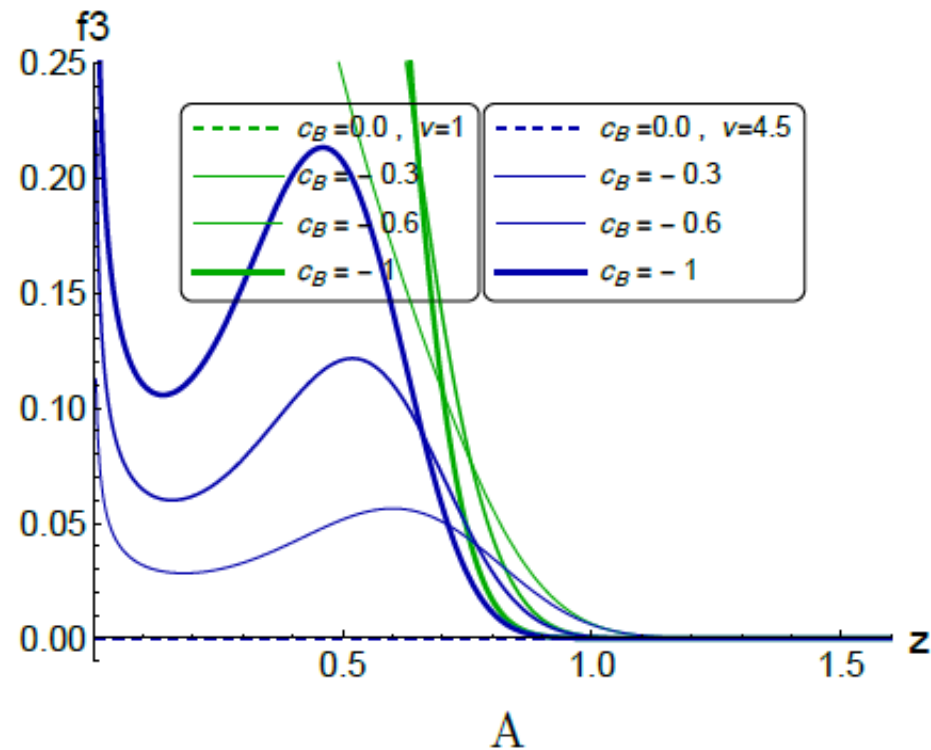


## 2<sup>nd</sup> coupling function:



The Null Energy Condition (NEC) should be satisfied by choosing the appropriate 2<sup>nd</sup> horizon.

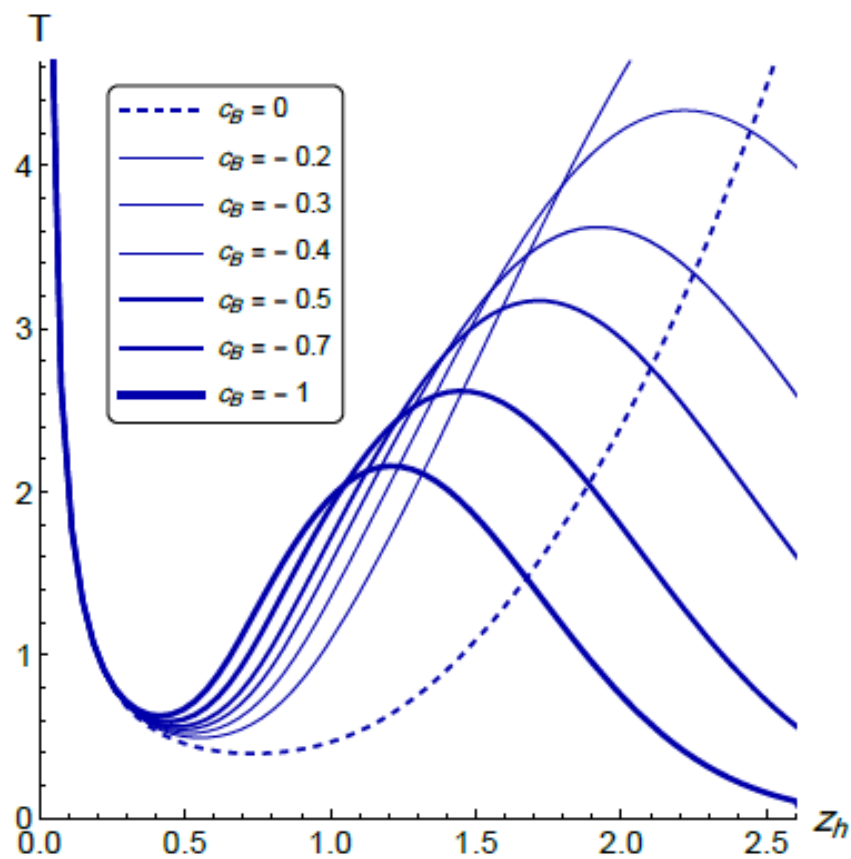
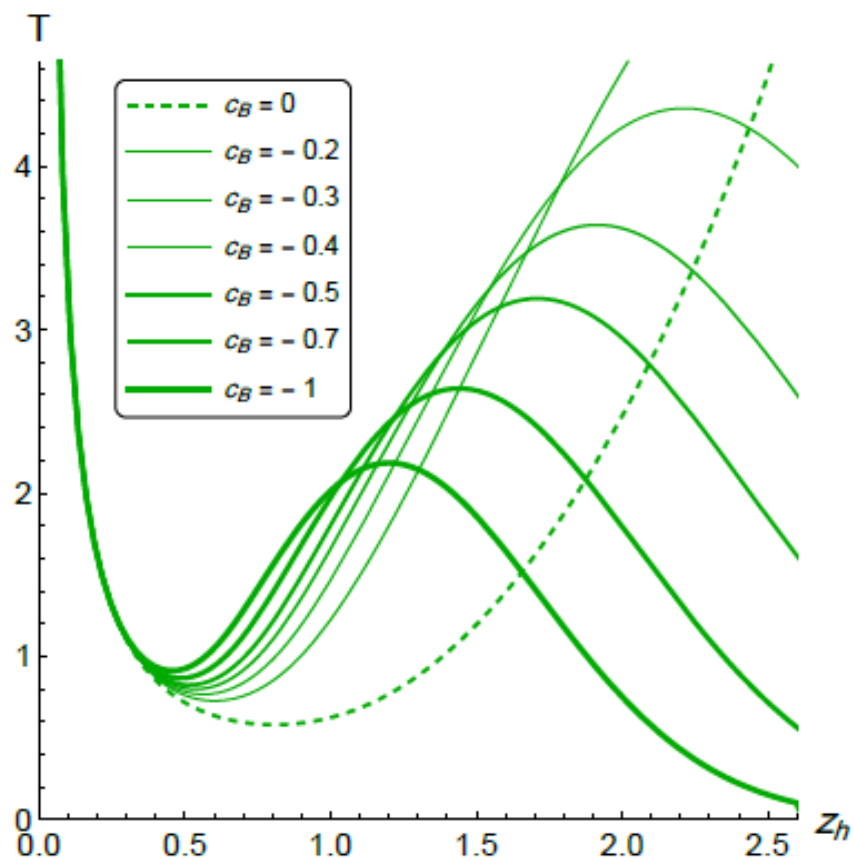
## 3rd coupling function:



The Null Energy Condition (NEC) should be satisfied by choosing the appropriate 2<sup>nd</sup> horizon.

# Temperature:

$$b(z) = e^{2A(z)} = e^{-cz^2/2 - 2(p - c_B q_3)z^4}$$



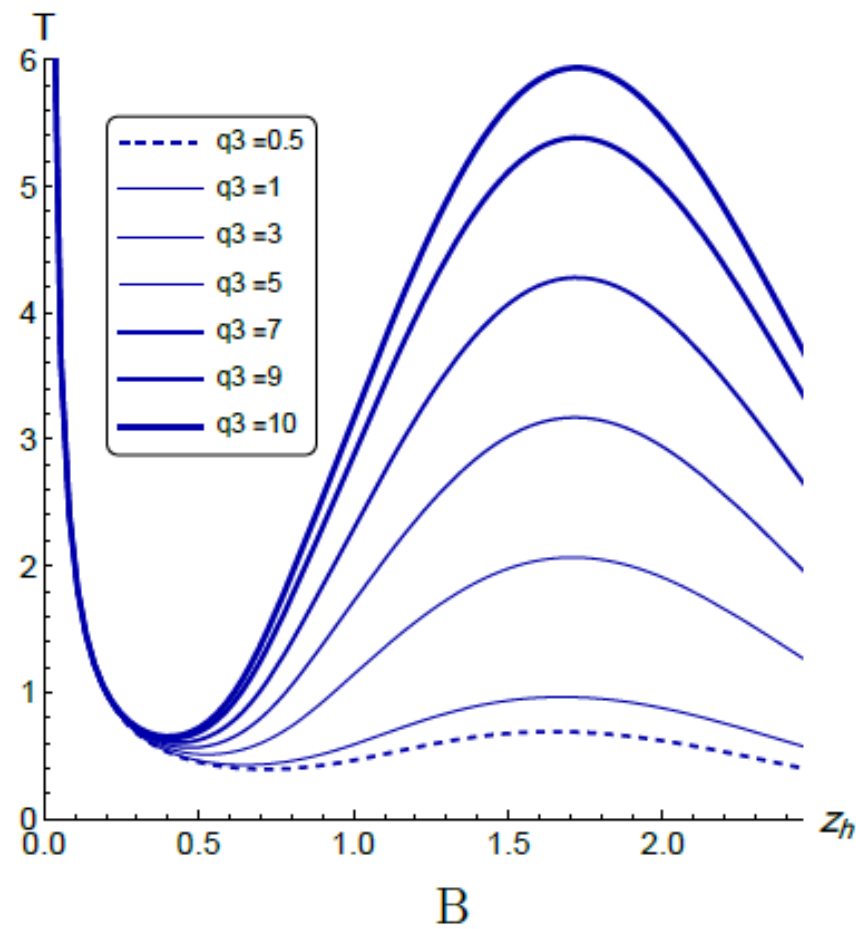
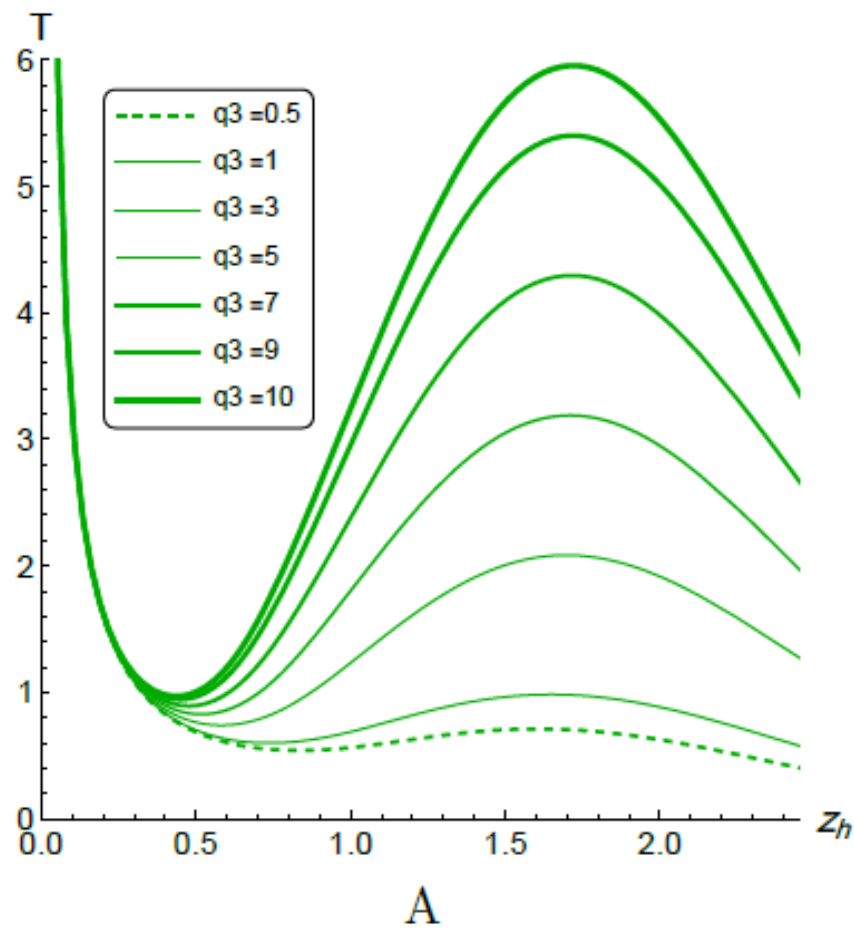
$$q_3 = 5$$

$$\mu = 0$$

Temperature:

$$\mu = 0$$

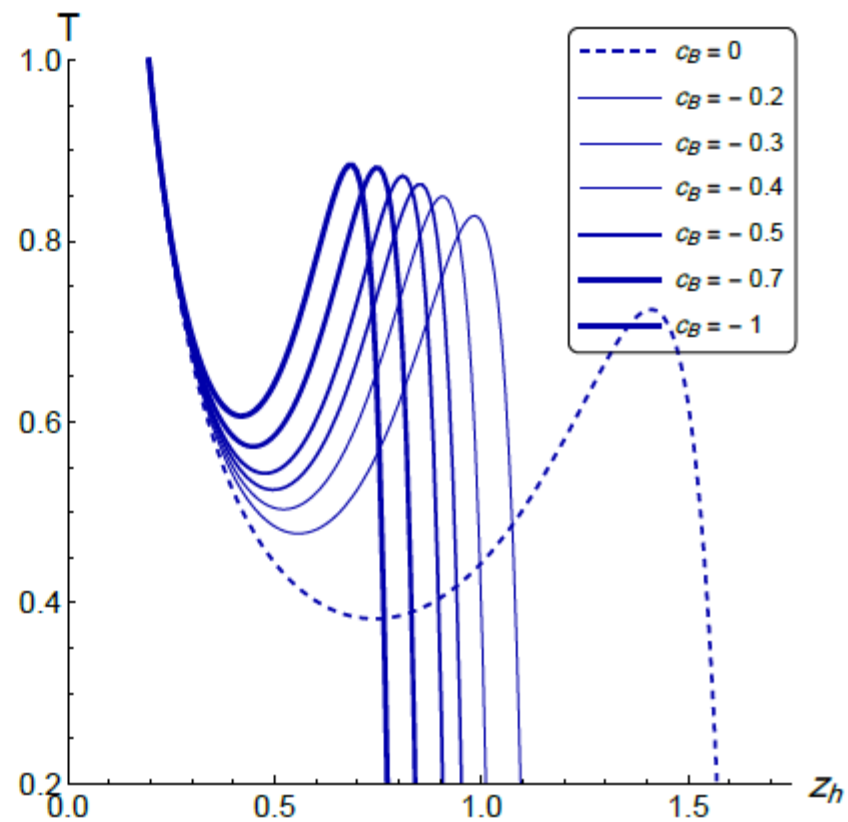
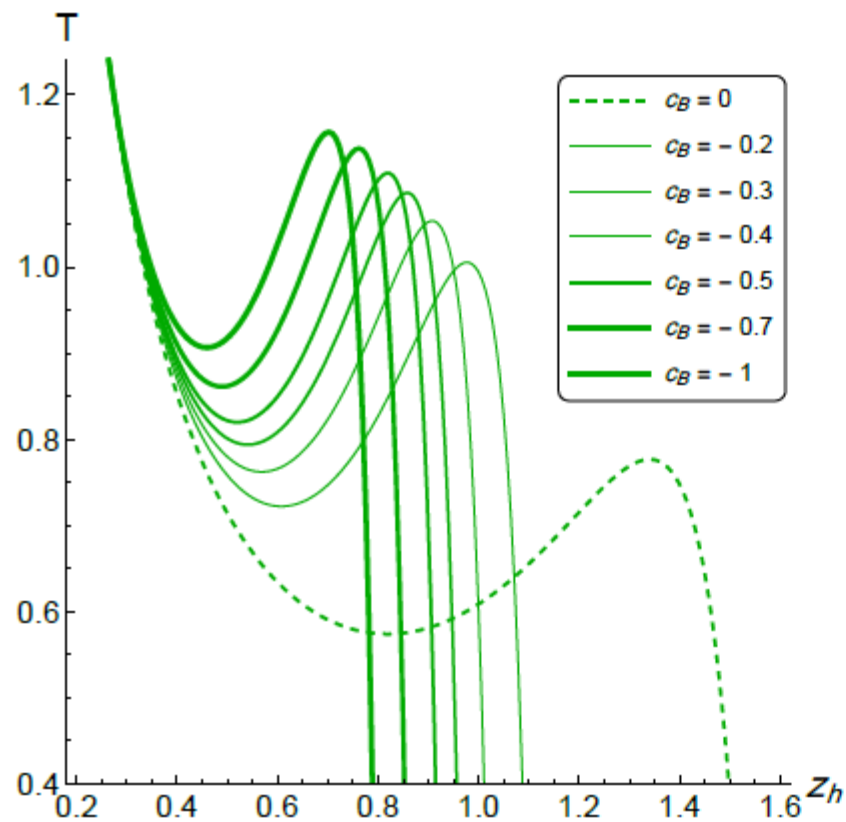
$$c_B = -0.5$$



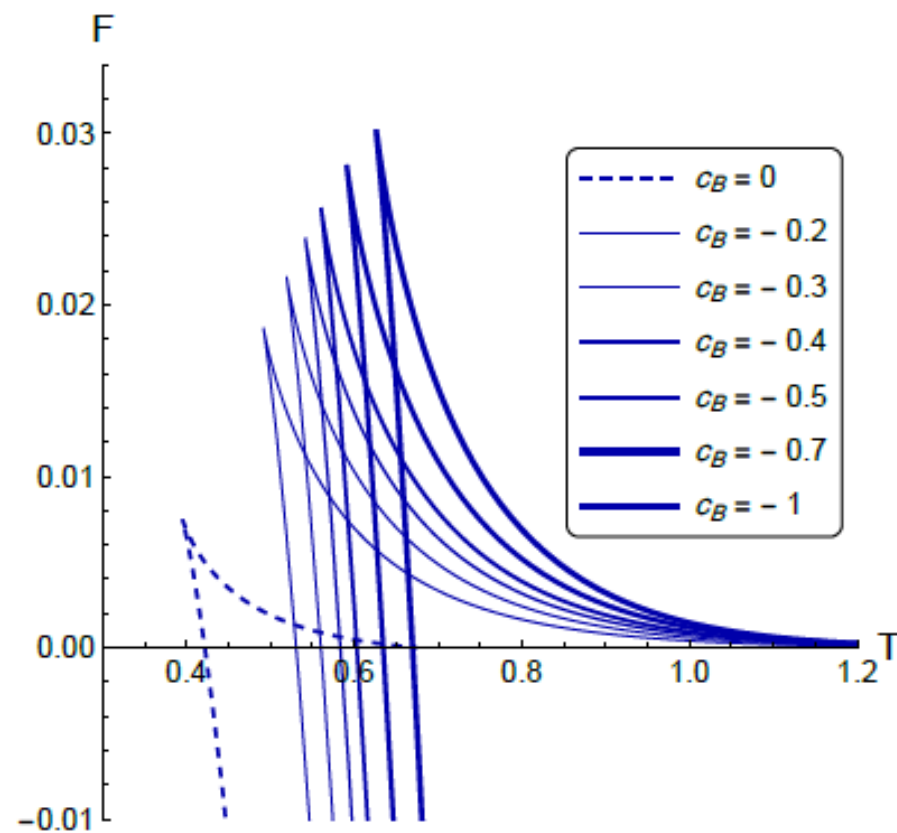
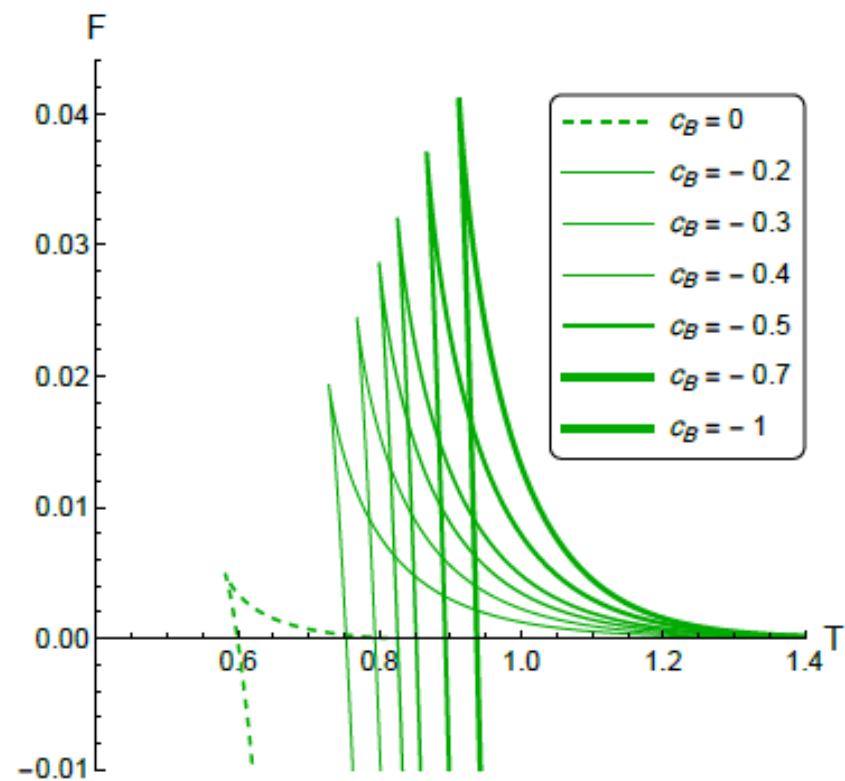
Temperature:

$$\mu = 0.3$$

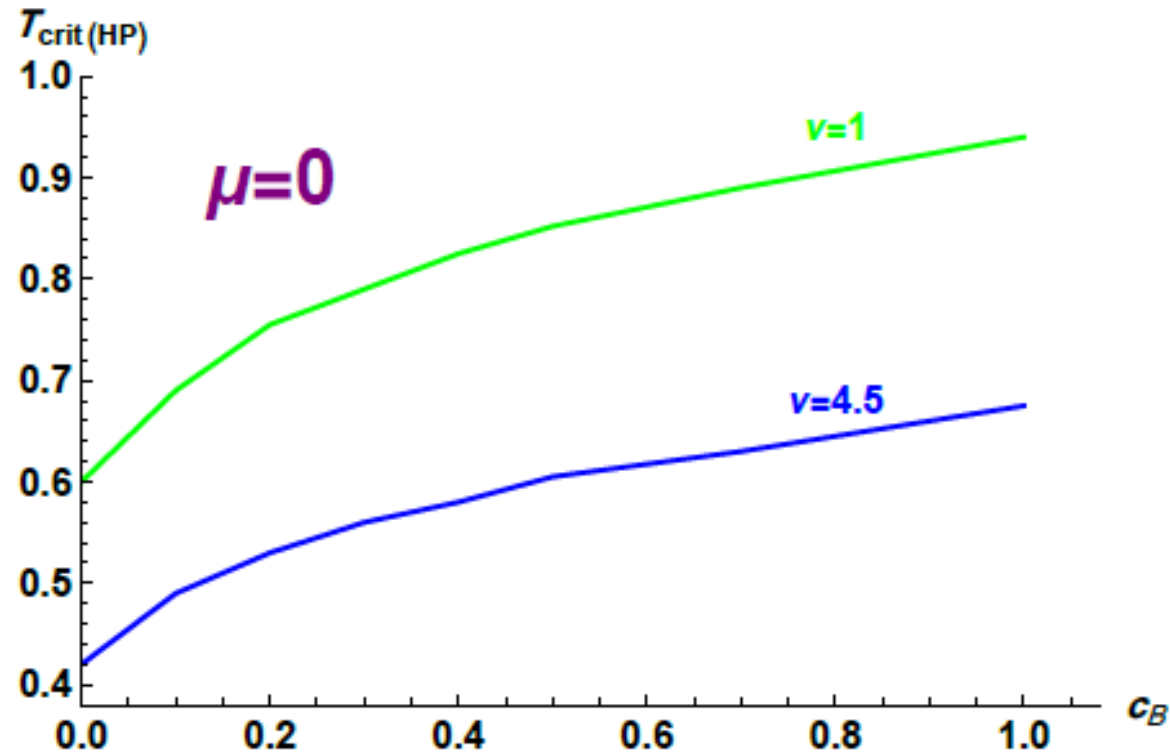
$$q_3 = 5$$



Free energy:  $\mu = 0$   $q_3 = 5$



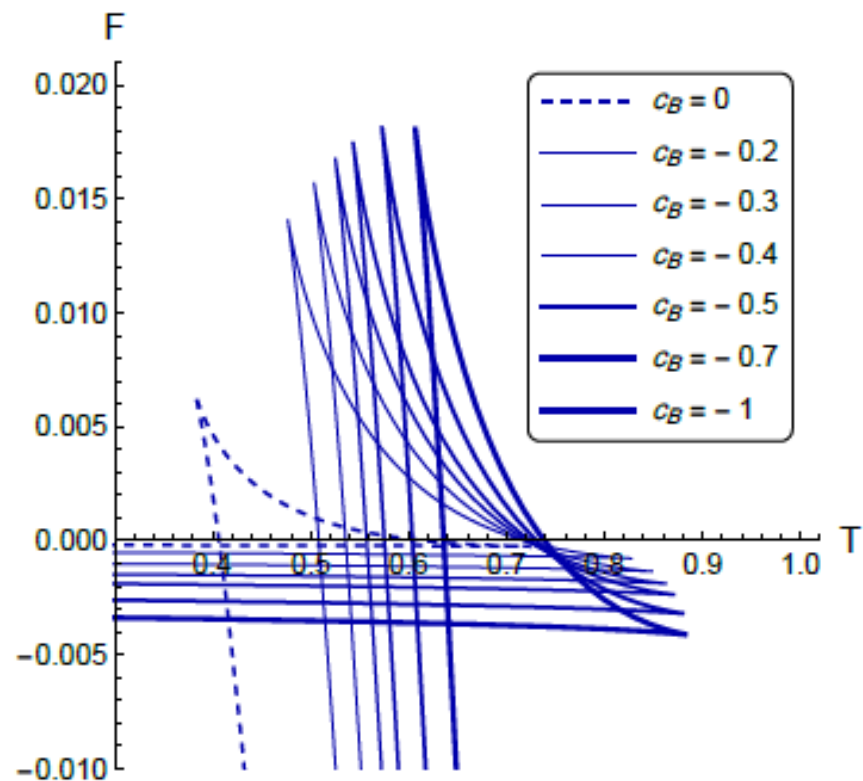
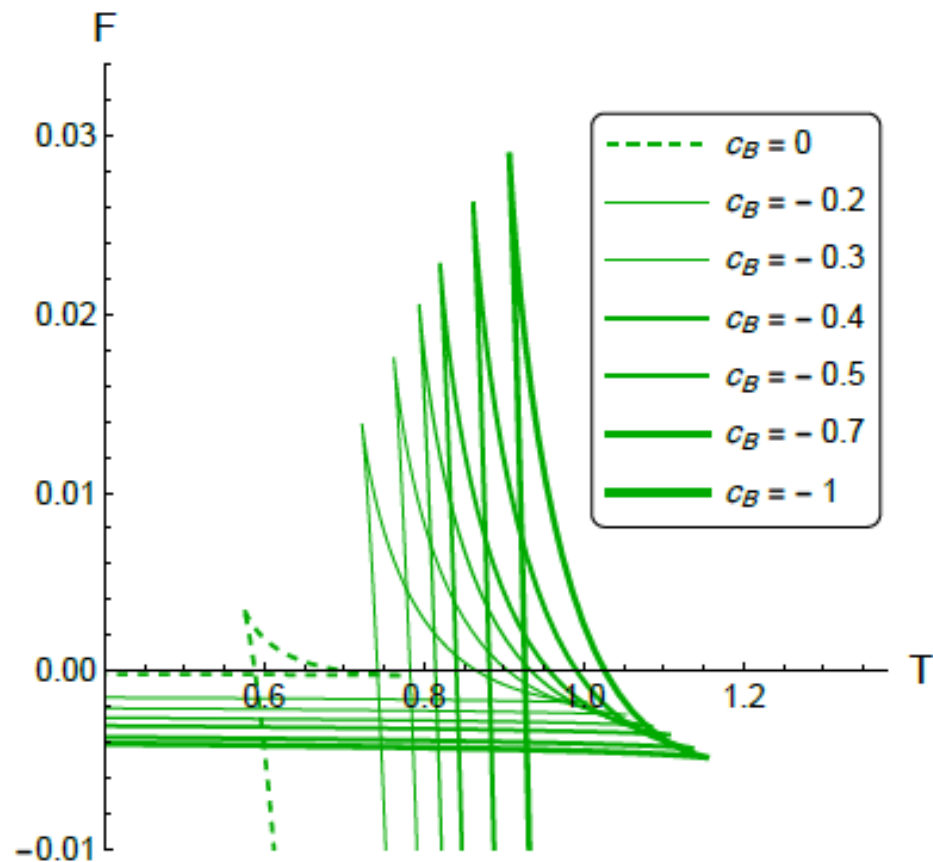
# Phase diagram:



$$q_3 = 5$$

MC phenomenon is obtained.

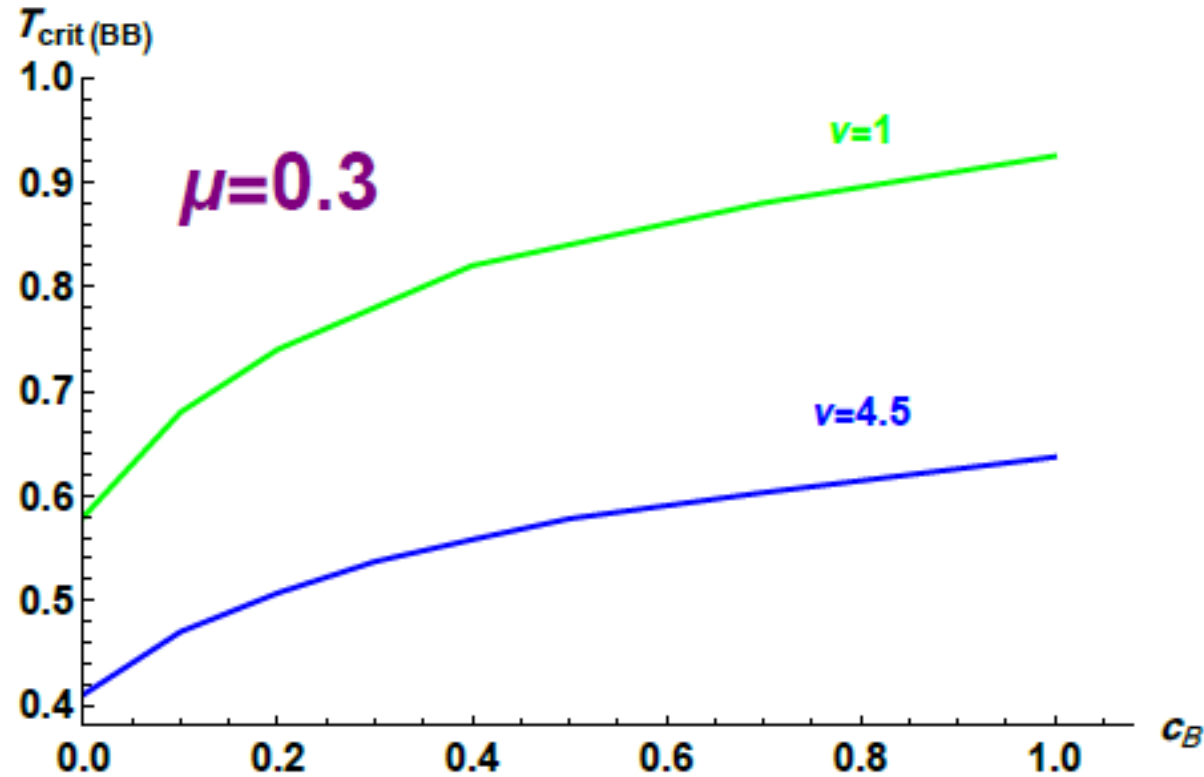
# Free energy:



$$q_3 = 5$$



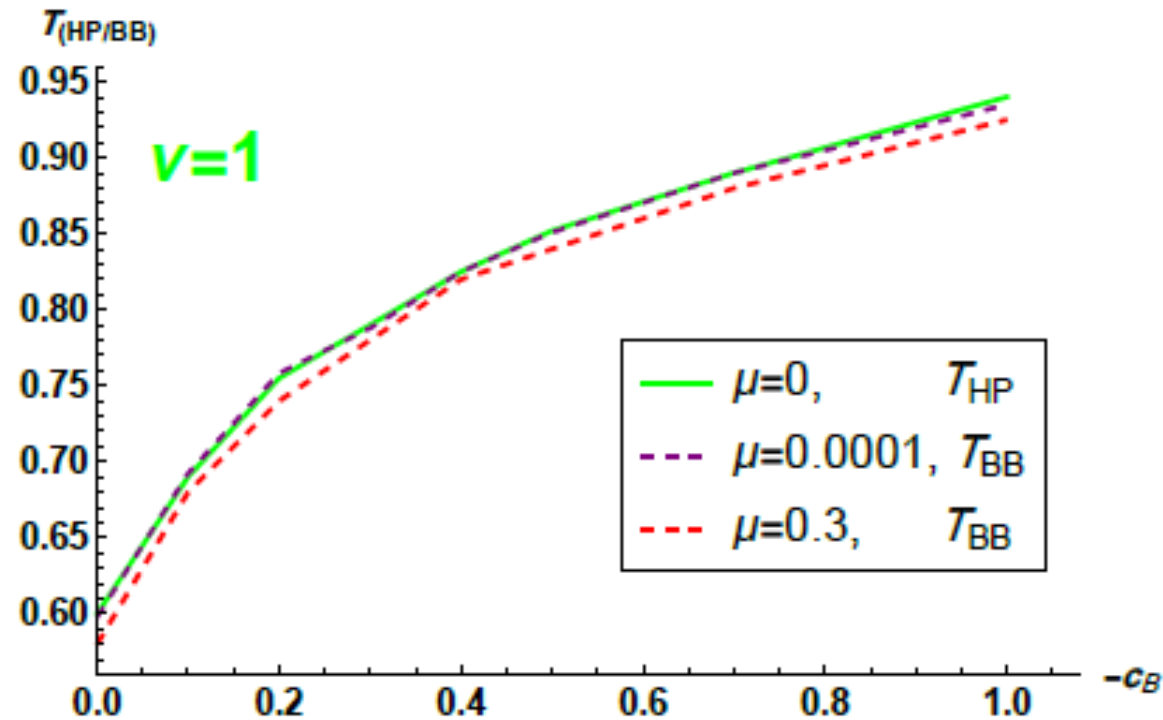
# Phase diagram:



$$q_3 = 5$$

MC phenomenon is obtained.

Phase diagram when:  $\mu \rightarrow 0$

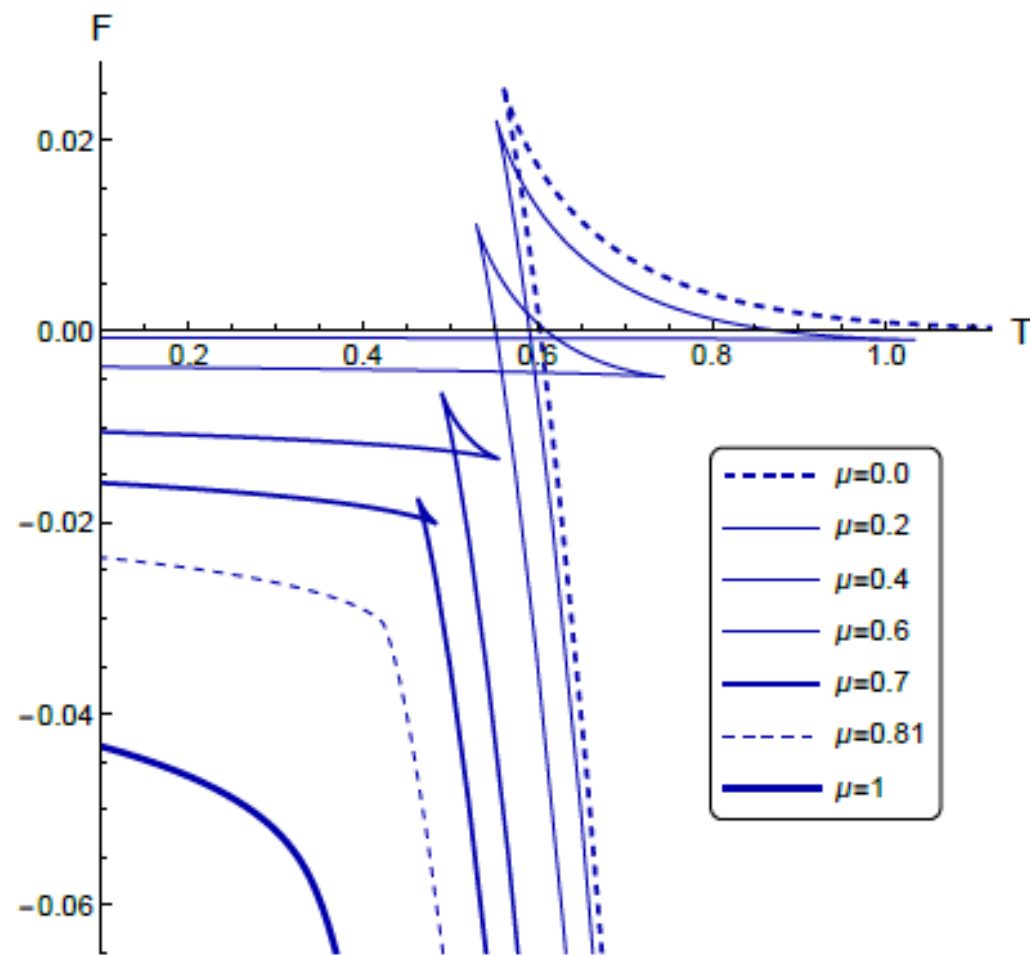
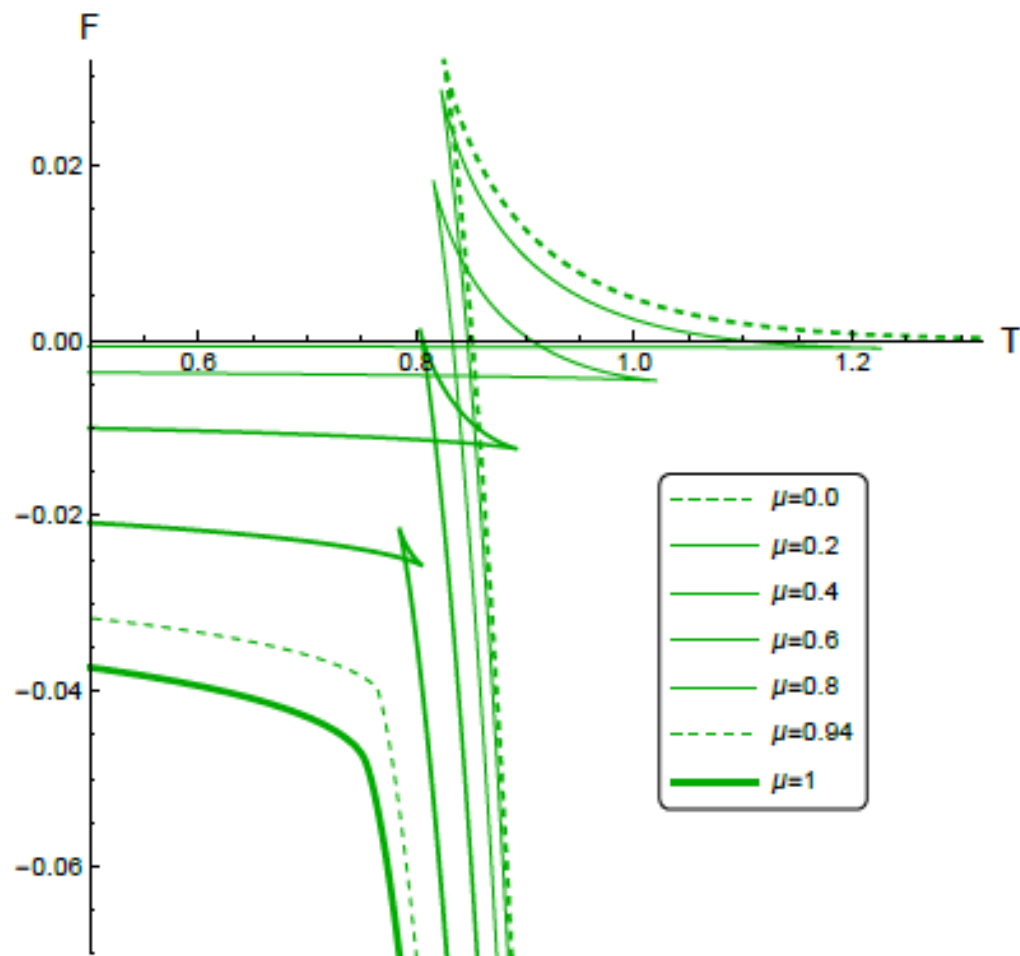


$$q_3 = 5$$

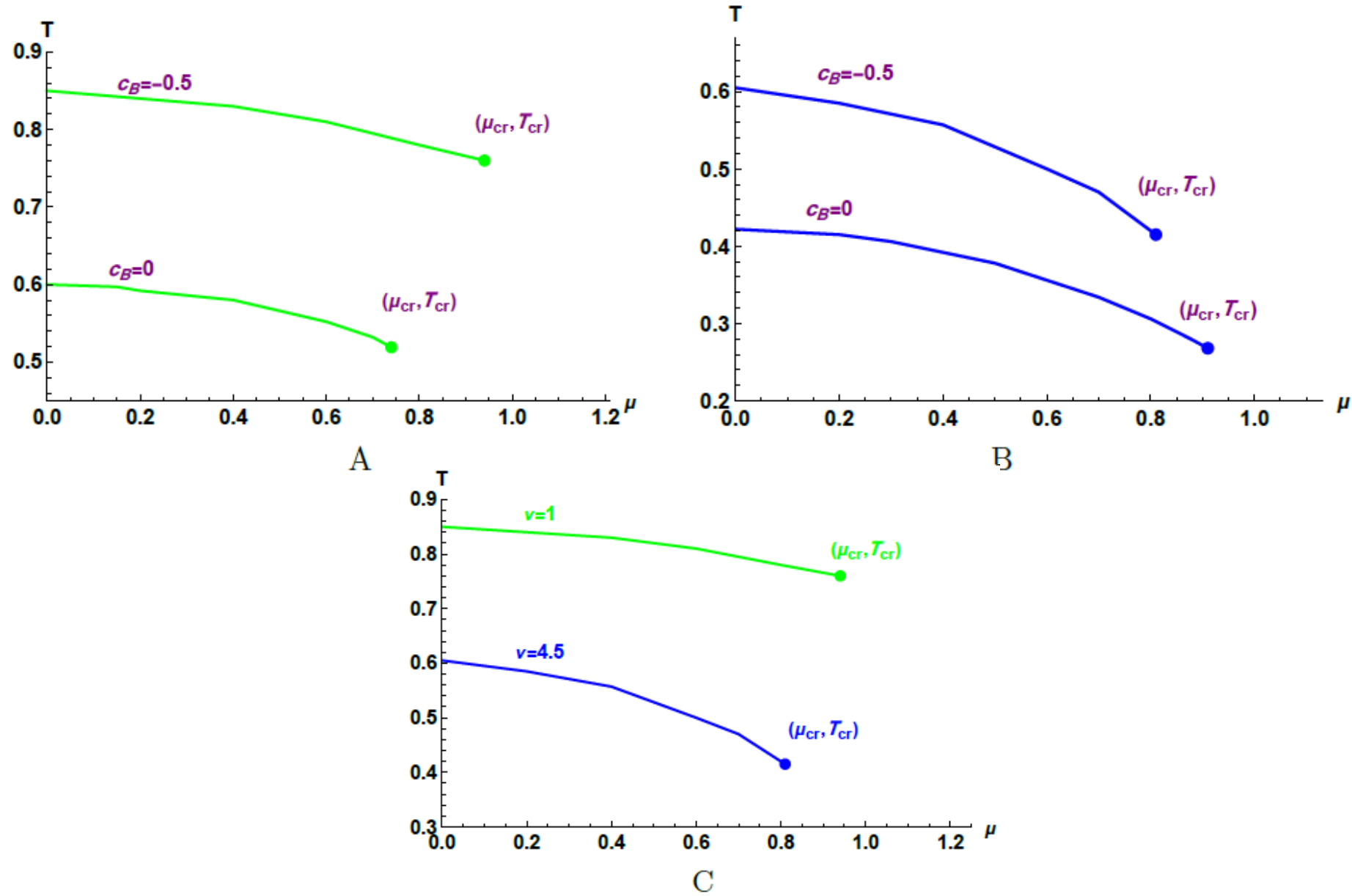
# Free energy:

$$c_B = -0.5$$

$$q_3 = 5$$



# Phase diagram for different cases of anisotropy:



# Summary:

- We developed the model to possess MC phenomenon.
- Increasing the primary anisotropy decreases the confinement/deconfinement and black hole black hole phase transition.
- Although, increasing anisotropy via magnetic field increases the critical transition temperatures.

# Works in progress:

- Renormalization group flow..... (MIAN)
- Unquenched set up for quarks .....(KU Leuven and NIT India)

Thanks...