

On strict monotonicity of the p -torsional rigidity over annuli

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Outline

1 Introduction

- Saint-Venant inequality
- Reverse Saint-Venant inequality

2 Monotonicity of Torsional Rigidity

- Monotonicity via Polarization: Complete Dirichlet boundary
- Monotonicity via Geometry and Shape calculus: Mixed boundary
- Neumann data of $w = u_s \circ \sigma_s - u_s$ on ∂B_R

3 Geometry of the torsion function

- Foliated Schwarz symmetry
- Monotonicity in the affine-radial directions
- Monotonicity in the axial direction

Introduction

For $0 < r < R < \infty$ and $0 \leq s < R - r$, let $\Omega_s = B_R(0) \setminus \overline{B_r(se_1)}$ be an annular domain in \mathbb{R}^d , $d \geq 2$. For $p \in (1, \infty)$, we consider the p -torsion problem in Ω_s :

$$-\Delta_p u = 1 \text{ in } \Omega_s \quad (1a)$$

$$u = 0 \text{ on } \Gamma^s; \quad \frac{\partial u}{\partial n} = 0 \text{ on } \partial\Omega_s \setminus \Gamma^s, \quad (1b)$$

where $\Delta_p u = \operatorname{div}(|\nabla u|^{p-2} \nabla u)$ is the p -Laplace operator, and

$$\text{either } \Gamma^s = \partial\Omega_s \text{ or } \Gamma^s = \partial B_R(0) \text{ or } \Gamma^s = \partial B_r(se_1). \quad (2)$$

- The problem (1a) with (1b) admits a unique solution $u_s \in W_{\Gamma^s}^{1,p}(\Omega_s)$, and it is called the p -torsion function of Ω_s . Further, $0 < u_s \in C^{1,\alpha}(\overline{\Omega})$.
- The p -torsional rigidity $\mathcal{T}(s)$ of Ω_s is defined as

$$\mathcal{T}(s) = \left(\int_{\Omega_s} u_s \, dx \right)^{p-1} = \left(\int_{\Omega_s} |\nabla u_s|^p \, dx \right)^{p-1}. \quad (3)$$

- The p -torsion function u_s and p -torsional rigidity $\mathcal{T}(s)$ are given by the variational characterizations:

$$\max_{u \in W_{\Gamma_s}^{1,p}(\Omega_s)} \int_{\Omega_s} u \, dx - \frac{1}{p} \int_{\Omega_s} |\nabla u|^p \, dx = \frac{p-1}{p} \int_{\Omega_s} u_s \, dx; \quad (4)$$

$$\mathcal{T}(s) = \sup \left\{ \left[\int_{\Omega_s} u \, dx \right]^p : 0 \leq u \in W_{\Gamma_s}^{1,p}(\Omega_s) \text{ with } \int_{\Omega_s} |\nabla u|^p \, dx = 1 \right\}. \quad (5)$$

- Saint-Venant inequality:** for a bounded domain $\Omega \subset \mathbb{R}^d$ with $\Gamma^s = \partial\Omega_s$,

$$\mathcal{T}(\Omega) \leq \mathcal{T}(\Omega^*),$$

where Ω^* is a ball in \mathbb{R}^d with $|\Omega^*| = |\Omega|$.

- ▶ Conjectured by Saint-Venant [1855] for $d = 2 = p$.
- ▶ Pólya [1948]: For the beams with simply connected cross-sections.
- ▶ Pólya-Weinstein [1950]: For the beams with multiply-connected cross-sections.
- ▶ Talenti [1977]: Proved by a symmetrization argument that holds even for $p \in (1, \infty)$ and $d \geq 2$.

Reverse Saint-Venant inequality

- Reverse Saint-Venant inequality:

$$\mathcal{T}(0) < \mathcal{T}(s) \text{ for } 0 < s < R - r,$$

where $\Omega_s = B_R(0) \setminus \overline{B}_r(se_1)$ for $0 \leq s < R - r$.

- ▶ Hersch¹ & Wienberger² [1960's] proved for $d = 2 = p$, using the method of interior parallels.
- ▶ Anoop-Ashok [2020]³ proved for $p \in (1, \infty)$ and $d \geq 2$, using the method of interior parallels.

¹ J. Hersch. *Pacific J. Math.*, 13, 1963.

² Wienberger. *Pacific J. Math.*, 13, 1963.

³ T. V. Anoop and K. Ashok Kumar. *J. Math. Anal. Appl.*, 485(1), 2020.

Monotonicity of Torsional Rigidity

- We are interested in the behavior $s \mapsto \mathcal{T}(s)$ on $[0, R - r)$.
- For the first eigenvalue problem $-\Delta_p u = \gamma_1(s)|u|^{p-2}u$ in Ω_s with (1b):
 - ▶ **For $\Gamma^s = \partial\Omega_s$:** Harrell⁴ & Kesavan⁵ ($p = 2$), and Anoop-Bobkov-Sasi⁶ ($1 < p < \infty$) proved the monotonicity of $\gamma_1(s)$ by showing that

$$\gamma_1'(0) = 0 \text{ and } \gamma_1'(s) < 0 \text{ for } s > 0.$$

- ▶ **For $\Gamma^s = \partial B_r(se_1)$:** Anoop-Ashok-Kesavan⁷ proved for $p = 2$.
- ▶ **For $\Gamma^s \subseteq \partial\Omega_s$ as in (2):** Anoop-Ashok⁸ proved when $p > \frac{2d+2}{d+2}$ for certain translations and rotations, using the methods of polarization.

⁴ E. M. Harrell, II, P. Kröger, and K. Kurata. *SIAM J. Math. Anal.*, 33(1), 2001.

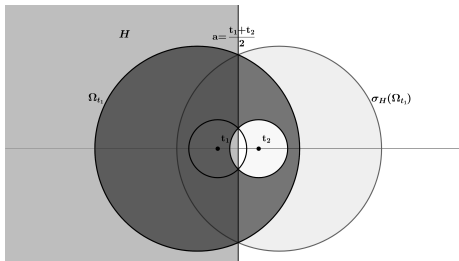
⁵ S. Kesavan. *Proc. Roy. Soc. Edinburgh Sect. A*, 133(3), 2003.

⁶ Anoop, T. V., V. Bobkov, and S. Sasi. *Trans. Amer. Math. Soc.*, 370(10), 2018.

⁷ Anoop, T. V., Ashok Kumar, K., and Kesavan, S. *J. Differ. Equ.* 298, 2021.

⁸ Anoop, T. V. and Ashok Kumar, K. *Adv. Differential Equations* 28(7-8), 2023.

Monotonicity: when $\Gamma^s = \partial\Omega_s$ via Polarization



- Take $0 \leq s < t < R - r$ and $a := \frac{s+t}{2}$. Let $H_a = \{(x_1, x') \in \mathbb{R}^d : x_1 < a\}$ and $\sigma_a(x_1, x') = (2a - x_1, x')$ be the reflection with respect to $\partial H_a = \{x_1 = a\}$.
- For the maximizer $u_s \in W_{\Gamma^s}^{1,p}(\Omega_s)$ of $\mathcal{T}(s)$, \tilde{u}_s is the zero extension to \mathbb{R}^d .

Define

$$P_a \tilde{u}_s(x) = \begin{cases} \max \{ \tilde{u}_s(x), \tilde{u}_s(\sigma_a(x)) \}, & \text{for } x \in H_a, \\ \min \{ \tilde{u}_s(x), \tilde{u}_s(\sigma_a(x)) \}, & \text{for } x \in H_a^c. \end{cases} \quad (6)$$

- $\text{supp}(P_a \tilde{u}_s) \subseteq \Omega_t$, and $v := P_a \tilde{u}_s|_{\Omega_t} \in W_{\Gamma^t}^{1,p}(\Omega_t)$ with $\int_{\Omega_t} v \, dx = \int_{\Omega_s} u_s \, dx$, and $\int_{\Omega_t} |\nabla v|^p \, dx = \int_{\Omega_s} |\nabla u_s|^p \, dx$.

Monotonicity via Polarization

- By the variational characterization (5),

$$\mathcal{T}(t) \geq \left(\int_{\Omega_t} v \, dx \right)^p = \left(\int_{\Omega_s} u_s \, dx \right)^p = \mathcal{T}(s).$$

- For the strict inequality, we use a version of the strong comparison principle.

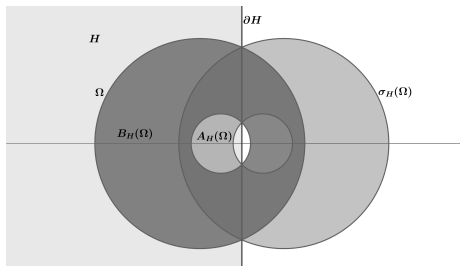
Strong Comparison Principle, Sciunzi⁹

Let $B \subset \mathbb{R}^d$ be a bounded smooth domain, and $\frac{2d+2}{d+2} < p < \infty$. Let $u, v \in \mathcal{C}^1(\overline{B})$ be positive distributional solutions of $-\Delta_p w = 1$ in B . If $v \geq u$ in B , then either $v > u$ in B or $v \equiv u$ in B .

- We find a ball $B \subset \Omega_s \cap H_a$ in which v, u_s satisfy a contradictory properties to the SCP.
- If $\mathcal{T}(s) = \mathcal{T}(t)$, then $v, u_s \in \mathcal{C}^1(\overline{B})$ are positive distributional solutions in B . This leads to a contradiction to the SCP. Thus the strict inequality $\mathcal{T}(s) < \mathcal{T}(t)$ holds.

⁹ Sciunzi, B., *Commun. Contemp. Math.*, 16(6), 2014.

Existence of a ball $B \subset \Omega_s \cap H_a$



- $B_H(\Omega_s) = B_R(0) \setminus B_R(\sigma_a(0))$, and $A_H(\Omega_s) = \overline{B}_r(se_1) \setminus \overline{B}_r(\sigma_a(se_1))$ have non-empty interiors in $\Omega_t \cap H_a$.
- $\Omega_t \cap H_a = \mathcal{M}_{u_s} \sqcup \mathcal{N}_{u_s}$ with $\mathcal{M}_{u_s} = \{x \in \Omega_t \cap H_a : v(x) > u_s(x)\} \supsetneq A_H(\Omega_s)$ and $\mathcal{N}_{u_s} = \{x \in \Omega_t \cap H_a : v(x) = u_s(x)\} \supsetneq B_H(\Omega_s)$.
- $\Omega_t \cap H_a$ is connected, then $\emptyset \neq \mathcal{N}_{u_s} \cap \partial \mathcal{M}_{u_s} \ni x_0$ and $B := B_\rho(x_0) \subset \Omega_s \cap H_a$.
- $v > u_s$ in $B \cap \mathcal{M}_{u_s} \neq \emptyset$ and $v = u$ in $B \cap \mathcal{N}_{u_s} \neq \emptyset$.

Monotonicity of Torsional Rigidity

Theorem

Let $0 < r < R < \infty$ and $p \in (1, \infty)$. For $0 \leq s < R - r$, let $\mathcal{T}(s)$ be the p -torsional rigidity of $\Omega_s = B_R(0) \setminus \overline{B_r}(se_1)$ with $\Gamma_s = \partial\Omega_s$. Then $s \mapsto \mathcal{T}(s)$ is increasing on $[0, R - r)$. Further, if $p > \frac{2d+2}{d+2}$ then $s \mapsto \mathcal{T}(s)$ is strictly increasing on $[0, R - r)$.

Monotonicity of Torsional Rigidity

Main ingredients in proof of Harrell & Kesavan: Let $u > 0$ be the first eigenfunction corresponding to $\lambda_1(s)$ with $\|u\|_2 = 1$.

- ① Shape derivative formula for $\lambda'_1(s)$,

$$\lambda'_1(s) = - \int_{\partial B_r} \left(\frac{\partial u}{\partial n}(x) \right)^2 n_1(x) \, dS(x).$$

Derived using the Hadamard perturbation formula

$$d\lambda_1(\Omega_s; V) = - \int_{\partial\Omega_s} \left(\frac{\partial u}{\partial n}(x) \right)^2 V(x) \cdot n(x) \, dS(x),$$

where V is a smooth vector field on \mathbb{R}^d , and $n(x) = (n_1(x), \dots, n_d(x))$ is an outward normal to Ω_s .

- ② Strong maximum principle and Hopf's boundary point lemma

Shape derivative of $\mathcal{T}(s)$: when $p = 2$

Let $u_s > 0$ be the torsion function of Ω_s with Γ^s as in (2), i.e.,

$$\mathcal{T}(s) = \int_{\Omega_s} u_s \, dx = \int_{\Omega_s} |\nabla u_s|^2 \, dx$$

- Hadamard perturbation formula: For a smooth vector field V on \mathbb{R}^N ,

$$\begin{aligned} d\mathcal{T}(\Omega_s; V) &= \int_{\Gamma^s} \left(\frac{\partial u_s}{\partial n} \right)^2 V \cdot n \, dS \\ &\quad - \int_{\partial\Omega_s \setminus \Gamma^s} \left(|\nabla u_s|^2 - 2u_s \right) V \cdot n \, dS \end{aligned}$$

- Shape derivative formula when $\Gamma^s = \partial B_r(se_1)$: Take V that fixes the outer boundary ∂B_R and translates the inner ball $B_r(se_1)$ along e_1 -axis:
 $V(x) = \rho(x)e_1$ with $0 < \rho \in \mathcal{C}_c^\infty(B_R(0))$ with $\rho \equiv 1$ in a neighborhood of $B_r(se_1)$.

$$\mathcal{T}'(s) = \int_{\partial B_r} \left(\frac{\partial u_s}{\partial n}(x) \right)^2 n(x) \cdot e_1 \, dS(x).$$

Monotonicity of $\mathcal{T}(s)$: Approach of Harrell & Kesavan

- ① Rewriting the shape derivative formula:

$$\mathcal{T}'(s) = \int_{\partial B_r \cap \{x_1 > s\}} \left[\left(\frac{\partial u_s}{\partial n}(x) \right)^2 - \left(\frac{\partial u_s}{\partial n}(\sigma_s(x)) \right)^2 \right] n_1(x) \, dS(x),$$

where $\sigma_s(x)$ is the reflection with respect to the affine-hyperplane $\{x_1 = s\}$.

- ② Observe $\frac{\partial u_s}{\partial n}(\sigma_s(x)), \frac{\partial u_s}{\partial n}(x) < 0$ (by Hopf's), and $n_1(x) < 0$ for $x \in \partial B_r \cap \{x_1 > s\}$.

- ③ If

$$\frac{\partial u_s}{\partial n}(\sigma_s(x)) - \frac{\partial u_s}{\partial n}(x) < 0 \text{ for } x \in \partial B_r \cap \{x_1 > s\}, \quad (7)$$

then we get $\mathcal{T}'(s) > 0$.

Monotonicity of $\mathcal{T}(s)$

- Consider $w = u_s \circ \sigma_s - u_s$ on $\widetilde{\Omega}_s = \Omega_s \cap \{x_1 > s\}$. Then w satisfies:

$$-\Delta w = 0 \text{ in } \widetilde{\Omega}_s.$$

- On $H_s \cup (\partial B_r \cap \{x_1 > s\})$: $w(x) = u_s \circ \sigma_s(x) - u_s(x) = u_s(x) - u_s(x) = 0$.
- On $\partial B_R \cap \{x_1 > s\}$: $w = ?$ or $\frac{\partial w}{\partial n} = ?$
- Since $\frac{\partial u_s}{\partial n} = 0$ on $\partial B_R(0)$, it is natural to expect the sign of $\frac{\partial w}{\partial n}$ on $\partial B_R(0)$.
- Here, the finer geometric properties of the torsion function help us to confirm

$$\frac{\partial w}{\partial n} \geq 0 \text{ on } \partial B_R \cap \{x_1 > s\}.$$

Neumann data of $w = u_s \circ \sigma_s - u_s$ on ∂B_R

- Observe that, for $x \in \partial B_R$,

$$\frac{\partial w}{\partial n}(x) = \frac{\partial u_s \circ \sigma_s}{\partial n}(x) - \frac{\partial u_s}{\partial n}(x) = \frac{\partial u_s \circ \sigma_s}{\partial n}(x).$$

- Recall that, for $x = (x_1, x') \in \mathbb{R}^d$,

$$\sigma_s(x) = (2s - x_1, x') = 2se_1 + \sigma_0(x).$$

- So $\nabla(u_s \circ \sigma_s)(x) = \nabla u_s(\sigma_s(x))A_0$, where $A_0 = \text{diag}(-1, 1, \dots, 1)$ is the reflection matrix with respect to $\{x_1 = 0\}$.

Neumann data of $w = u_s \circ \sigma_s - u_s$ on ∂B_R

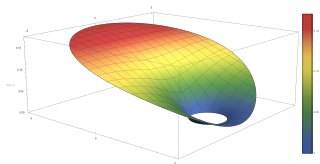
- For $x \in \partial B_R$, the unit outward normal is $n(x) = \frac{x}{R}$. Then

$$\begin{aligned} R \nabla(u_s \circ \sigma_s)(x) \cdot n(x) &= \nabla u_s(\sigma_s(x)) A_0 \cdot x = \nabla u_s(\sigma_s(x)) \cdot \sigma_0(x) \\ &= \nabla u_s(\sigma_s(x)) \cdot (\sigma_s(x) - 2se_1). \end{aligned}$$

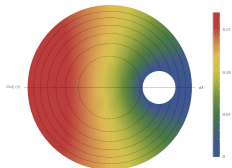
- Write

$$R \frac{\partial w}{\partial n}(x) = \varphi(\sigma_s(x)) - s \frac{\partial u_s}{\partial x_1}(\sigma_s(x)),$$

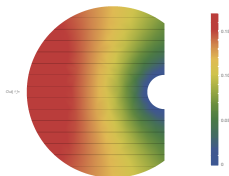
where $\varphi(x) = \nabla u_s(x) \cdot (x - se_1)$ for $x \in \Omega_s$.



(a)



(b)



(c)

Figure: (a) The torsion function of $B_5(0) \setminus \overline{B_1(3e_1)} \subset \mathbb{R}^2$. (b) The foliated Schwartz symmetry, (c) Monotonicity along the e_1 -direction up to $x < 3$.

Theorem

Let $\Omega_s = B_R(0) \setminus \overline{B_r(se_1)}$ be an annular domain with $0 < s < R - r$. Let $u_s > 0$ be the torsion function and $\mathcal{T}(s)$ of Ω_s with $\Gamma^s = \partial B_r(se_1)$ be the torsional rigidity. Then

- i u_s has the foliated Schwarz symmetry in Ω_s with respect to $-\mathbb{R}^+ e_1$;
- ii u_s is strictly increasing along all affine-radial directions from se_1 in Ω_s , i.e.,

$$\nabla u_s(x) \cdot (x - se_1) > 0 \text{ for } x \in \overline{\Omega_s} \setminus \{\pm Re_1\};$$

- iii u_s is strictly decreasing in the e_1 -direction on the sub-region $\{x \in \Omega_s : x_1 < s\}$ of Ω_s , i.e.,

$$\frac{\partial u_s}{\partial x_1} < 0 \text{ on } \{x \in \Omega_s : x_1 < s\}.$$

- The affine-radial monotonicity of u_s implies that $\varphi > 0$ in Ω_s ,
- The axial monotonicity of u gives that

$$\frac{\partial u_s}{\partial x_1} < 0 \text{ for } x \in \Omega_s \cap \{x_1 < s\}.$$

- Therefore, for $x \in \partial B_R \cap \{x_1 > s\}$,

$$\frac{\partial w}{\partial n}(x) = \frac{1}{R} \varphi(\sigma_s(x)) - \frac{s}{R} \frac{\partial u_s}{\partial x_1}(\sigma_s(x)) \geq 0.$$

Monotonicity of $\mathcal{T}(s)$

- The map w satisfies the following boundary value problem on $\widetilde{\Omega}_s := \Omega_s \cap \{x_1 > s\}$:

$$\left. \begin{aligned} -\Delta w &= 0 \text{ in } \widetilde{\Omega}_s, \\ w &= 0 \text{ on } H_s \cup (\partial B_r \cap \{x_1 > s\}), \\ \frac{\partial w}{\partial n} &\geq 0 \text{ on } \partial B_R \cap \{x_1 > s\}, \end{aligned} \right\} \quad (8)$$

- We derived a variant of strong maximum principle that holds for the boundary value problem above.

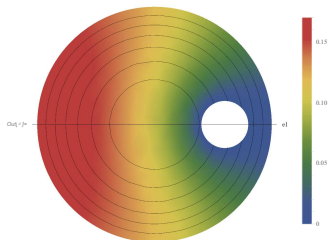
Monotonicity of $\mathcal{T}(s)$

- Since $\frac{\partial w}{\partial n} > 0$ on $\partial B_R \setminus \{\pm Re_1\}$, SMP implies that $w > 0$ in $\widetilde{\Omega}_s$.
- Hopf's lemma implies that $\frac{\partial w}{\partial n} < 0$ on $\partial B_r \cap \{x_1 > s\}$ since $w = 0$.
- For $x \in \partial B_r \cap \{x_1 > s\}$,

$$0 > \frac{\partial w}{\partial n}(x) = \frac{\partial u_s}{\partial n}(\sigma_s(x)) - \frac{\partial u_s}{\partial n}(x). \quad (9)$$

- Therefore $\mathcal{T}'(s) > 0$.

Foliated Schwarz symmetry of u



- Let $-\mathbb{R}^+ e_1$ be the half-ray from 0 in $-e_1$ -direction.
- Let (r, θ) be the polar-coordinates given by $-\mathbb{R}^+ e_1$ with

$$r(x) = |x| \text{ and } \cos(\theta(x)) = -\frac{x \cdot e_1}{|x|}, \quad \theta \in [0, \pi]. \quad (10)$$

- A map $0 \leq u_s : \Omega_s \rightarrow \mathbb{R}$ is foliated Schwarz symmetric in Ω_s with respect to $-\mathbb{R}^+ e_1$, if and only if
 - ① u depends only on (r, θ) ,
 - ② for fixed $r > 0$, u decreases in θ . Equivalently, for $x, y \in \Omega_s$ with $|x| = |y| = r$ and $\theta(x) \geq \theta(y)$ we have $u(x) < u(y)$.

Foliated Schwarz symmetrization via Polarization

- Let H be an open affine half-space in \mathbb{R}^N such that $-e_1 \in H$ and $se_1 \in \partial H$, and let σ_H be the reflection with respect to ∂H .
- For $0 \leq u : \mathbb{R}^d \rightarrow \mathbb{R}$, the polarization $P_H(u) : \mathbb{R}^d \rightarrow \mathbb{R}$ is defined as

$$P_H(u)(x) = \begin{cases} \max\{u(x), u(\sigma_H(x))\}, & \text{for } x \in H, \\ \min\{u(x), u(\sigma_H(x))\}, & \text{for } x \in H^c. \end{cases} \quad (11)$$

For $0 \leq u : \Omega_s \rightarrow \mathbb{R}$ let \tilde{u} be the zero extension of u to \mathbb{R}^d . The polarization of u is defined as $P_H(u) := P_H(\tilde{u})|_{\Omega_s}$.

- If $u \in H^1(\Omega_s)$ with $u|_{\partial B_r} = 0$, then $P_H(u) \in H^1(\Omega_s)$ with $P_H(u)|_{\partial B_r} = 0$ and

$$\|u\|_{2,\Omega} = \|P_H(u)\|_{2,\Omega_s} \text{ and } \|\nabla u\|_{2,\Omega} = \|\nabla P_H(u)\|_{2,\Omega_s}.$$

- The function u has foliated Schwarz symmetry with respect to $-\mathbb{R}e_1$ in Ω_s iff $P_H(u) = u$ for any affine half-space H with $-e_1 \in H$ and $se_1 \in \partial H$.
- Take $0 < u_s$ to be the torsion function. By the uniqueness,

$$P_H(u_s) = u_s \text{ in } \Omega_s, \quad \forall H.$$

- The foliated Schwarz symmetry of u_s gives that: For any $x \in \overline{\Omega_s}$ and $\eta \in S^{d-1}$ with $\eta \cdot x = 0$,

$$\nabla u_s(x) \cdot \eta > 0 \text{ iff } \eta_1 < 0.$$

- Neumann condition $0 = R \frac{\partial u_s}{\partial n} = \nabla u(x) \cdot x$ for $x \in \partial B_R(0)$. Since $|\nabla u_s| \neq 0$, we get $\frac{\partial u_s}{\partial x_1} < 0$ on $\partial B_R(0) \setminus \{\pm Re_1\}$.

Monotonicity in the affine-radial directions

- For $x \in \mathbb{R}$, the vector $x - se_1$ is called as the affine-radial direction from se_1 .
- The map $\varphi(x) := \nabla u_s(x) \cdot (x - se_1)$ for $x \in \Omega_s$ satisfies:

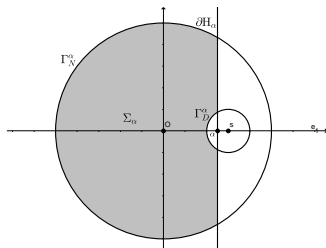
$$-\Delta\varphi = 2u_s > 0 \text{ in } \Omega_s. \quad (12)$$

- On $\partial B_R(0)$: $\varphi(x) = R\nabla u_s(x) \cdot \frac{x}{R} - s\frac{\partial u_s}{\partial x_1} \geq 0$.
- For $\partial B_r(se_1)$: The vector $x - se_1$ is an inward direction to Ω_s and $u_s(x) = 0$. Hopf's boundary lemma gives that $\varphi(x) \geq 0$. Therefore

$$\varphi(x) \geq 0 \text{ on } \partial\Omega_s.$$

- By SMP, $\varphi > 0$ in Ω_s .

Monotonicity in the axial direction

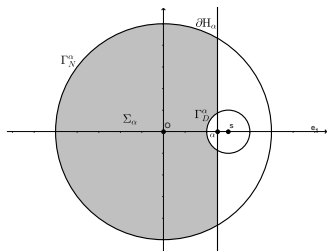


- For $\alpha \in (-R, R)$, the α -cap Σ_α is

$$\Sigma_\alpha := \Omega_s \cap \{x_1 < \alpha\}.$$

- First, by fSS of u_s , observe that: for $\alpha \leq 0$

$$\frac{\partial u_s}{\partial x_1} < 0 \text{ in } \overline{\Sigma_\alpha} \setminus \{-Re_1\}. \quad (*)$$



- $(-R, 0] \subset \mathcal{A} := \{\alpha \in (-R, R) : (*) \text{ holds}\}.$

-

$$\frac{\partial u_s}{\partial x_1} < 0 \text{ on } \partial B_r \cap \{x_1 < s\}, \text{ and } \frac{\partial u_s}{\partial x_1} \geq 0 \text{ on } \partial B_r \cap \{x_1 \geq s\}$$

- Smoothness of u_s , and compactness of $\Omega_s \cap \{x_1 = \alpha\}$: for any $\alpha \in \mathcal{A}$ with $\alpha < s$ there exists $\epsilon > 0$ such that $(*)$ holds for $\alpha + \epsilon$. Therefore

$$\sup \mathcal{A} = s.$$

